The Energy-Energy Correlator At Small Angles

Lance Dixon (SLAC)

LD, Ian Moult, HuaXing Zhu, 1905.01310

Uppsala University

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The EEC

- Energy-energy correlation (EEC) in e^+e^- annihilation:

one of first infrared safe event-shapes defined in QCD, over 40 years ago Basham, Brown, Love, S. Ellis, PRD, PRL 1978

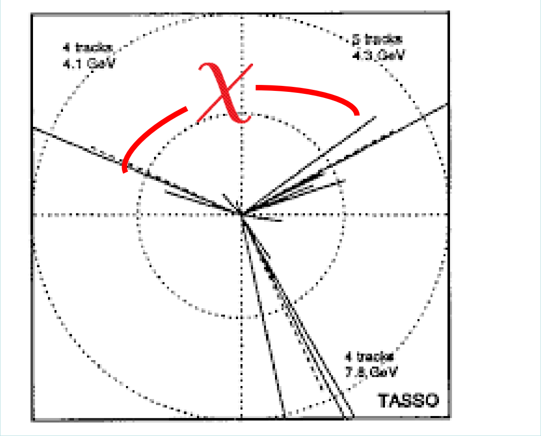
$$\frac{d\Sigma}{d\cos\chi} = \sum_{\text{partons } i,j} \int d\sigma \, \frac{E_i E_j}{Q^2} \delta(\cos\theta_{ij} - \cos\chi)$$

Collinear parton splitting

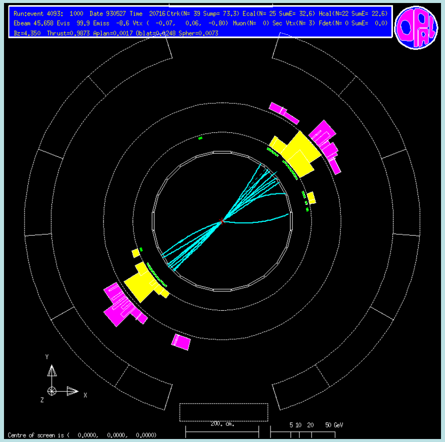
$$E_i \to x E_i + (1-x) E_i$$

preserves observable.
So does soft emission.

