

Particle Detectors

Lecture 7 Physics 152

Lance Dixon

(thanks again to Colin Jessop)

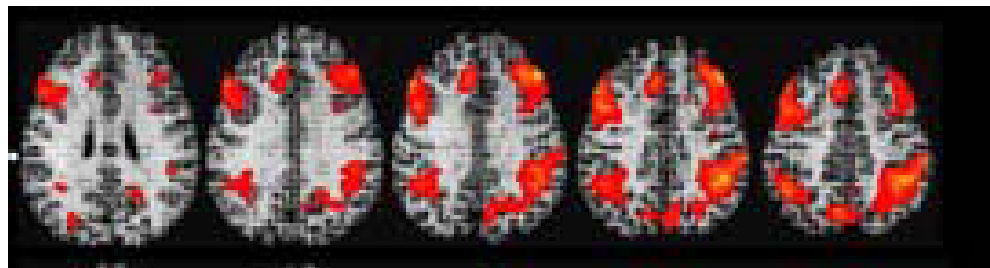
Particle Detectors

Technology developed for particle physics, but has many other applications. For example:

Medical Imaging: X-rays, PET Scans

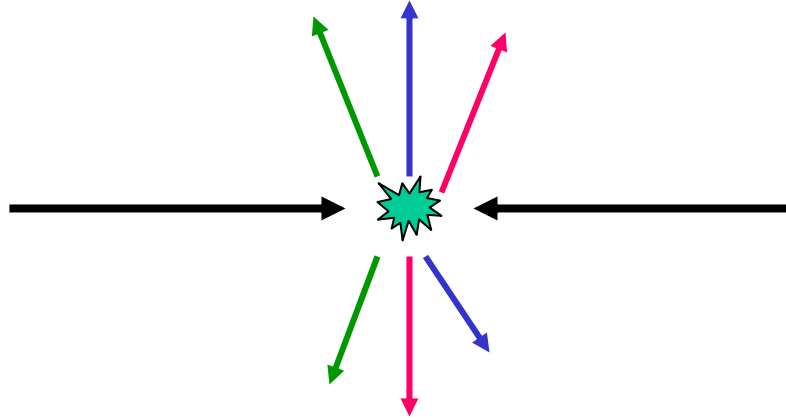
Structure determination in biology and chemistry

Airport Scanners: Explosive detection



A positron emission tomography (PET) scan of brain activity while remembering

Particle Physics Events

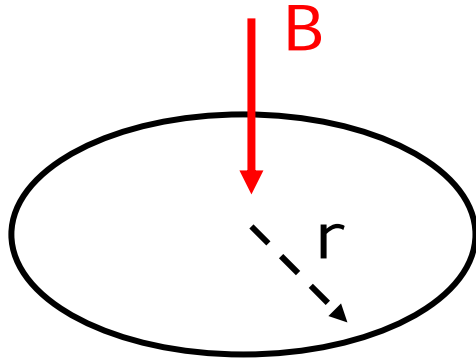


Want to measure 4-vector and identity of all particles produced

Many different particles are produced but **all decay rapidly** to

electrons, photons, protons, neutrons, pions, kaons, muons, and neutrinos

Momentum of Charged Particle



$$p = qrB$$

Charged particle is bent in a magnetic field.

Measuring the trajectory (r) of a charged particle in a known magnetic field gives the momentum - called "Tracking"

Energy Loss of Heavy Charged Particle in Matter

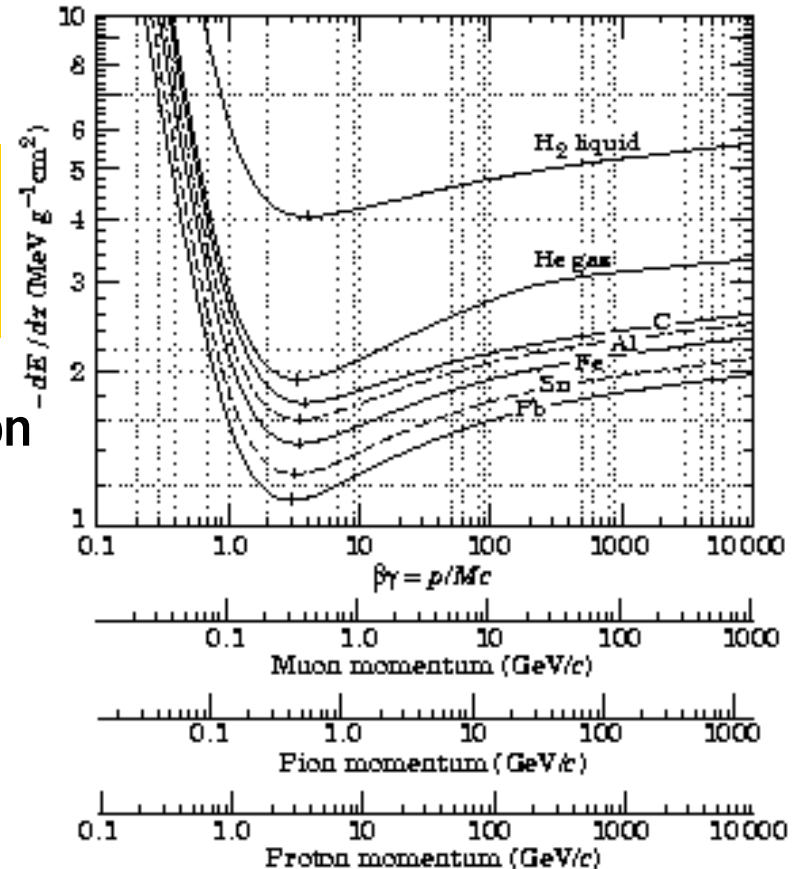
Moderately relativistic charged particles (**except electrons**) lose energy primarily by ionization as they traverse matter.

Bethe-Bloch equation

$$-\frac{dE}{dx} = Kz^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\ln \left(\frac{2m_e c^2 \beta^2 \gamma^2}{I} \right) - \beta^2 - \frac{\delta}{2} \right]$$

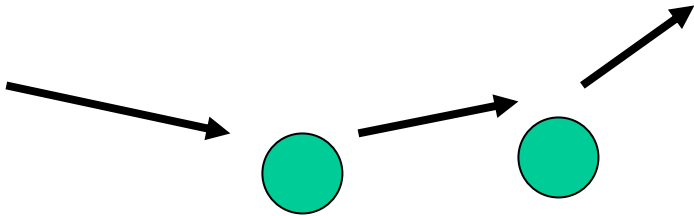
Energy loss per unit distance x is function of velocity and not mass (i.e. same for all particles)

charge z scattering off atom:
 atomic no. Z , atomic mass A ,
 mean ionization potential I .
 K and δ are constants



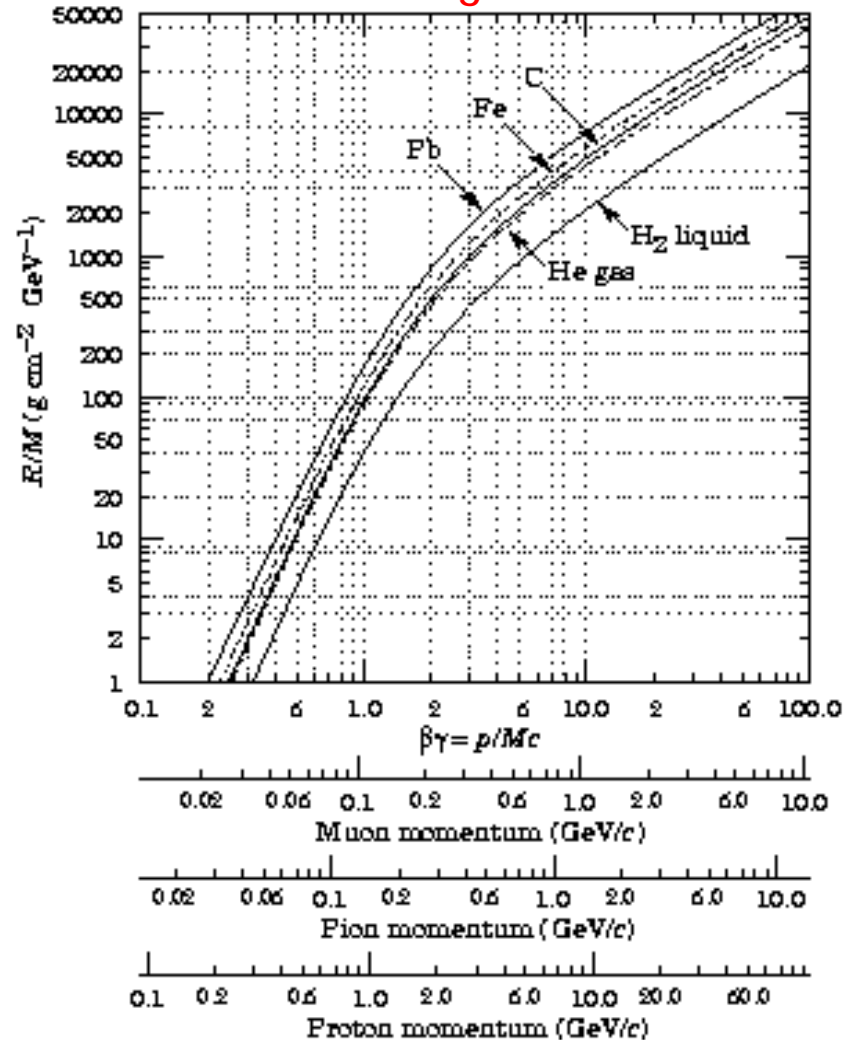
Multiple Scattering

Particle is deflected after many coulomb scatters

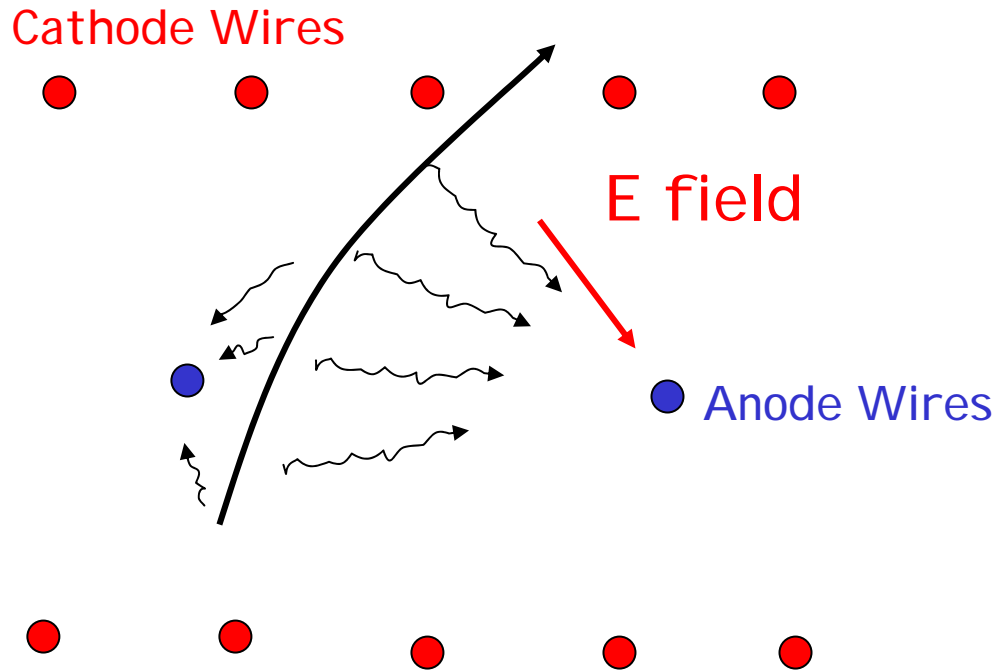


To measure trajectory of a particle, measure its ionization trail as it passes through a medium with low mass (to minimize multiple scattering) inside a magnetic field.

