

Physics 212 – Statistical Mechanics

Quiz #1 - Solution

1. (a) From the partition function

$$Z = \sum_s e^{-\beta H[s]} = \sum_s e^{\beta J \sum_i S_i S_{i+1} + \beta H \sum_i S_i}$$

we have

$$M = \mu \sum_i \langle S_i \rangle = \frac{\sum_s e^{-\beta H} \mu \sum_i S_i}{\sum_s e^{-\beta H}}$$

Then

$$\chi = \frac{\sum_s e^{-\beta H} \mu \sum_i S_i \cdot \beta \mu \sum_j S_j}{\sum_s e^{-\beta H}} - \frac{(\sum_s e^{-\beta H} \mu \sum_i S_i)(\sum_s e^{-\beta H} \beta \mu \sum_j S_j)}{(\sum_s e^{-\beta H})^2}$$

so that

$$\chi = \beta \mu^2 \sum_{ij} (\langle S_i S_j \rangle - \langle S_i \rangle \langle S_j \rangle)$$

In the high-temperature phase,  $\langle S_j \rangle = 0$ . Also,  $\langle S_i S_j \rangle$  is translation invariant,

$$\langle S_i S_j \rangle = \langle S_{i-j} S_0 \rangle$$

Then

$$\chi = \beta \mu^2 N \sum_j \langle s_j s_0 \rangle$$

(b) As shown in class, the partition function at  $H = 0$  is given by

$$\begin{aligned} Z &= \sum_s \prod_{\text{bonds}} \cosh(\beta J) (1 + s_i s_{i+r} \tanh \beta J) \\ &= 2^N (\cosh \beta J)^{2N} \sum_{n=0}^{\infty} b_n z^n \end{aligned}$$

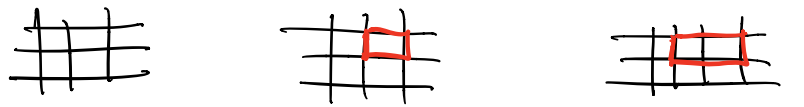
where  $b_n$  is the number of colorings of the lattice with  $n$  colored bonds such that there is an even number of colored bonds at each vertex. Also

$$\sum_s e^{-\beta H} s_j s_0 = 2^N (\cosh \beta J)^{2N} \sum_{n=0}^{\infty} c_n^j z^n$$

where  $c_n^j$  is the the number of colorings of the lattice with  $n$  colored bonds such that there is an even number of colored bonds at each vertex – except that, at the vertices 0 and  $j$ , there are an odd number of colored bonds. We wish to compute

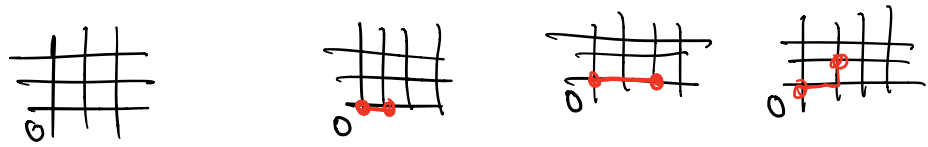
$$\chi = \beta \mu^2 \frac{\sum_{n=0}^{\infty} \sum_j c_n^j z^n}{\sum_{n=0}^{\infty} b_n z^n}$$