Data from wide range of CM energies →

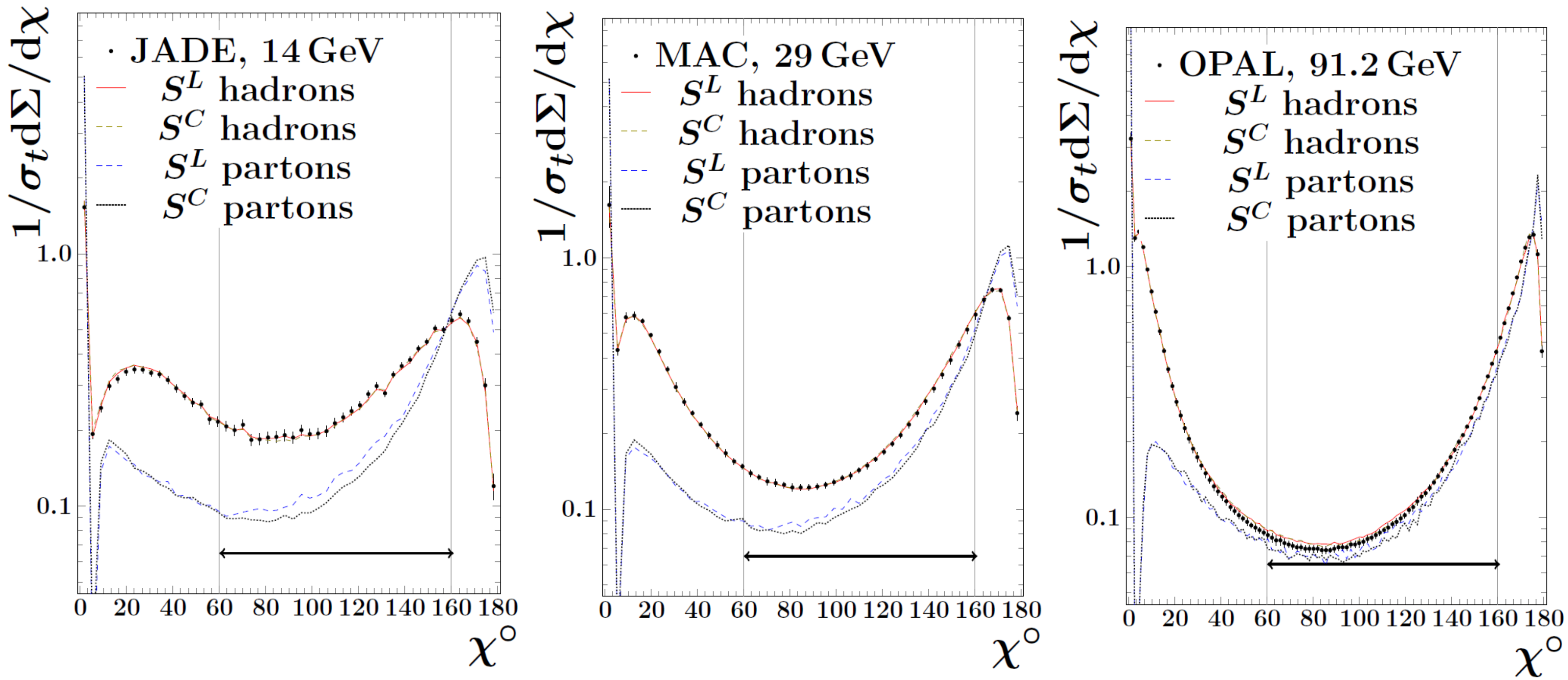


$Q \approx 22$ GeV



$Q = 91$ GeV

Evolution with energy clearly visible



data reviewed in Kardos et al, 1804.09146

Why the EEC?

- Many event-shape variables to choose from:
  thrust, oblateness, C parameter, heavy jet mass, angularity, jet rates, …
- EEC among the simplest analytically
- Angle χ lives on a compact domain, $[0,\pi]$:
  large logarithms on **both** ends can be resummed.
- Might get enough precision to measure α_s precisely
- As $\chi \to 0$, probe jet substructure. Can generalize to computable LHC jet substructure variables, correlating multiple small angles Moult, Necib, Thaler, 1609.07483
- **Gravitons** couple to **energy**, so AdS/CFT holography can be used to compute at strong gauge coupling (in planar N=4 SYM, not QCD) Hofman, Maldacena, 0803.1467

Why the EEC (cont.)?

- Energy-momentum tensor a fundamental object in any QFT, but also in conformal field theories.
- Alternative methods of study in CFT, especially N=4 SYM:
- Mellin representation of a four-point correlation function → analytic results in N=4 SYM at NLO, NNLO

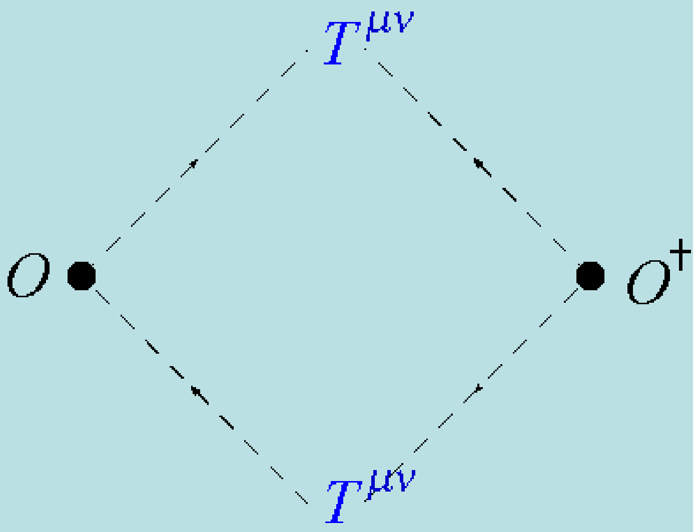
Belitsky, Hohenegger, Korchemsky, Sokatchev, Zhiboedov, 1309.0769, 1309.1424, 1311.6800, 1409.2502;

Henn, Sokatchev, Yan, Zhiboedov, 0903.05314; Korchemsky, 1905.01444

- Using properties of "ANEC"

light-ray operators

Kologlu, Kravchuk, Simmons-Duffin, Zhiboedov, 1905.01311

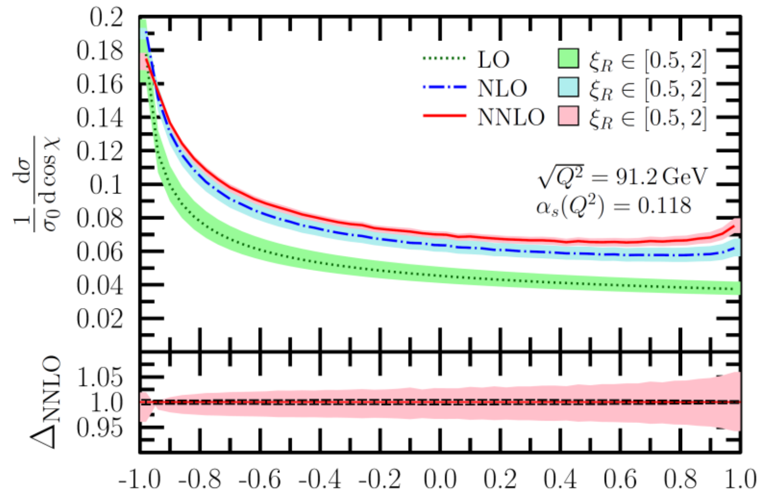


Numerical results in QCD

- EEC computed at NLO numerically in 1980s and 1990s
  Richards, WJ Stirling, Ellis, 1982, 1983; Ali, Barreiro, 1982, 1984; Schneider, Kramer, Schierholz, 1984; Falck, Kramer, 1989; Kunszt, Nason, Marchesini, Webber, LEP Yellow Book, 1989; Glover, Sutton, 1994; Clay, Ellis, 1995; Kramer, Spiesberger, 1996; Catani, Seymour, 1996 [EVENT2].
- Computed numerically at NNLO only 3 years ago