Range



Drift Chambers



Ionized Electrons drift to anodes at constant velocity if voltage is high enough

Drift Chambers

Point of ionization is determined from

$$v_{drift} \ll c$$

$$\sim \text{cm}/\mu\text{s}$$

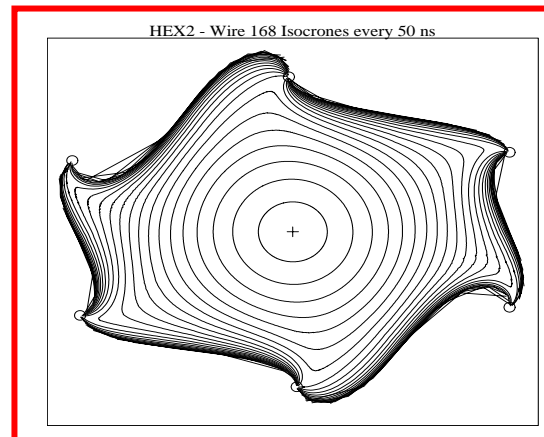
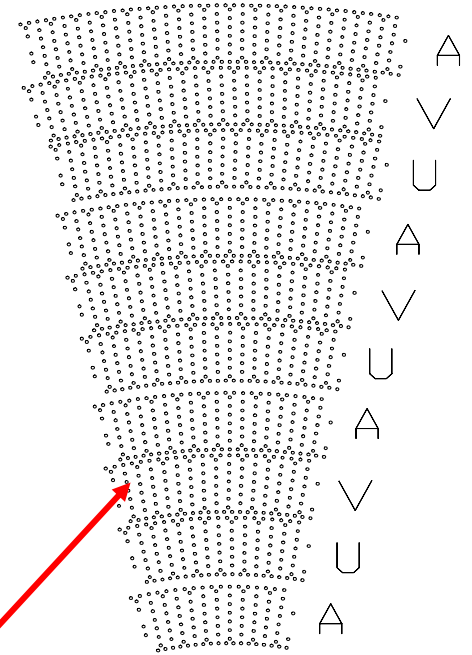
$$x_{drift} = \frac{v_{drift}}{t_{drift}}$$

Drift time is difference between anode pulse arrival time, and particle creation time.

Many layers of wires formed into **cells**

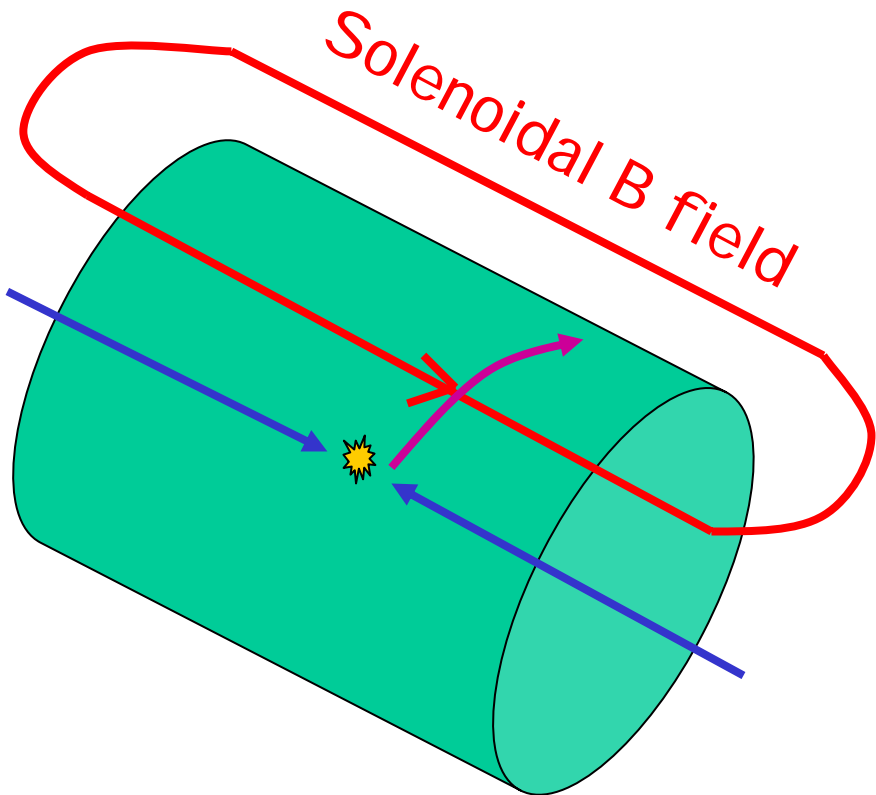
Helium/I sobutane Gas

I sobutane reduces diffusion of electrons as they drift



Drift Cell (isochrones are perp. To field)

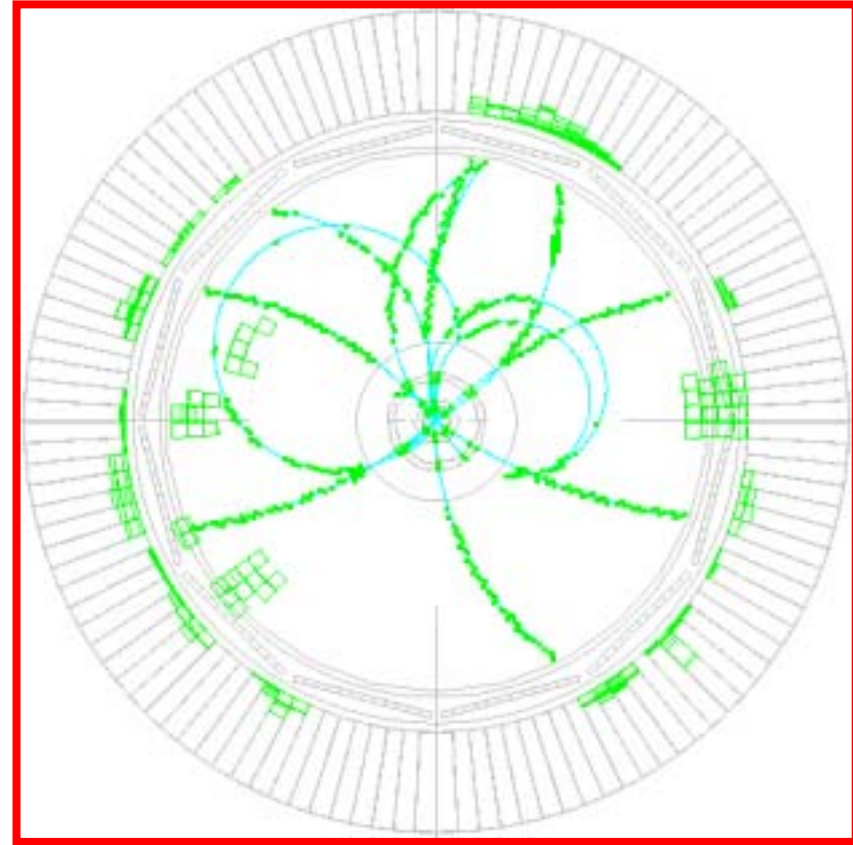
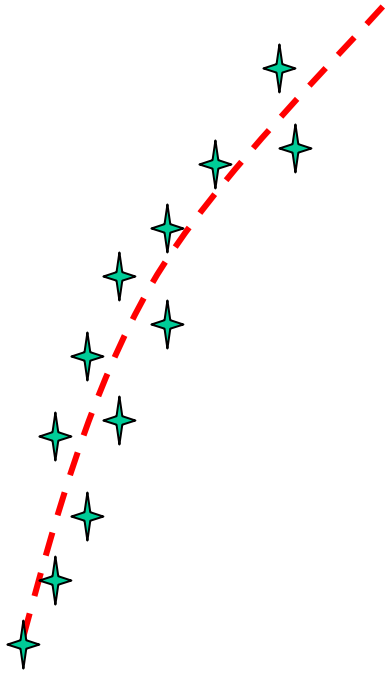
Drift Chambers



40,000 wires in BaBar drift Chamber with B field parallel to beam axis

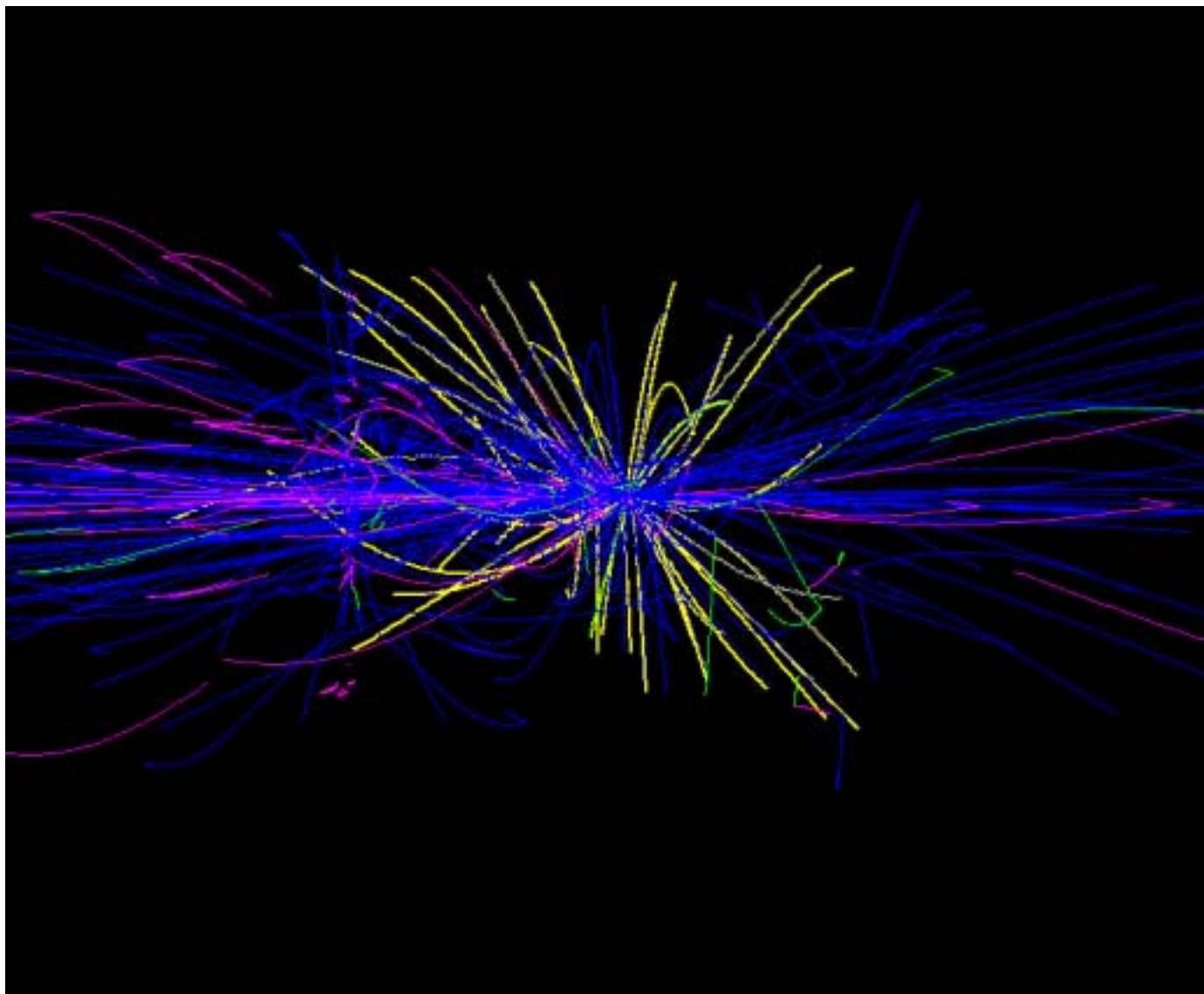
Drift Chambers

Pattern Recognition to find a "road" of hits



χ^2 fit to helix with correction for dE/dx to measure momentum.
Precision of 80 microns achievable

Tracks in 3d from a proton-antiproton collision

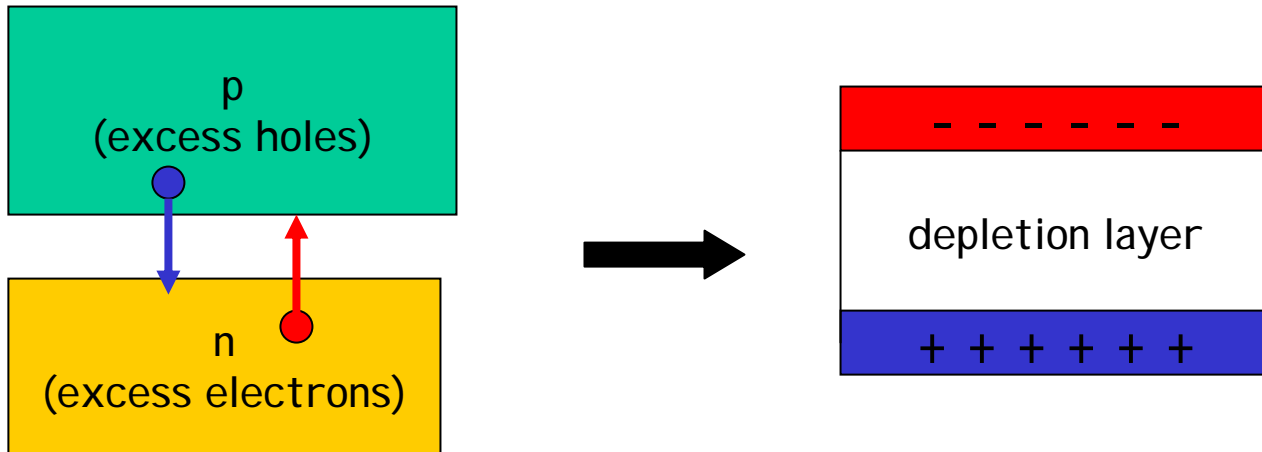


Solid State Tracking Devices

Silicon Devices have been developed in the last decade.
Increase spatial precision from $O(100 \mu\text{m})$ to $O(10 \mu\text{m})$

Ionization detection in silicon (3.6 eV per electron hole pair)
rather than gas (30 eV ionization potential)

p doped (excess holes) in contact with n doped (excess electrons)
forms a **diode**. Apply a **reverse-bias voltage**, $V_n > V_p$.
Very little current flow (few electrons in p doped region).