The value of the denominator in this expression is



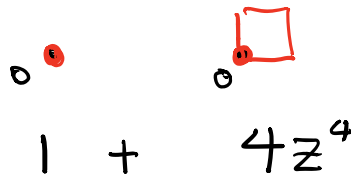
$$1 + N z^4 + O(z^6)$$

The first few terms in the numerator are

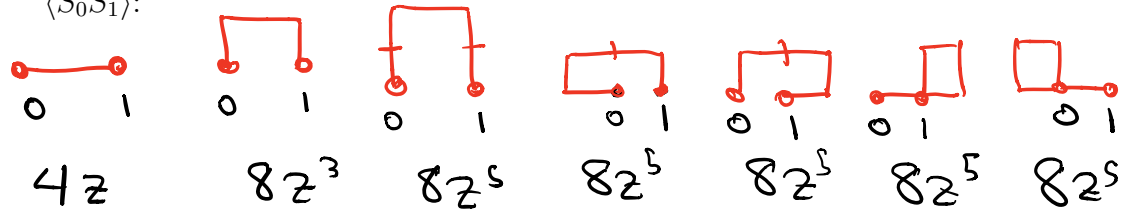


$$1 + 4z + 4 \cdot 3z^2 + \dots$$

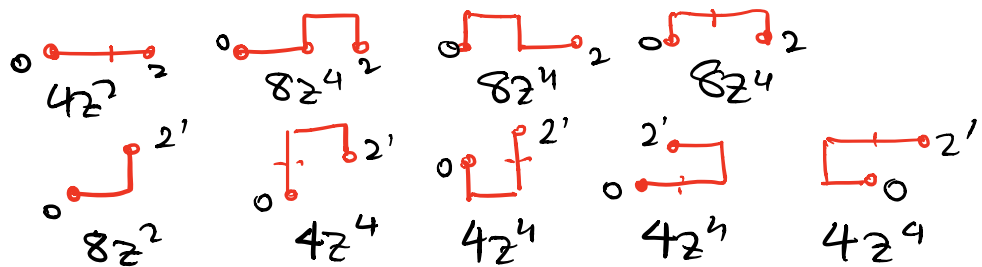
There are many ways to count the colorings in higher orders. Brute force is not so painful. Here are the diagrams that contribute to  $\langle S_0 S_0 \rangle$  through order  $z^5$ :



plus reflections; the total number of diagrams of each type is indicated by the number written below. Here are the diagrams that contribute to  $\langle S_0 S_1 \rangle$ :



Here are the diagrams that contribute to  $\langle S_0 S_2 \rangle$  and  $\langle S_0 S_{2'} \rangle$ :



We see that the diagrams resemble random walks starting at 0. Also the diagrams with  $j$  separated from 0 by  $n$  or more steps contain no closed loops up to order  $z^{n+4}$ .

Given this understanding, here is an easier way to sum the diagrams. Let's refer to walks starting from 0 and not backtracking as "snakes". The number of snakes with  $n$  steps is

$$1, 4, 4 \cdot 3, 4 \cdot 3^2, 4 \cdot 3^3, 4 \cdot 3^4, \dots$$

$$= 1, 4, 12, 36, 108, 324, \dots$$

The snakes include all allowed diagrams, but some diagrams appear two or more times. This can only happen if the diagrams contain closed loops. The first example appears in order  $z^4$ . The loop diagram above appears in the snakes



and so is overcounted by 4. At order  $z^4$  there are also diagrams contributing to  $\langle S_0 S_0 \rangle$  in which the closed loop does not touch the point 0.



In all, the coefficient of  $z^4$  is

$$108 - 4 + N - 4 = 100 + N$$

We can apply the same logic to  $n = 5$ . The snakes



do not correspond to any allowed coloring, an overcount of 8. The pairs of snakes



yield the same coloring. This is an overcount of  $2 \cdot 8 = 16$ . There are also  $4(N-6)$  disconnected diagrams with a square that does not touch the line emerging from 0



In all, the value of  $\sum_j c_5^j$  is

$$324 - 8 - 16 + 4N - 24 = 276 + 4N$$

In all

$$\chi/\beta N^2 = \frac{1 + 4z + 12z^2 + 36z^3 + (100+N)z^4 + (276+4N)z^5}{1 + Nz^4 + \dots}$$

or

$$\chi/\beta N^2 = 1 + 4z + 12z^2 + 36z^3 + 100z^4 + 276z^5 + \dots$$

The  $N$ 's cancel from the coefficients, as they must.

(c) The Taylor expansion of  $A/(z_* - z)^\gamma$  is

$$\frac{A}{(z_* - z)^\gamma} = \frac{A}{z_*^\gamma} \frac{1}{(1 - z/z_*)}^\gamma = \frac{A}{z_*^\gamma} \left( 1 + \gamma \frac{z}{z_*} + \frac{\gamma(\gamma+1)}{2!} \left(\frac{z}{z_*}\right)^2 + \dots \right)$$

Then the ratio of successive coefficients is

$$\frac{a_n}{a_{n-1}} = \frac{\gamma + n - 1}{n} \frac{1}{z_*}$$

Asymptotically

$$\frac{a_n}{a_{n-1}} \rightarrow \frac{1}{z_*}$$

(d) Trying this on the series coefficients in (b), the ratios of successive series coefficients are

$$\frac{a_n}{a_{n-1}} = 4, 3, 3, 2.78, 2.76, \dots$$

This points to

$$\frac{1}{z_*} \approx 2.8$$

or

$$\tanh(\beta J) \approx 0.39 \qquad \frac{T_c}{J} \approx 2.7$$

In principle, if we had enough terms in the series, we could use

$$\frac{a_{n-1}}{a_{n-2}} - \frac{a_n}{a_{n-1}} = \frac{\gamma-1}{n(n-1)}$$

to estimate  $\gamma$ . The actual value of  $\gamma$  for the 2-dimensional Ising model is  $7/4 = 1.75$ . This implies that the ratios of successive terms should increase asymptotically. This short series does not yet show that behavior.