  Del Duca, Duhr, Kardos, Somogyi, Trocsanyi, 1603.08927

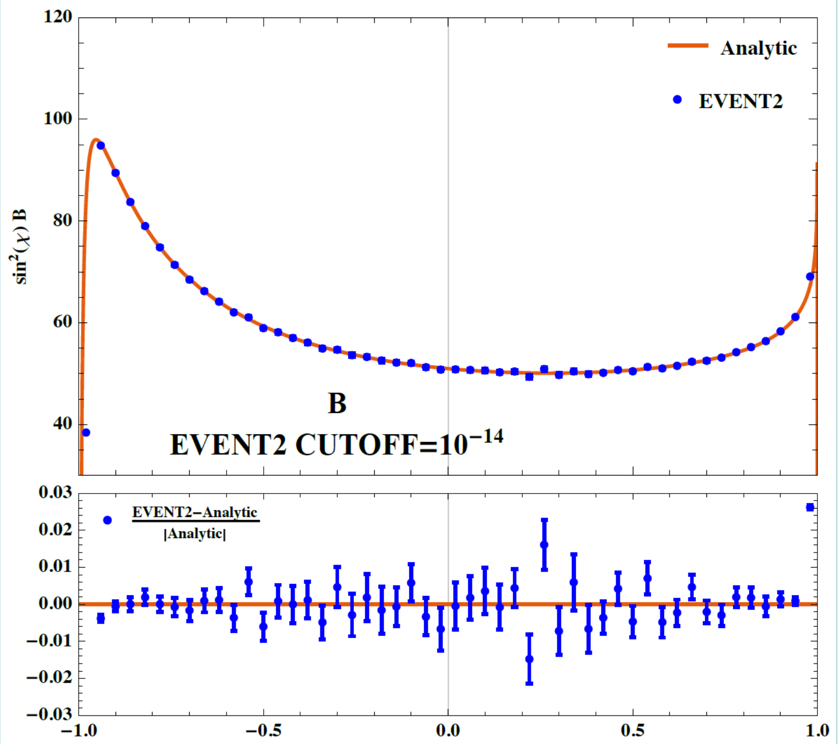
- Can now compute analytically at NLO in QCD
  LD, Luo, Shtabovenko, Yang, Zhu, 1801.03129



Why analytic?

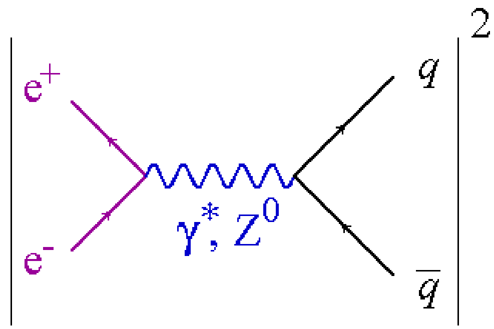
- Validate accuracy of numerical QCD results.
- Compare with analytic NLO result in N=4 SYM

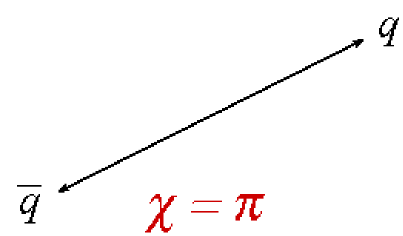
Belitsky, Hohenegger, Korchemsky, Sokatchev, Zhiboedov, 1309.0769, 1309.1424, 1311.6800



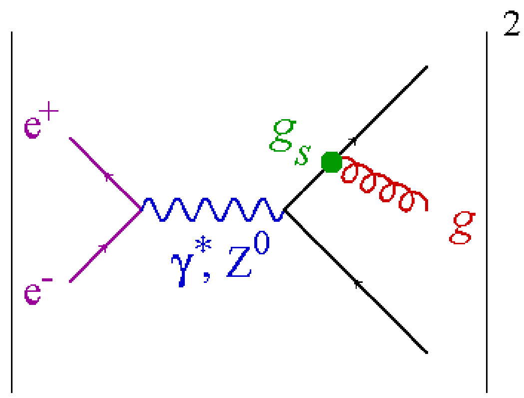
- Study limits as $\chi \to 0,\pi$ to aid resummation of large logarithms there.