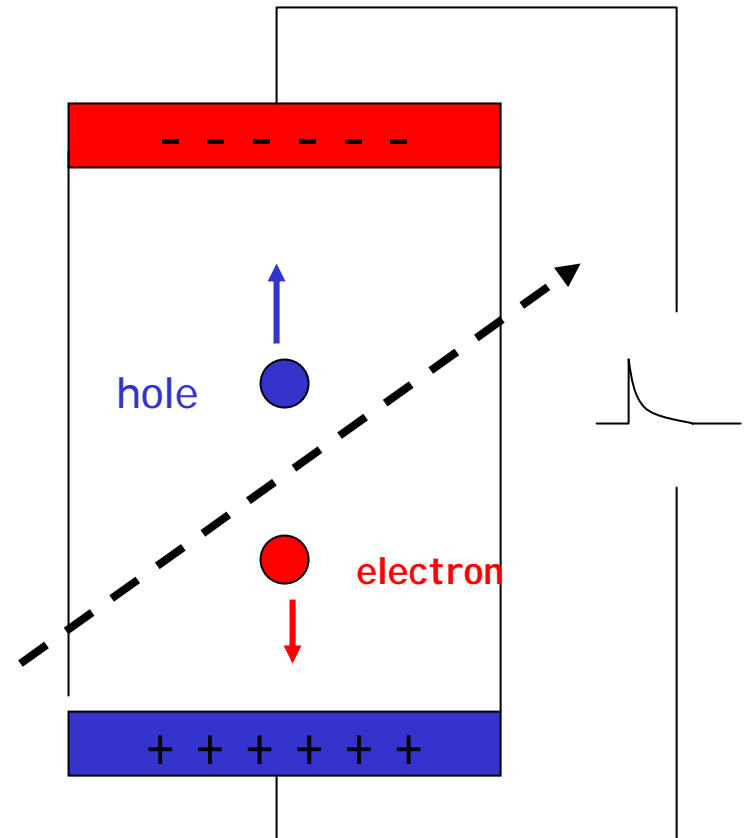


Solid State Tracking Devices

Ionizing particle creates electron/hole pair

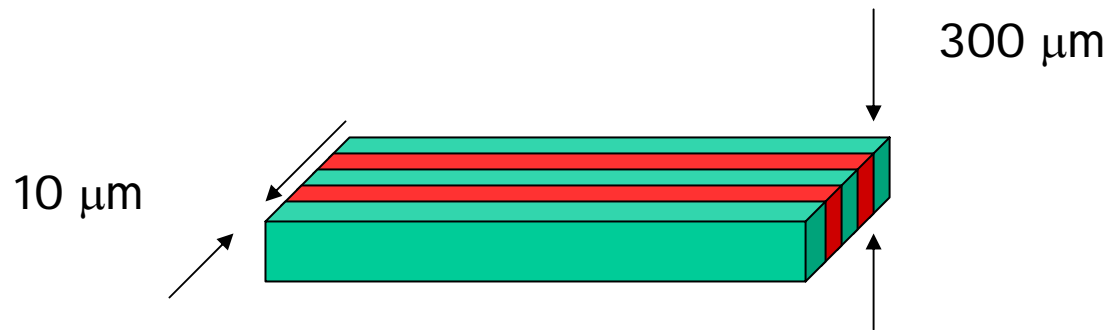
Electron/hole drift to the -ve (p) /+ve (n) surfaces, where they are majority carriers and can escape to electrodes.

If connect to electrical circuit get a pulse of charge



Solid State Tracking Devices

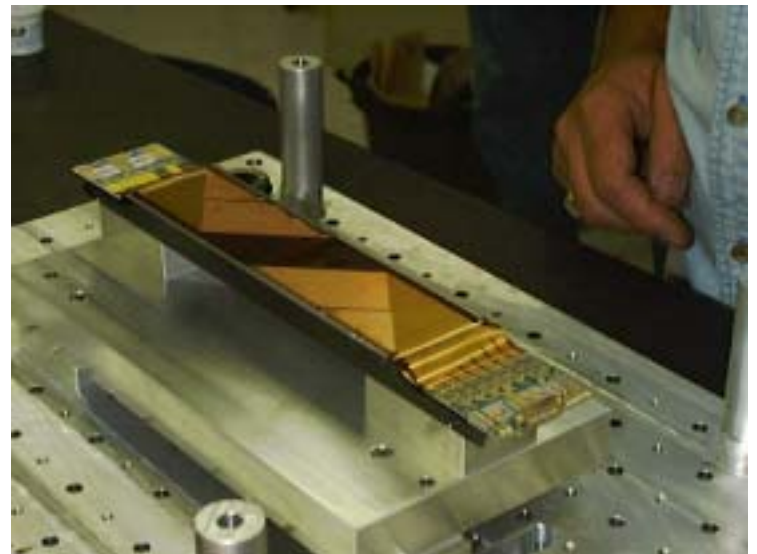
Detectors constructed of strips of width $10\ \mu\text{m}$ to give $O(10\ \mu\text{m})$ precision



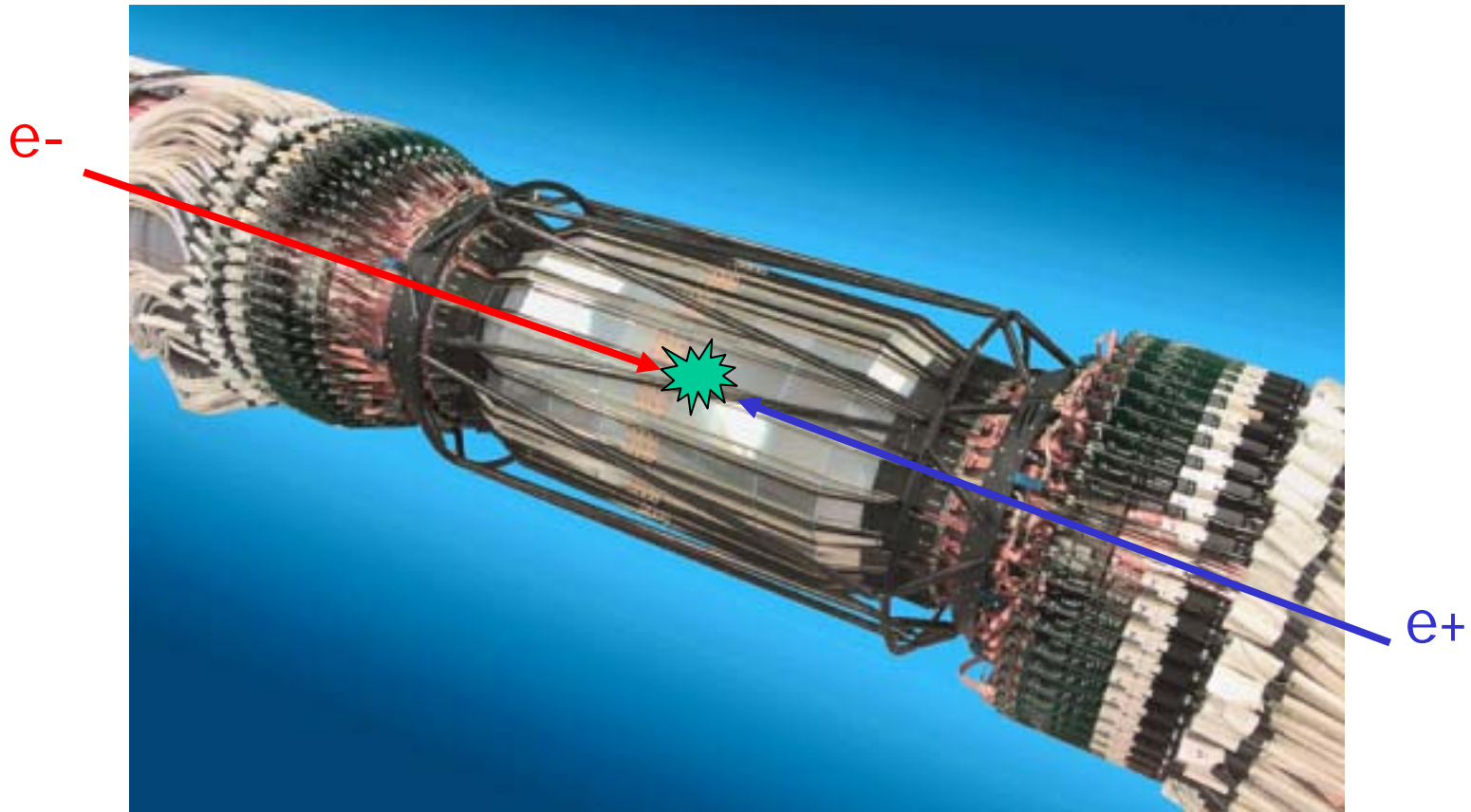
Construct detector of layers of thin strips.

Some of the electronics is incorporated in the silicon

Very light support structure (low multiple scattering)



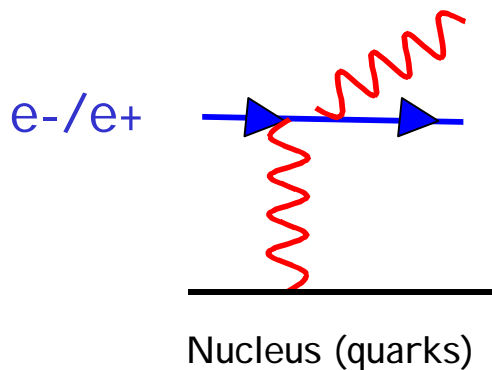
Solid State Tracking



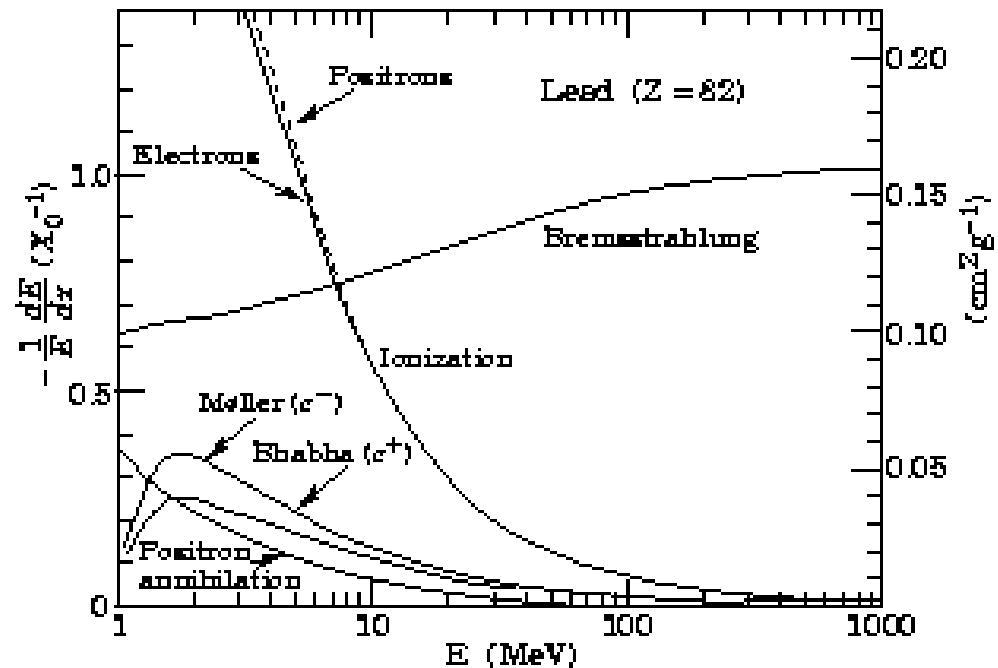
The BaBar Silicon Detector
5 layers, $r = 3.2$ to 14.4 cm