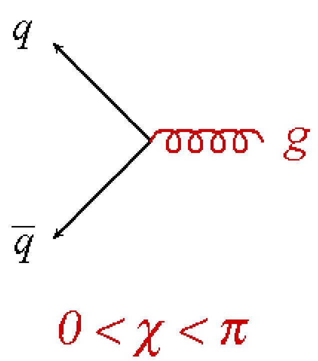
LO EEC for $0 < \chi < \pi$ is $O(\alpha_s)$





$$\alpha_s = \frac{g_s^2}{4\pi}$$

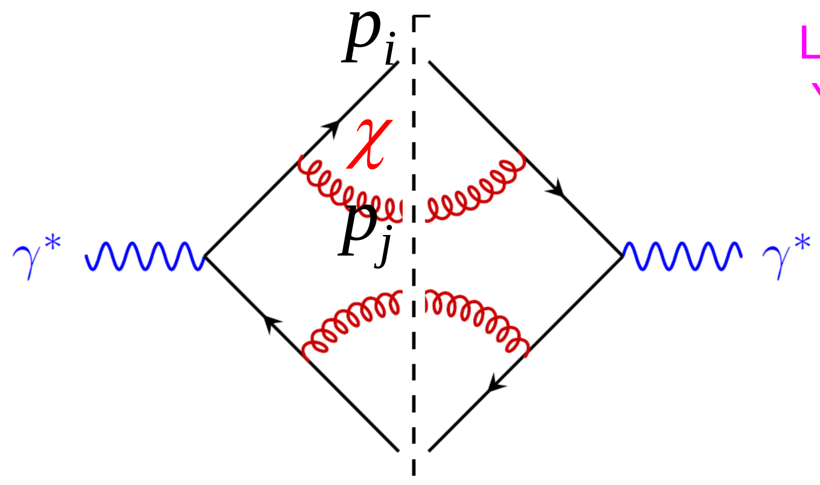




How to compute at NLO in QCD?

LD, Luo, Shtabovenko, Yang, Zhu, 1801.03129

Sample NLO real emission contribution



- Interference method with Feynman diagrams
- Reverse unitarity: Treat all momenta as loop momenta, put all cut momenta on shell and impose $\delta(\cos\theta_{ij} - \cos\chi)$
- IBPs/Laporta algorithm Chetyrkin, Tkachov (1981), Laporta (2001)
- Differential equations for master integrals Gehrmann, Remiddi (2000)

can all be solved in terms of polylogarithms.

Same method works also for Higgs → gg → hadrons

Luo, Shtabovenko, Yang, Zhu, 1903.07277

Structure of QCD result

$$\frac{1}{\sigma_0}\frac{d\Sigma}{d\cos\chi} = \frac{\alpha_s(\mu)}{2\pi}A(z) + \left(\frac{\alpha_s(\mu)}{2\pi}\right)^2 \left(\beta_0 \log\frac{\mu}{Q}A(z) + B(z)\right) + \mathcal{O}(\alpha_s^3)$$

$$z = \frac{1}{2}(1-\cos\chi) \;\in\; [0,1]$$

LO result fits on one line:

Basham, Brown, Love, S. Ellis, 1978

$$A(z) = C_F \frac{3-2z}{4(1-z)z^5}[3z(2-3z) + 2(2z^2 - 6z + 3)\ln(1-z)]$$

NLO result will be expressed in terms of classical polylogarithms:

$$\mathrm{Li}_n(u) = \int_0^u \frac{dt}{t}\mathrm{Li}_{n-1}(t), \qquad \mathrm{Li}_1(t) = -\ln(1-t)$$

Belitsky et al. method for N=4 SYM

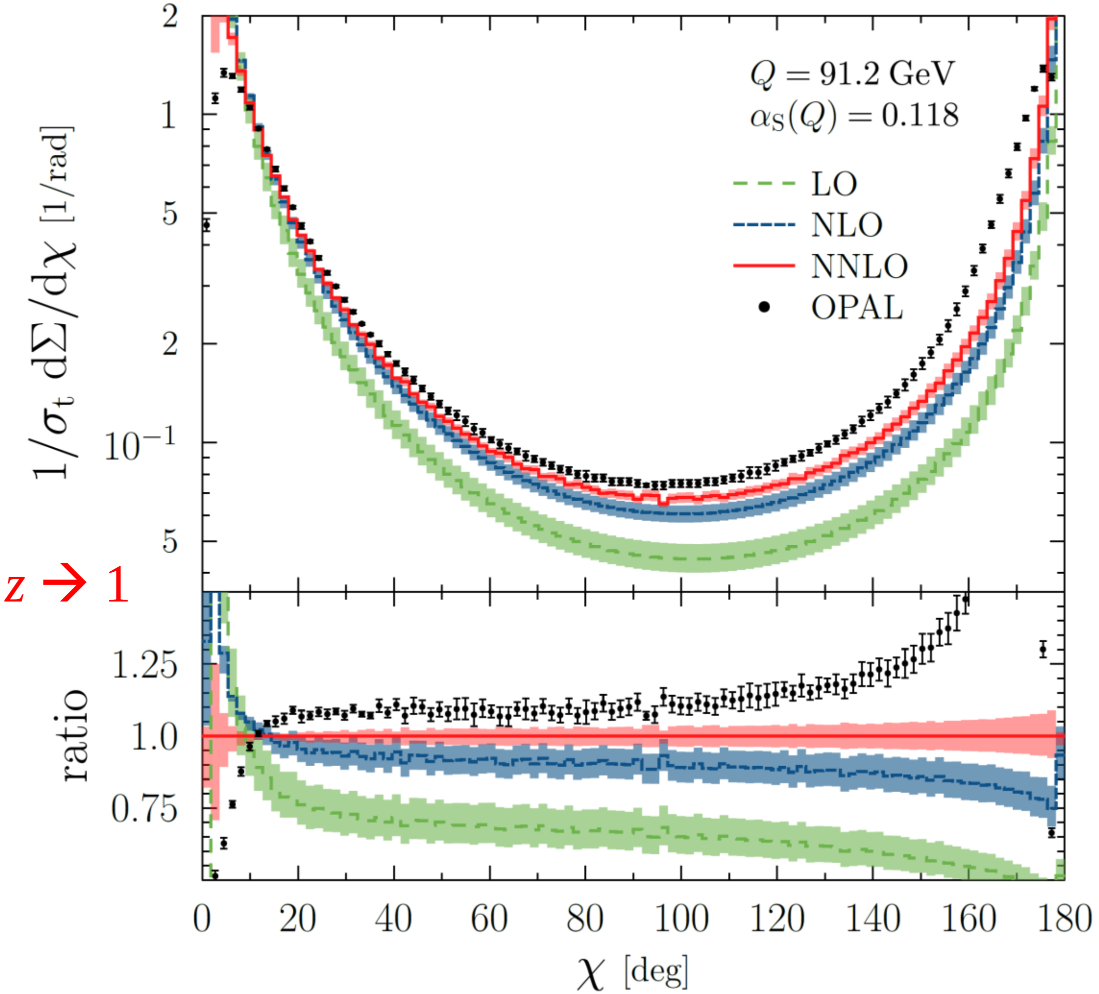
- Very different from "QCD method", which uses dimensional regularization; divergences cancel between virtual and real
- Exploit conformal invariance of 4-point function with two "energy flow operators"

$$\langle \mathcal{E}(\vec{n}_1)\mathcal{E}(\vec{n}_2)\rangle_q = \int d^4x\, e^{iq\cdot x} \langle 0|O^\dagger(x)\mathcal{E}(\vec{n}_1)\mathcal{E}(\vec{n}_2)O(0)|0\rangle$$

$$\mathcal{E}(\vec{n}) = \int_{-\infty}^{\infty} d\tau \lim_{r\to\infty} r^2 n^i T_{0i}(t=\tau+r, r\vec{n})$$

- Analytically continue from Euclidean to physical region using double Mellin transform
- No infrared divergences at any step!
- Recently pushed to NNLO (semi-analytic): Henn, Sokatchev, Yan and Zhiboedov, 1903.05314