Measuring Particle Energy

Electron/Positron Energy Loss



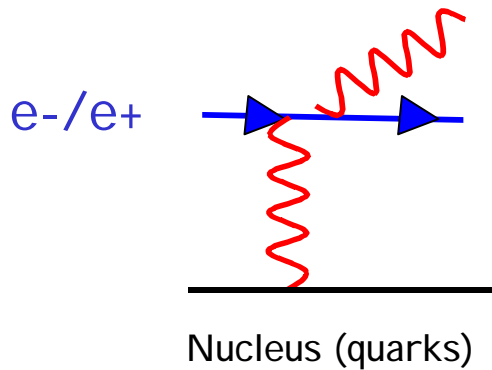
Bremsstrahlung
(radiation of photon)



Electron energy loss primarily by Bremsstrahlung at $E > 20$ MeV and ionization below 20 MeV

Radiation Length

Electron/Positron Energy Loss



Bremsstrahlung
(radiation of photon)

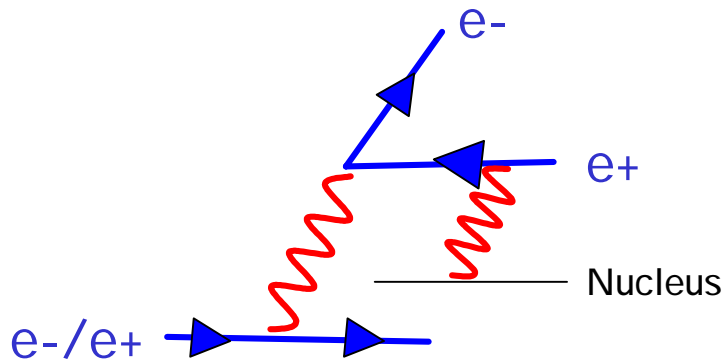
$$\frac{dE}{dx_{Brem}} = -\frac{E}{X_0}$$

$$X_0 = \frac{180A}{Z^2} g \text{ cm}^{-2}$$

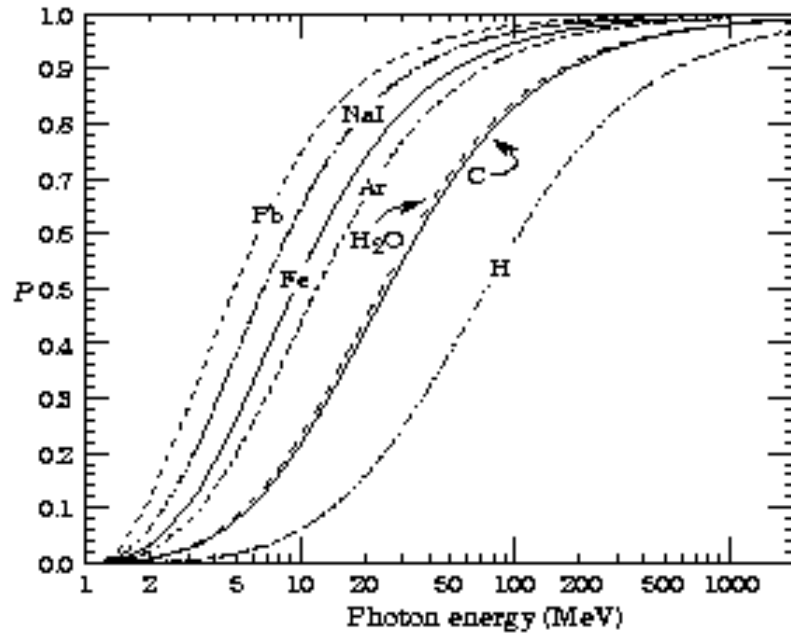
Radiation probability depends on radiation length X_0

Measuring Particle Energy

Photon Energy Loss

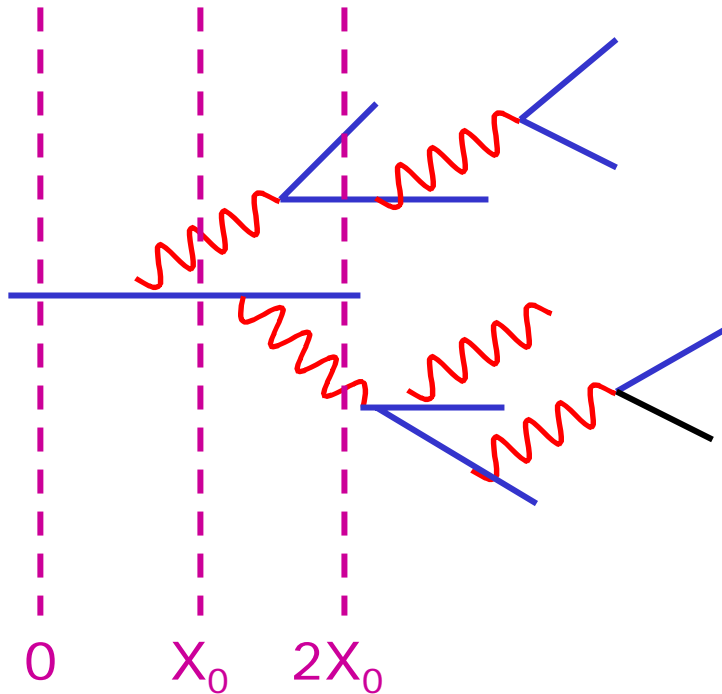


Pair Production



Photon energy loss primarily pair production at $E > 20$ MeV and Compton scattering ($\gamma e \rightarrow \gamma e$) below 20 MeV

Electromagnetic Showers



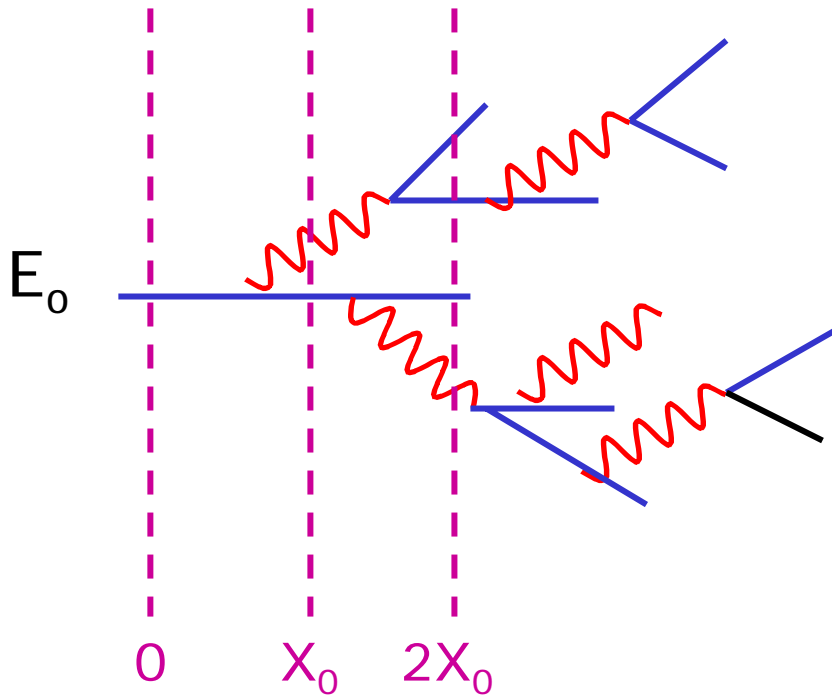
Assume

1. Each electron $E > E_c$ travels $1 X_0$ and gives up 50% E to photon
2. Each photon travels $1 X_0$ and pair produces with 50% E to each
3. Electrons with $E < E_c$ stop

Electrons/positrons or photons will form a cascade by the combination of Bremsstrahlung and pair production