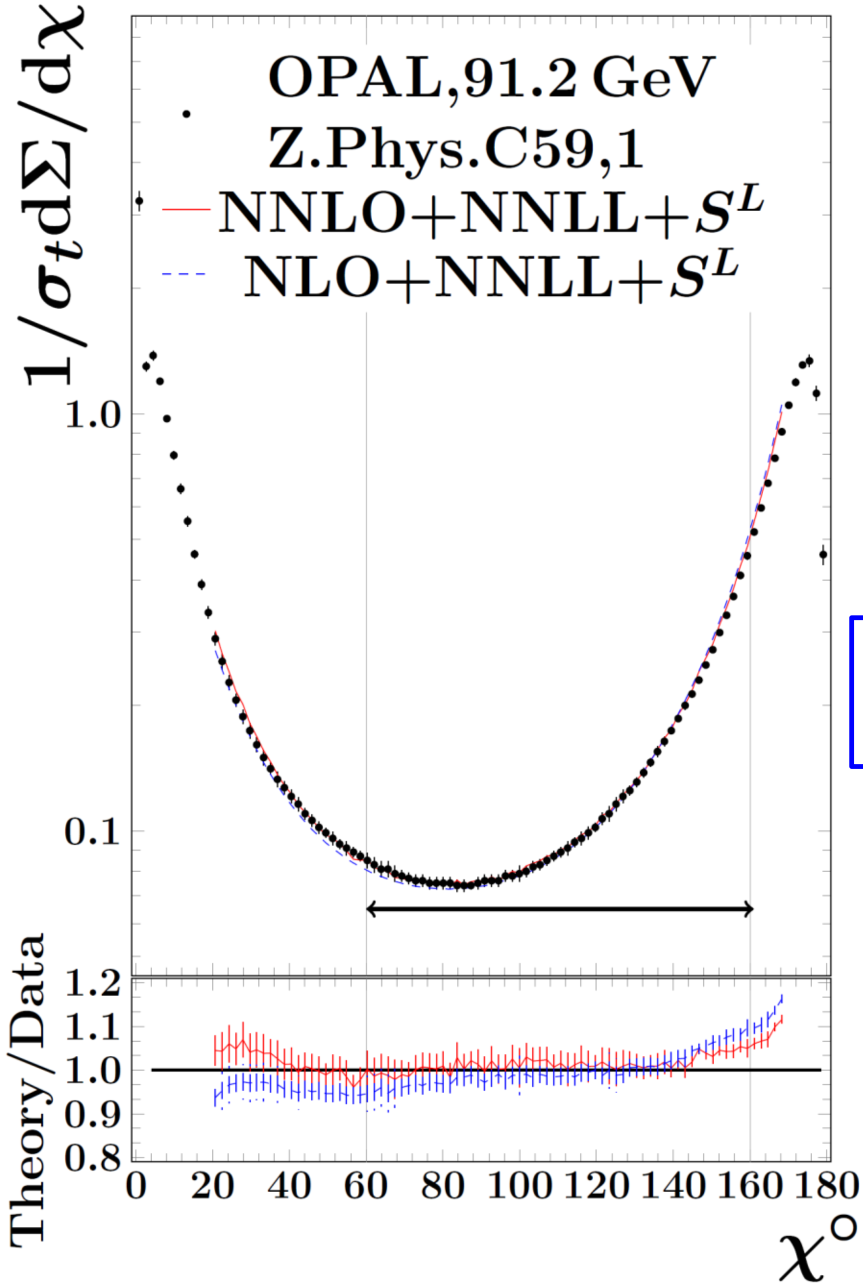
Fixed order QCD vs. Z pole data



z → 1

z → 0

Tulipant, Kardos, Somogyi, 1708.04093



To measure strong coupling α_s :
Add NNLL $z \rightarrow 1$ resummation
+ MC estimate of
nonperturbative contributions

Kardos, Kluth, Somogyi, Tulipant, Verbytskyi, 1804.09146

$$\alpha_s(M_Z) = 0.11750 \pm 0.00018(exp.) \pm 0.00102(hadr.) \pm 0.00257(ren.) \pm 0.00078(res.)$$

Competitive measurement of α_s

Still room for theory improvement:
→ NNNLO (approx.?)
+ NNNLL $z \rightarrow 1$ resummation (soon)
+ NNLL $z \rightarrow 0$ resummation

Back-to-back limit, $z \to 1$

$$
\begin{aligned}
B(z) = C_F \Bigg\{ & \frac{1}{1-z} \Bigg[\frac{1}{2} C_F \ln^3(1-z) + \ln^2(1-z) \left(\frac{11C_A}{12} + \frac{9C_F}{4} - \frac{N_f T_f}{3} \right) \\
& + \ln(1-z) \left(C_A \left(\frac{\zeta_2}{2} - \frac{35}{72} \right) + C_F \left(\zeta_2 + \frac{17}{4} \right) + \frac{N_f T_f}{18} \right) \\
& + C_A \left(\frac{11\zeta_2}{4} + \frac{3\zeta_3}{2} - \frac{35}{16} \right) + C_F \left(3\zeta_2 - \zeta_3 + \frac{45}{16} \right) + N_f T_f \left(\frac{3}{4} - \zeta_2 \right) \Bigg] \\
& + \left(\frac{C_A}{2} + C_F \right) \ln^3(1-z) + \ln^2(1-z) \left(\frac{27C_A}{8} + \frac{13C_F}{2} - \frac{N_f T_f}{2} \right) \\
& + \ln(1-z) \left[C_A \left(22\zeta_2 - \frac{2011}{72} \right) + C_F(47 - 19\zeta_2) + N_f T_f \left(\frac{361}{36} - 4\zeta_2 \right) \right] \\
& + C_A \left(\frac{6347\zeta_2}{80} - 21\zeta_2 \ln 2 - \frac{137\zeta_3}{4} - \frac{3305}{72} \right) \\
& + C_F \left(-\frac{1727\zeta_2}{20} + 42\zeta_2 \ln 2 + \frac{121\zeta_3}{2} + \frac{3437}{96} \right) \\
& + N_f T_f \left(-\frac{1747\zeta_2}{120} + 12\zeta_3 + \frac{2099}{144} \right) \Bigg\} + \mathcal{O}(1-z)
\end{aligned}
$$

leading power (first three lines)

first subleading power (remaining lines)