Electromagnetic Showers

Number of particle after
t radiation lengths= $N(t)$

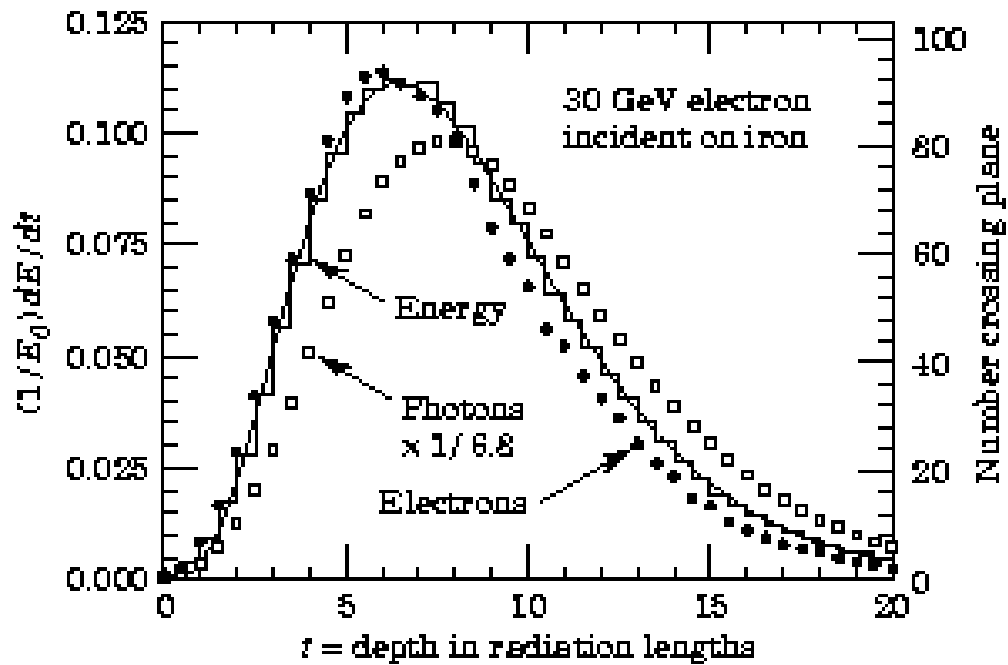


$$N(t) = 2^t = e^{t \ln 2}$$

The total track length of particles

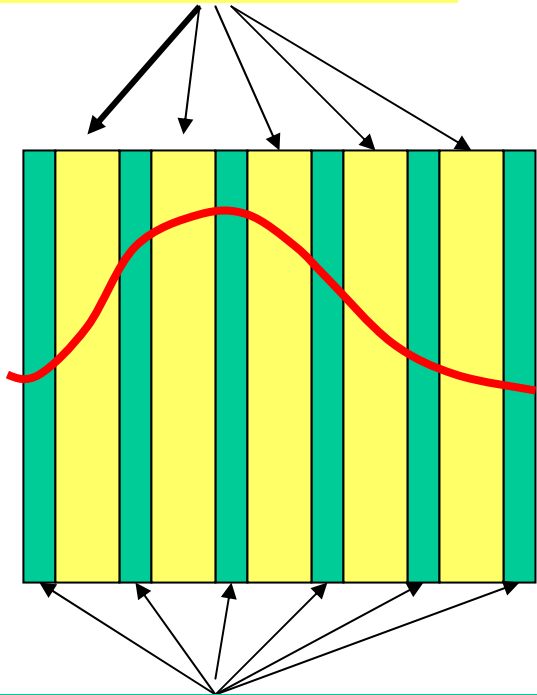
$$L = \frac{E_0}{E_c}$$

Shower Development



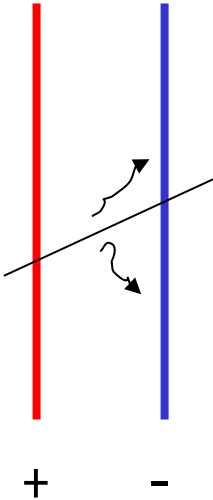
The Sampling Calorimeter

Lead- causes shower



Active Detector to measure total track length L

E.g. Ionization Chamber

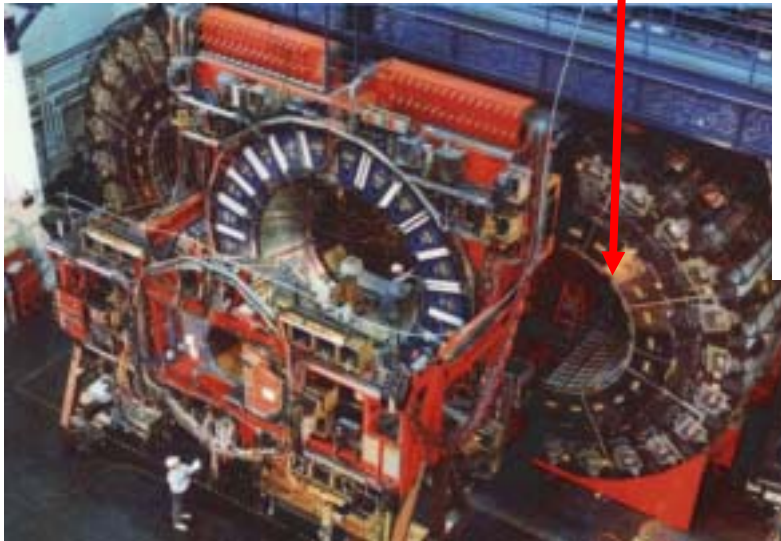
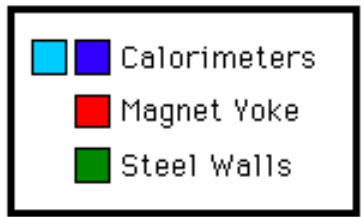
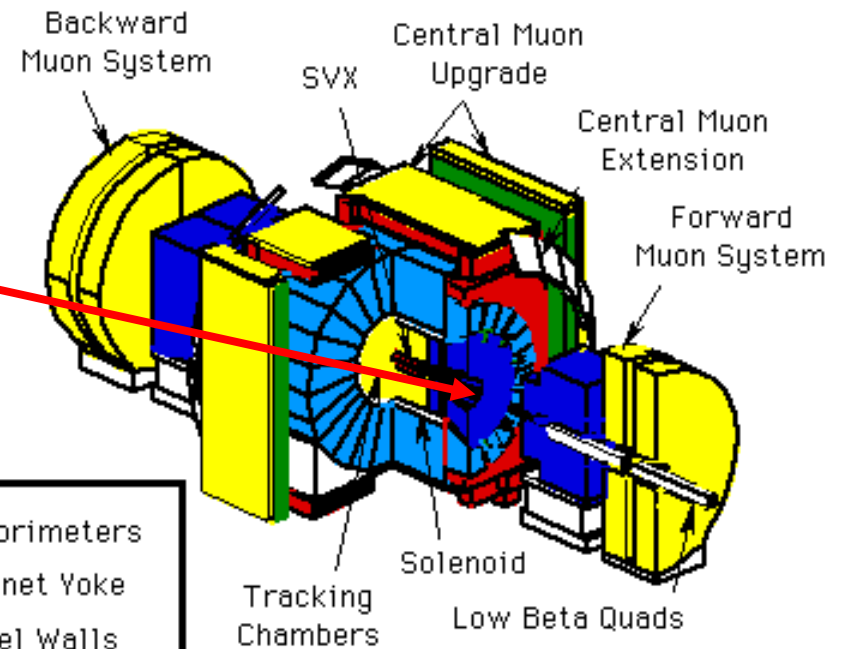


$$Q_{ionization} \propto L$$

Example of Sampling Calorimeter

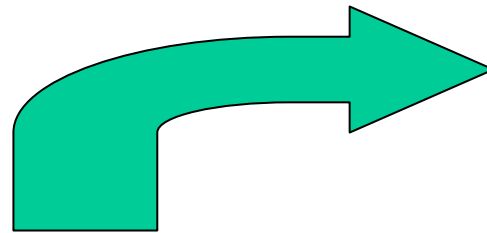
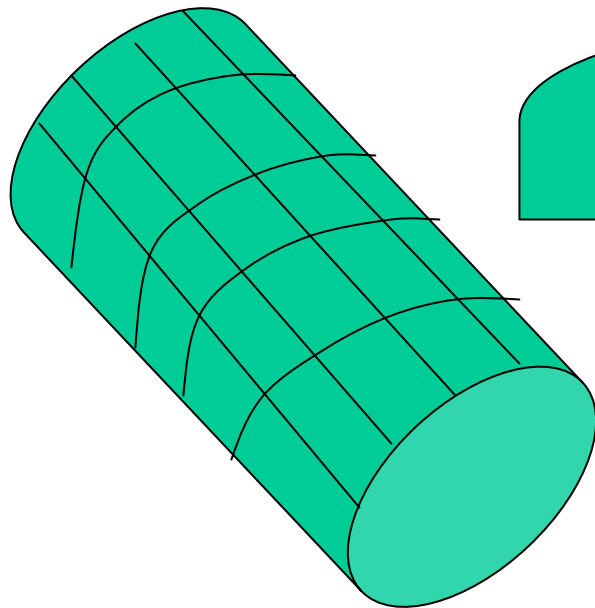
Sampling Calorimeter
(Lead/Scintillator)

CDF Detector



Sampling Calorimeters

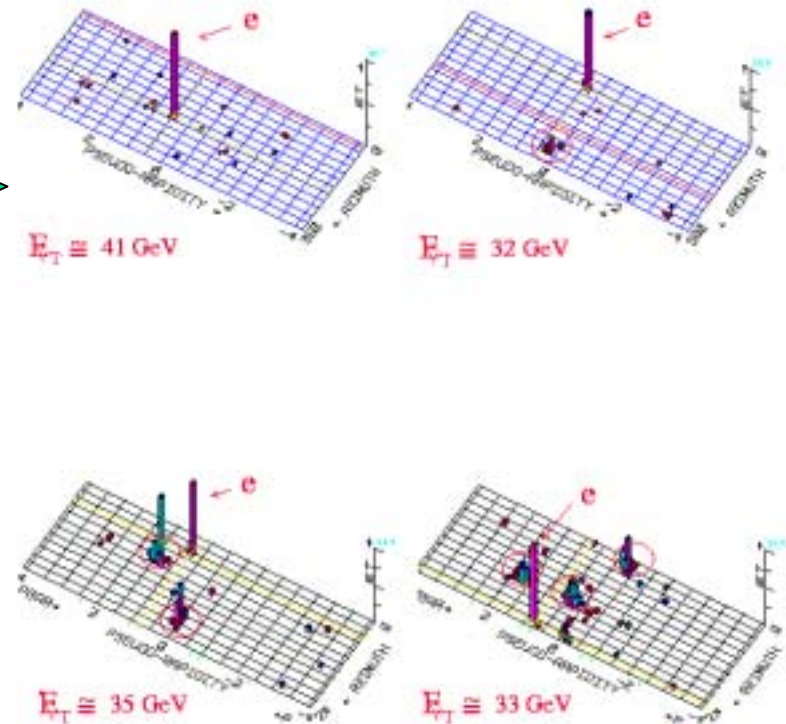
Calorimeters segmented into towers



Unwrap cylinder

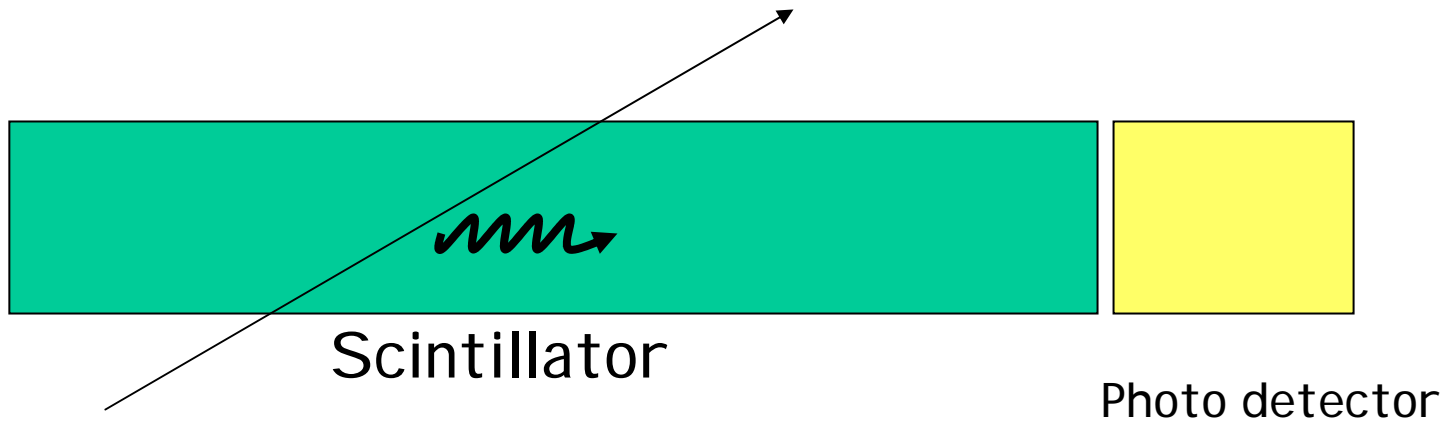
CDF:

W + 0,1,2,3 jet(s) Events



Scintillation

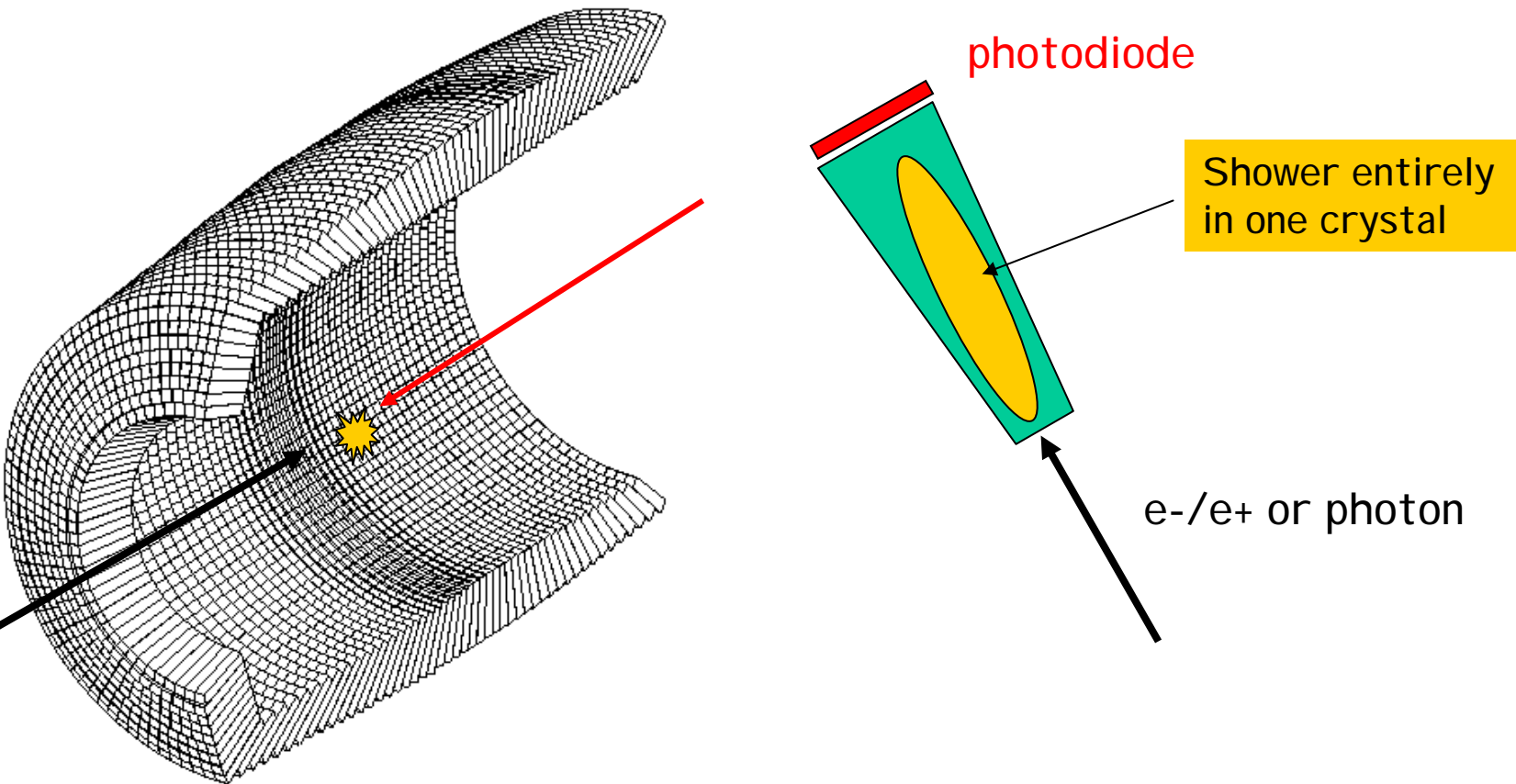
Particle causes excitation which results in light pulse



Number of photons proportional to track length

Crystal Calorimeters

Dense **scintillating crystals** like CsI (TI) can be used without showering medium - all the energy of particle is measured



The BaBar Calorimeter has 6580 crystals in a cylinder

Construction of the BaBar Calorimeter



Identifying Particles

Can distinguish photons and electrons

photon = electromagnetic shower without track

electron = electromagnetic shower with track

but

Protons, charged kaons, charged pions, muons

do not make a shower but do make a track.