- Double log behavior, $\ln^{2L+1}(1-z)/(1-z)$ characteristic of Sudakov suppression from soft/collinear gluon emission. Collins, Soper,…
- Coefficients of leading-power terms agree precisely with NNLL resummation DeFlorian, Grazzini,hep-ph/0407241

Intra-jet limit, $z \to 0$

LL

$$B(z) = C_F \Big\{ \frac{1}{z} \Big[\ln z \Big(-\frac{107 C_A}{120} + \frac{25 C_F}{32} + \frac{53 N_f T_f}{240} \Big) + C_A \Big(-\frac{25\zeta_2}{12} + \frac{\zeta_3}{2} + \frac{17683}{2700} \Big) + C_F \Big(\frac{43\zeta_2}{12} - \zeta_3 - \frac{8263}{1728} \Big) - \frac{4913 N_f T_f}{3600} \Big]$$

$$+ \ln z \Big[C_A \Big(\frac{33\zeta_2}{2} - \frac{703439}{25200} \Big) + C_F \Big(\frac{42109}{1200} - 21\zeta_2 \Big) + N_f T_f \Big(\frac{86501}{12600} - 4\zeta_2 \Big) \Big]$$

$$+ C_A \Big(\frac{213\zeta_2}{5} - \frac{101\zeta_3}{2} - \frac{26986007}{5292000} \Big) + C_F \Big(-\frac{1541\zeta_2}{30} + 65\zeta_3 + \frac{18563}{2700} \Big)$$

$$+ N_f T_f \Big(-\frac{46\zeta_2}{3} + 12\zeta_3 + \frac{2987627}{330750} \Big) \Big\} + \mathcal{O}(z)$$

NLL

leading power

first subleading power

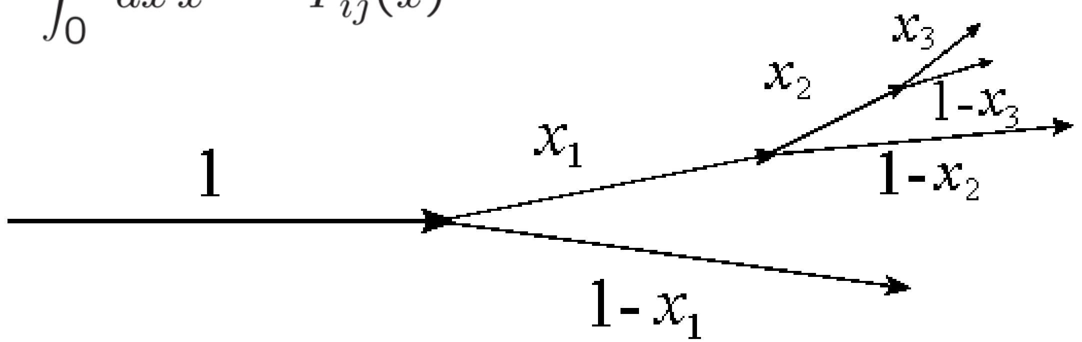
- Single log behavior, $\ln^L z / z$ characteristic of pure collinear observable.
- Leading log (LL) resummation first performed in "jet calculus" approach Konishi, Ukawa, Veneziano, Phys.Lett.1978,1979
- Coefficients of leading-power terms agree precisely with LL result Richards, Stirling, Ellis, NPB229, 317, 1983

LL resummation

- Limit dominated by collinear emission. At leading log, only a single Mellin moment N=3 of time-like splitting function (twist 2 anomalous dimension) dominates

Konishi, Ukawa, Veneziano; Richards, Stirling, Ellis

$$\gamma_{ij}^{(N)} \equiv -\int_0^1 dx\, x^{N-1}\, P_{ij}(x)$$



Energy weighting →
$$\int_0^1 dx\, x(1-x)\, P_{ij}(x) \;\to\; -\int_0^1 dx\, x^2\, P_{ij}(x) \;\equiv\; \gamma_{ij}^{(N=3)}$$

Momentum sum rule controls x^1 term,
→ can drop it.

$$\int_0^1 dx\, x\, P_{ij}(x) \;\equiv\; -\gamma_{ij}^{(N=2)}$$

LL resummed formula

Richards, Stirling, Ellis, NPB229, 317, 1983

$$\frac{1}{\sigma_0}\frac{d\sigma}{d\cos\chi} = \frac{\alpha_s(\sqrt{z}Q)}{16\pi\ z} \sum_{i,j=q,g} \gamma_{ij}^{(0)} \left[\frac{\alpha_s(\sqrt{z}Q)}{\alpha_s(Q)}\right]_{jq}^{-\gamma^{(0)}/\beta_0}$$

$$\gamma_{ij}^{(0)} = \begin{bmatrix} \frac{25}{6}C_F & -\frac{7}{15}n_f \\ -\frac{7}{6}C_F & \frac{14}{5}C_A + \frac{2}{3}n_f \end{bmatrix} \qquad \beta_0 = \frac{11}{3}C_A - \frac{2}{3}n_f$$