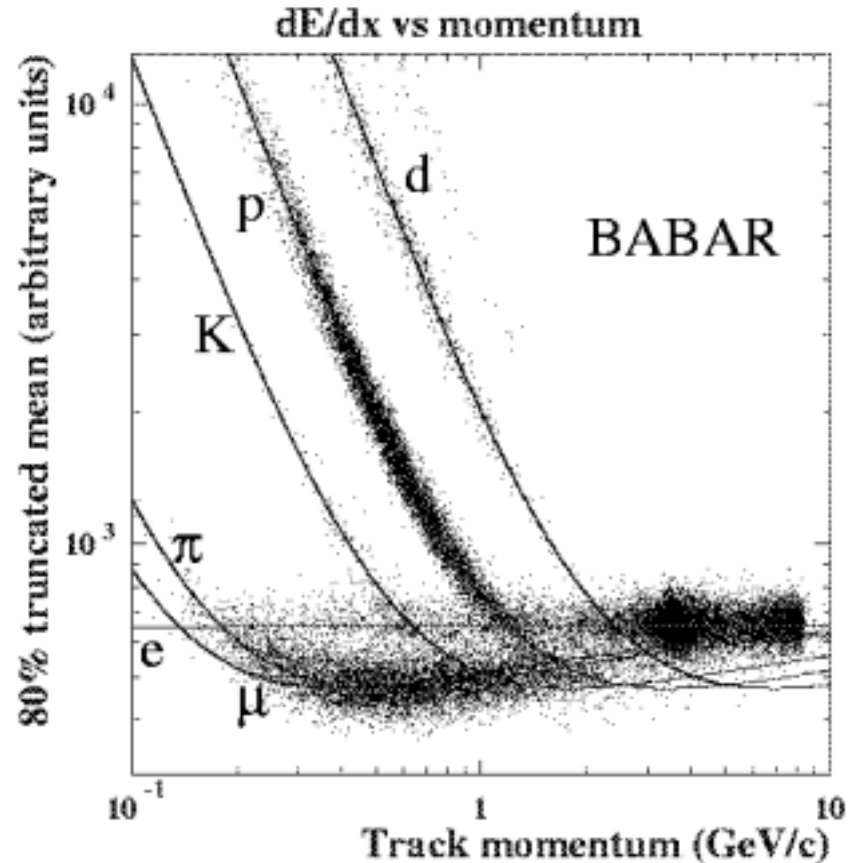
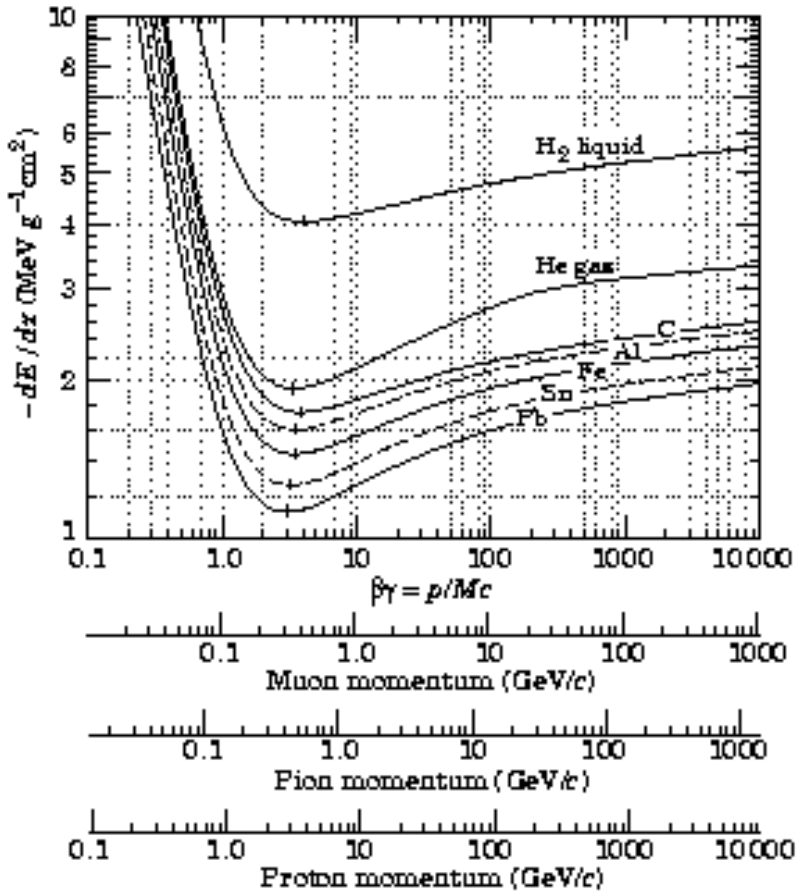
How to distinguish ?

Since $p = mv$ if we know the momentum then we need to measure the mass or velocity

1. dE/dx - less than 1 GeV, where dependence on β is steep
2. Cerenkov Detectors > 1 GeV

dE/dx using Drift Chamber

Recall dE/dx depends on velocity only



Measure dE/dx (velocity) and momentum to infer mass

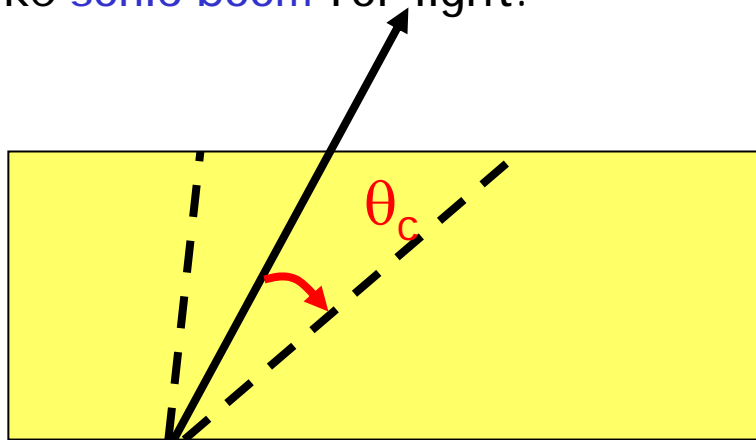
The Cerenkov Effect

Light in medium with refractive index n moves at velocity c/n .

If **particle velocity** $> c/n$

then Cerenkov light is emitted.

Like **sonic boom** for light.



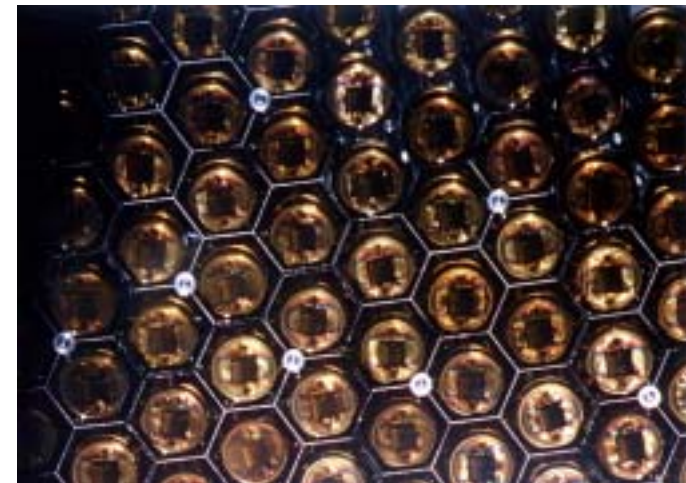
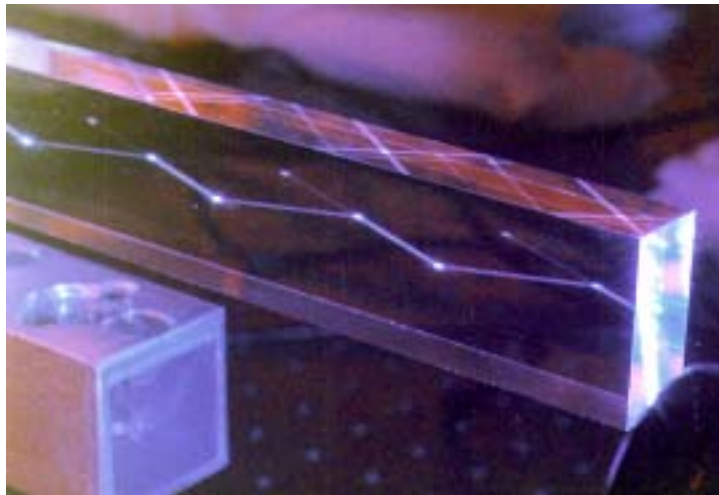
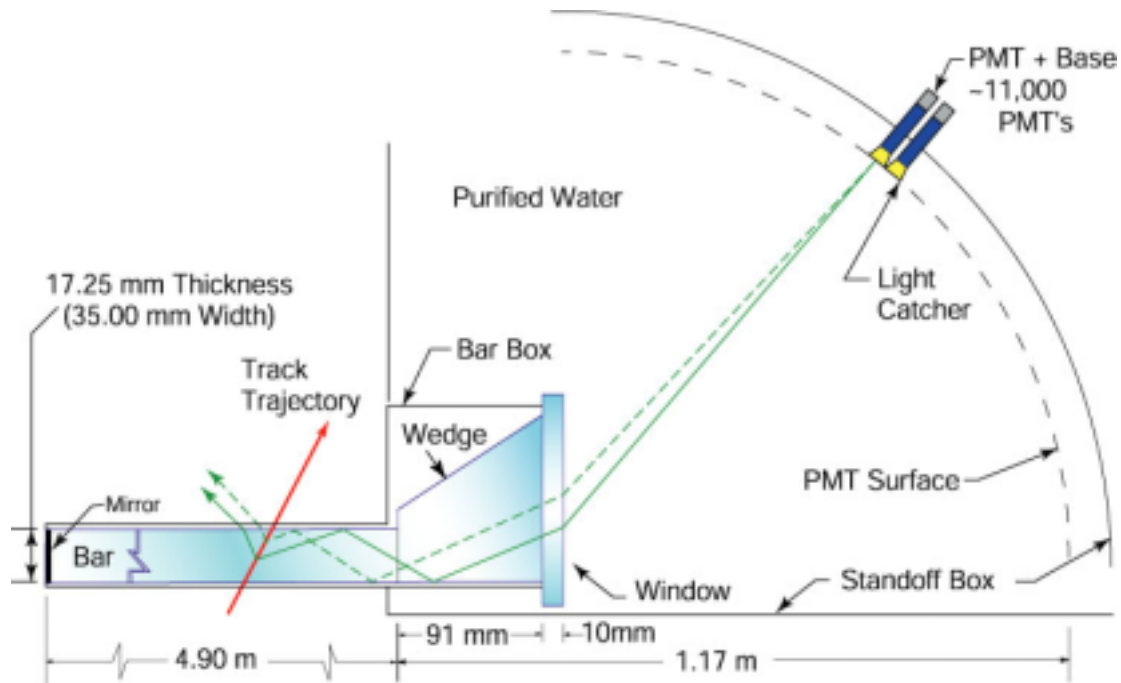
$$\theta_c = \cos^{-1}\left(\frac{1}{\beta n}\right)$$

Cone angle measures velocity

electrons in nuclear reactor pool emit Cerenkov light



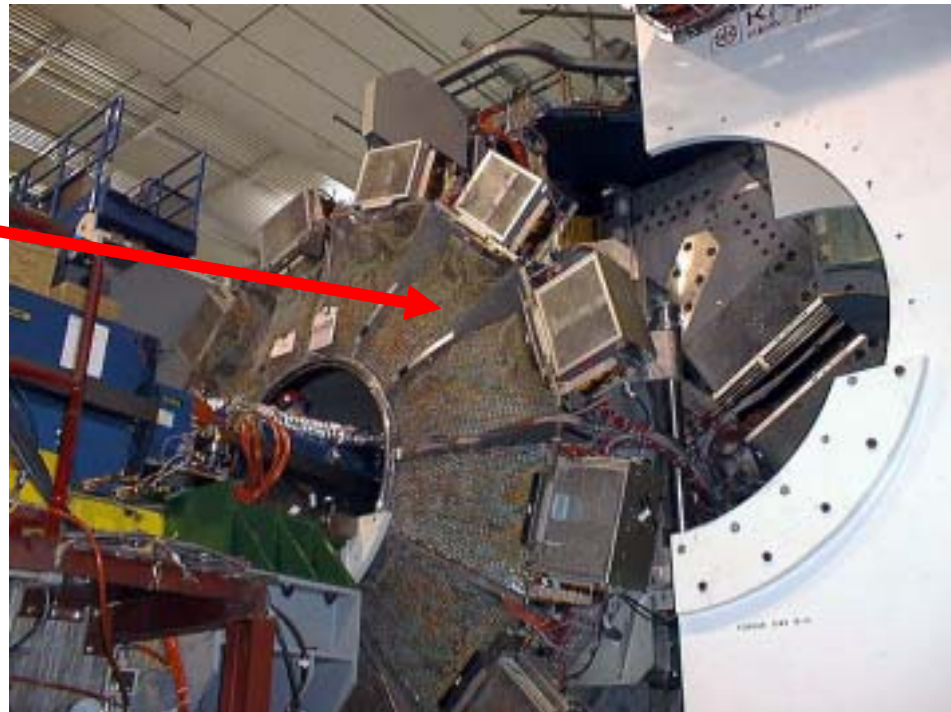
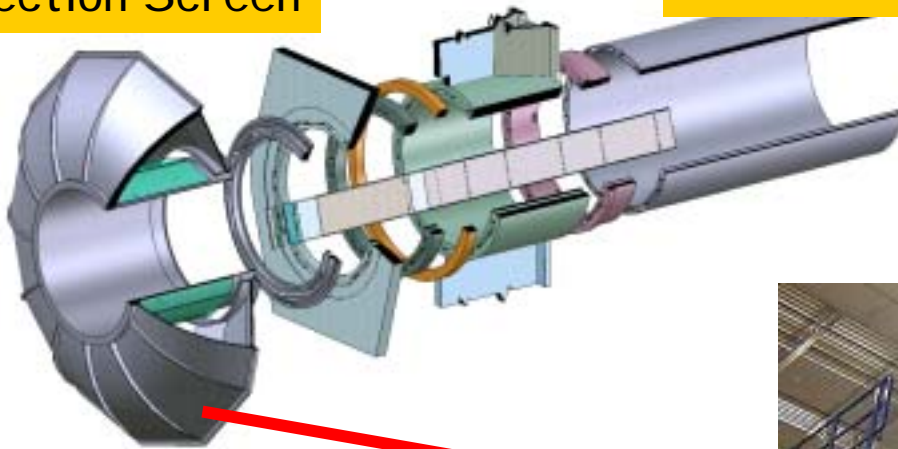
The BaBar Cerenkov Detector



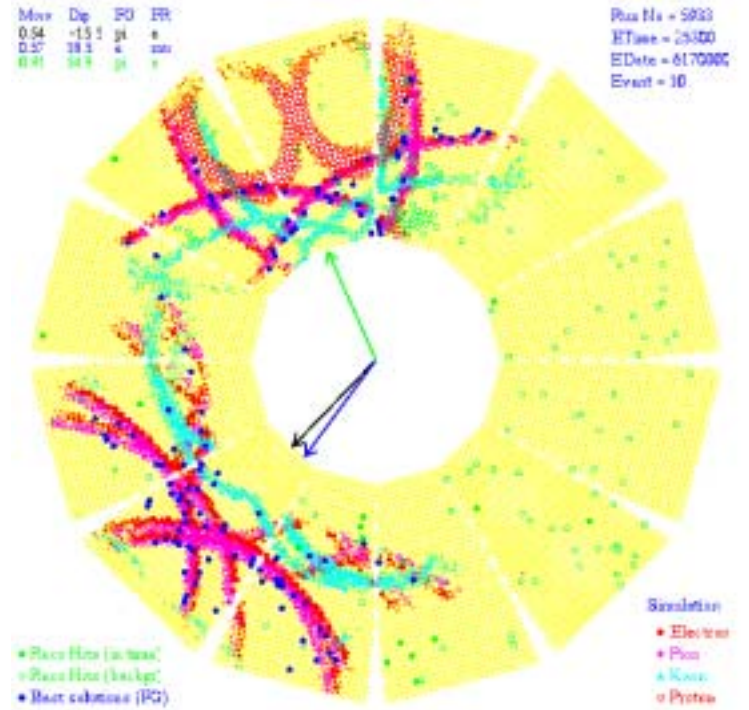
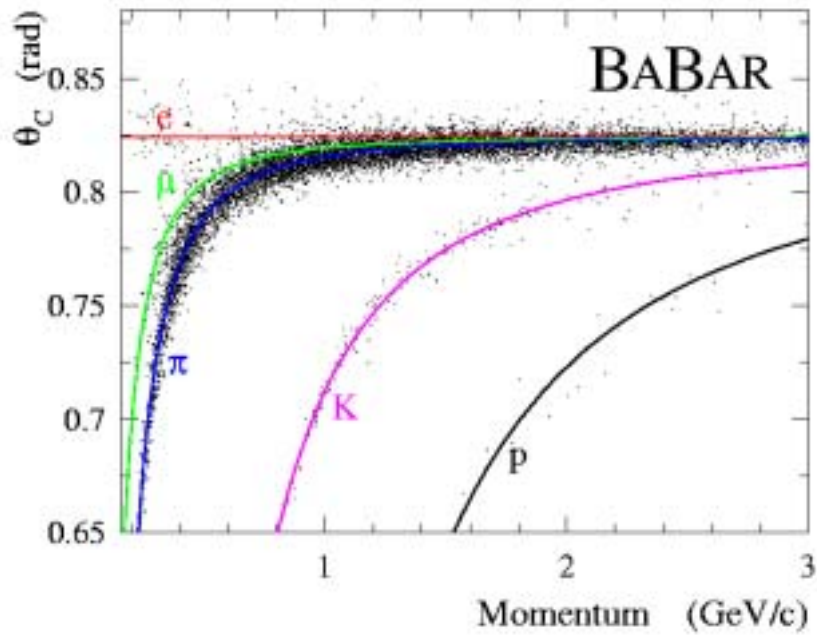
BaBar Cerenkov Detector Construction

Projection Screen

Quartz bars

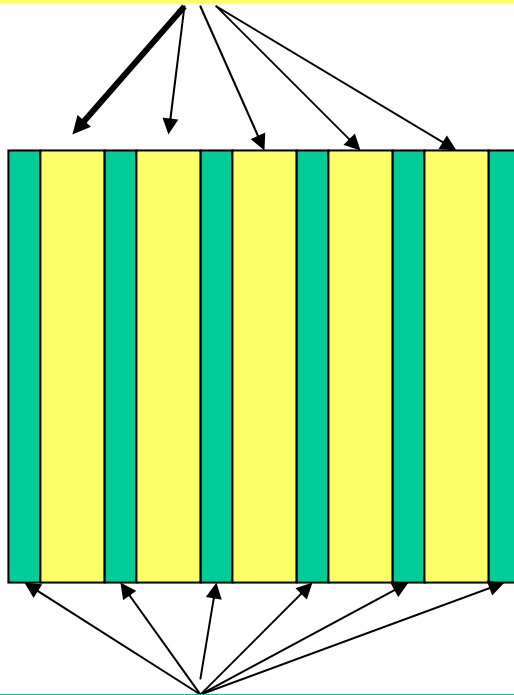


Event Showing Cerenkov Rings in BaBar



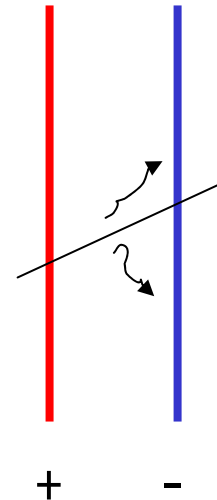
Muon I dentification

Iron shields other particles



Active Detector to measure total track length L

E.g Ionization Chamber



$$Q_{ionization} \propto L$$

Pions range out in iron due to strong interactions - only muons left

The BaBar Muon Detector

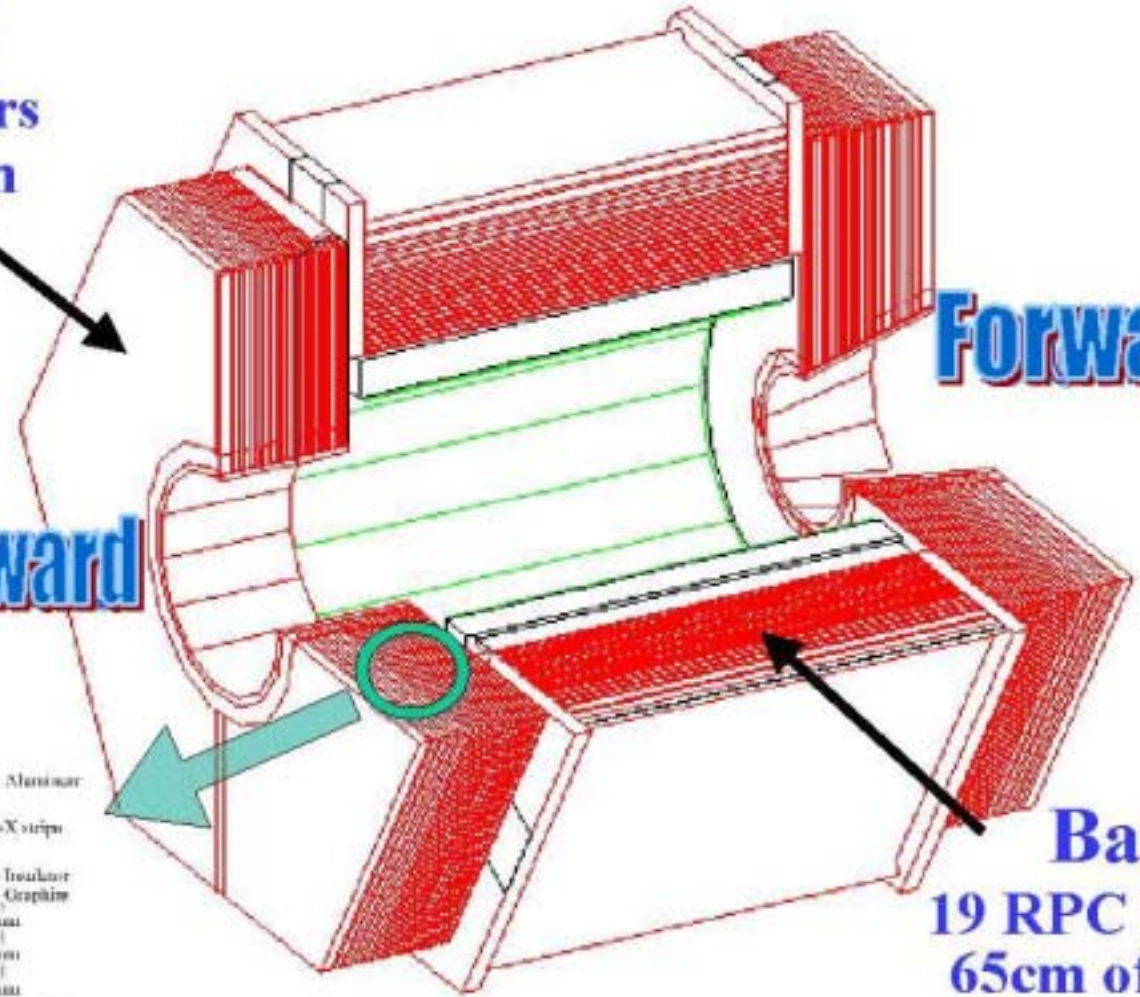
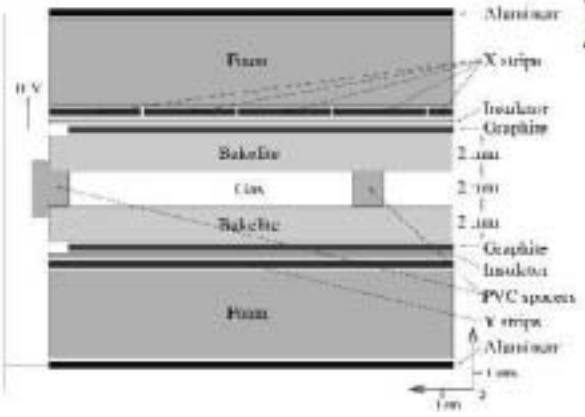
Endcap
18 RPC layers
60cm of iron

Forward

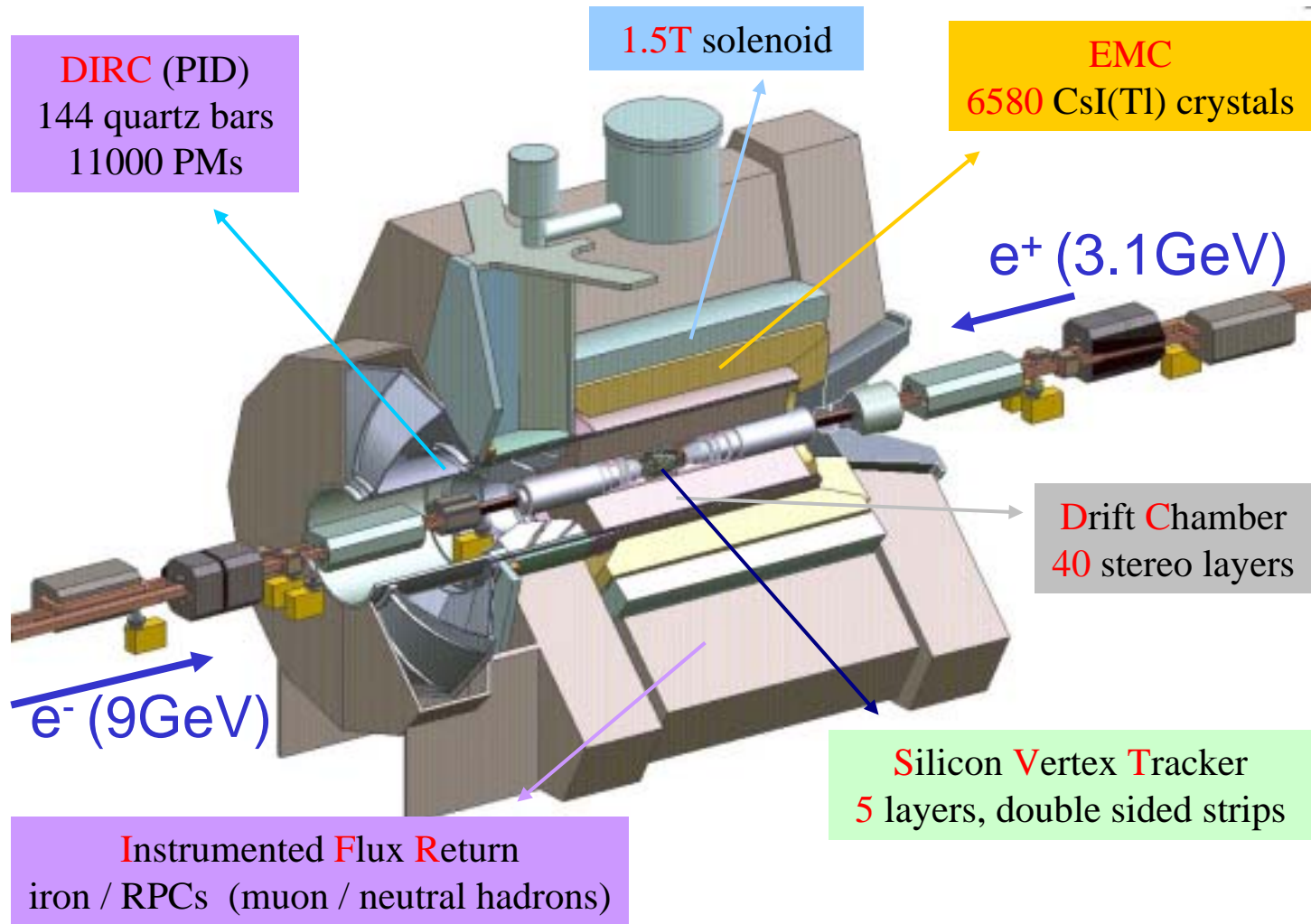
Backward

Barrel
19 RPC layers
65cm of iron

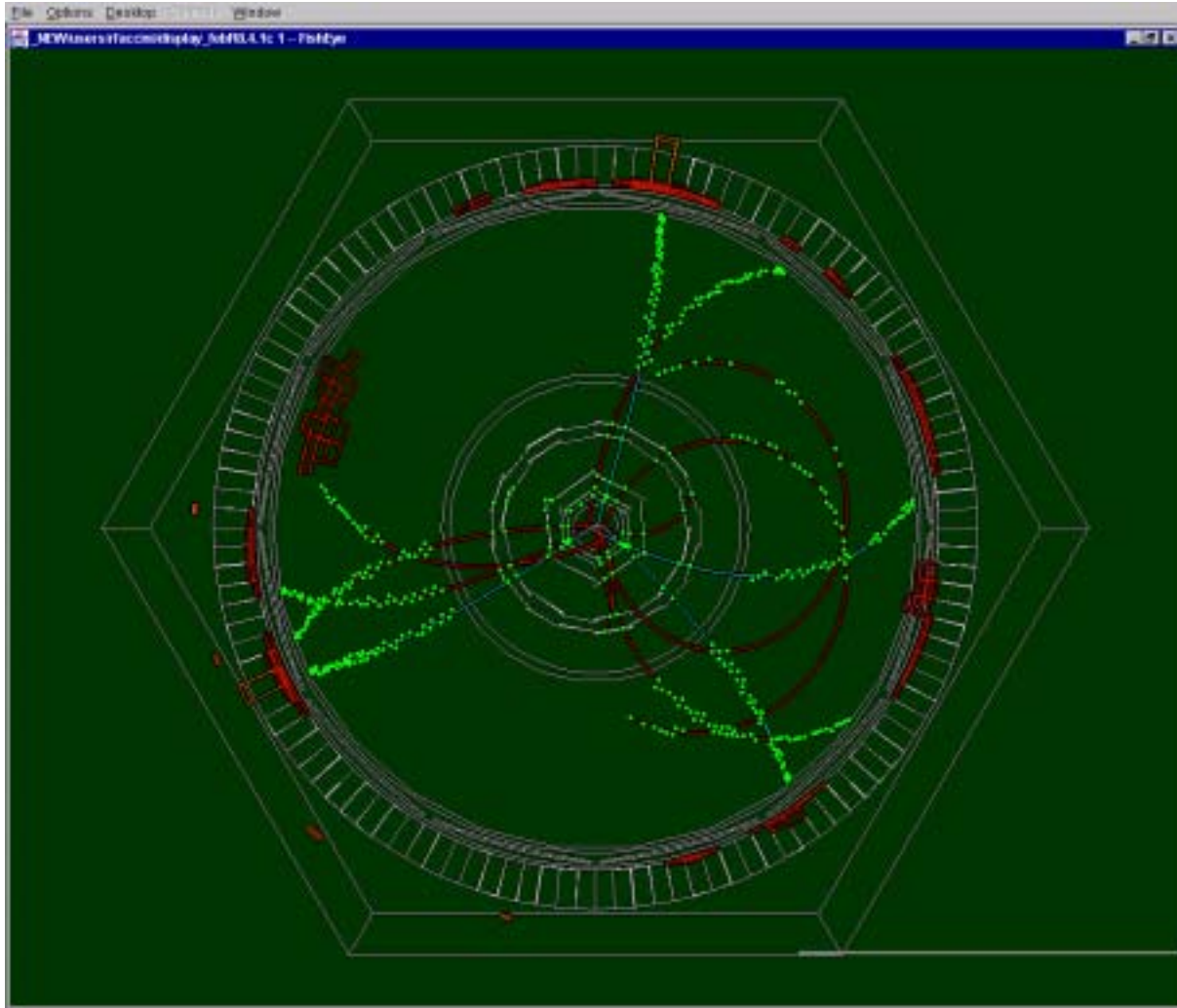
RPC section



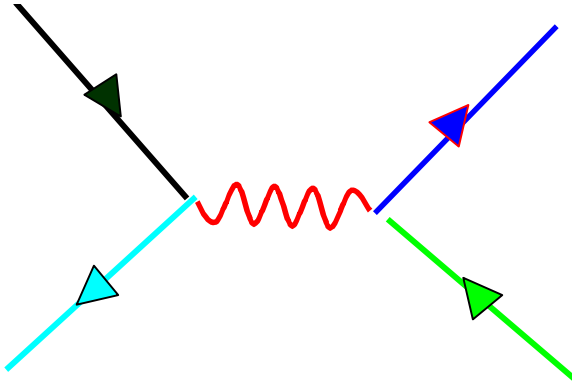
The Complete BaBar Detector



A matter-antimatter Event



$e^+e^- \rightarrow \mu^+\mu^-$ at PEP II

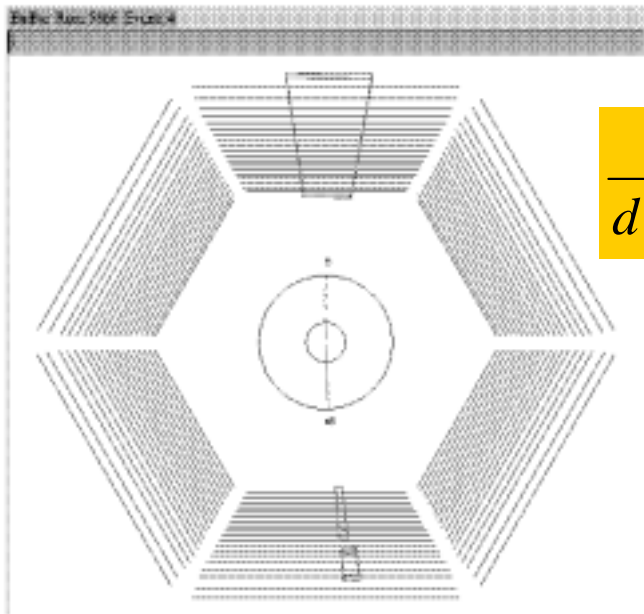


$$\frac{d\sigma}{d(\cos \vartheta)_{cm}} = \frac{\pi\alpha^2}{2s} (1 + \cos^2 \vartheta)$$

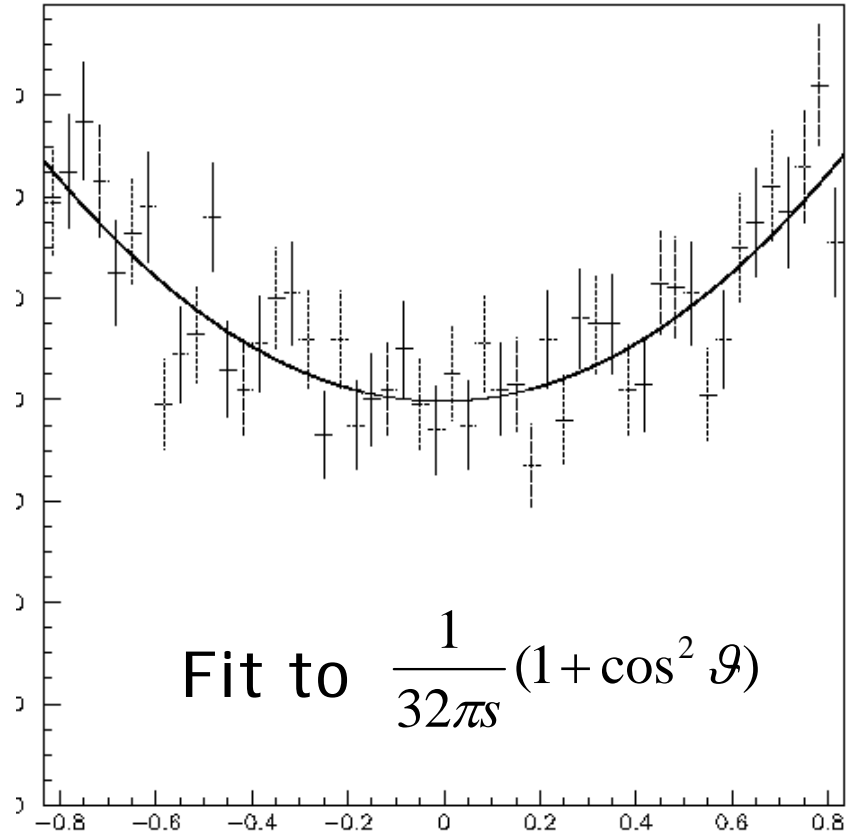
$$\sigma = \frac{4\pi\alpha^2}{3s}$$

At $s = (10.58 \text{ GeV})^2$ (PEP II)
 $\sigma(e^+e^- \rightarrow \mu^+\mu^-) = 0.78 \text{ nb}$

Muon Pair Production Data from BaBar



$$\frac{dE}{d(\cos \mathcal{G})_{cm}}$$



Fit to $\frac{1}{32\pi s} (1 + \cos^2 \mathcal{G})$

$$\cos \mathcal{G}_{cm}$$