One-loop (LO) $N=3$ time-like moments

To expand back into fixed order:

$$\frac{\alpha_s(\sqrt{z}Q)}{\alpha_s(Q)} = \left[1 + \beta_0 \frac{\alpha_s(Q)}{4\pi} \ln z\right]^{-1}$$

Beyond LL as $z \to 0$

LD, Moult, Zhu, to appear

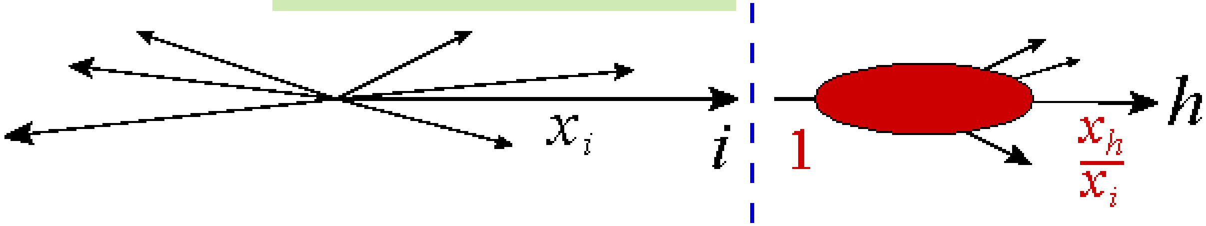
- Factorize on single parton states, similar to production of identified hadrons h with momentum $p_h = x \times Q/2$

$$\frac{d\sigma(e^+e^- \to h + X)}{dx} = \sum_{i=q,g} \int_0^1 dx_i \, \frac{d\sigma(e^+e^- \to i + X)}{dx_i} D_{i\to h}(x/x_i)$$

perturbative hard function, computed to NNLO + evolution

nonperturbative fragmentation function

…, Mitov, Moch, Vogt, 2006
Moch, Vogt, 0709.3899,
Almasy, Moch, Vogt, 1107.2263

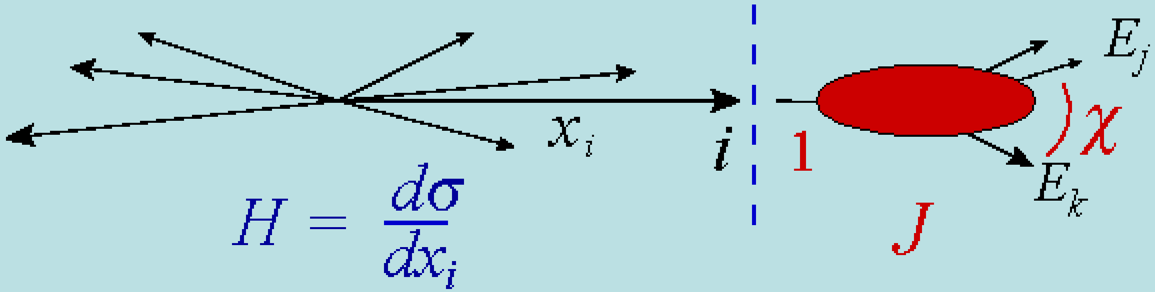


All orders factorization formula

Cumulant $\Sigma\left(z, \ln\frac{Q^2}{\mu^2}, \alpha_s(\mu)\right) \equiv \frac{1}{\sigma_0}\int_0^z dz' \frac{d\sigma}{dz}\left(z', \ln\frac{Q^2}{\mu^2}, \alpha_s(\mu)\right)$

$$\Sigma\left(z, \ln\frac{Q^2}{\mu^2}, \alpha_s(\mu)\right) = \int_0^1 dx\, x^2 \vec{J}\left(\ln\frac{zx^2Q^2}{\mu^2}, \alpha_s(\mu)\right) \dot{\vec{H}}\left(\ln\frac{Q^2}{\mu^2}, \alpha_s(\mu)\right)$$

- Reuses hard function



- Replaces nonperturbative fragmentation function with perturbative jet function J which includes the small angle EEC measurement
- J depends on its only physical scale: $q_T^2 \approx (\chi x Q/2)^2 \approx zx^2Q^2$

Evolution of jet function

- To resum large logs, evolve jet function from its natural scale, $\mu = \sqrt{z}Q$ up to natural scale of hard function, $\mu = Q$
- Hard function evolves with time-like splitting kernel, $P_T(y,\mu)$:

$$\frac{d\vec{H}(x)}{d\ln\mu^2} = -\int_x^1 \frac{dy}{y} \widehat{P}_T(y,\mu) \cdot \vec{H}(x/y)$$

- Σ is RGE invariant, i.e. independent of μ
- Leads to evolution equation for J:

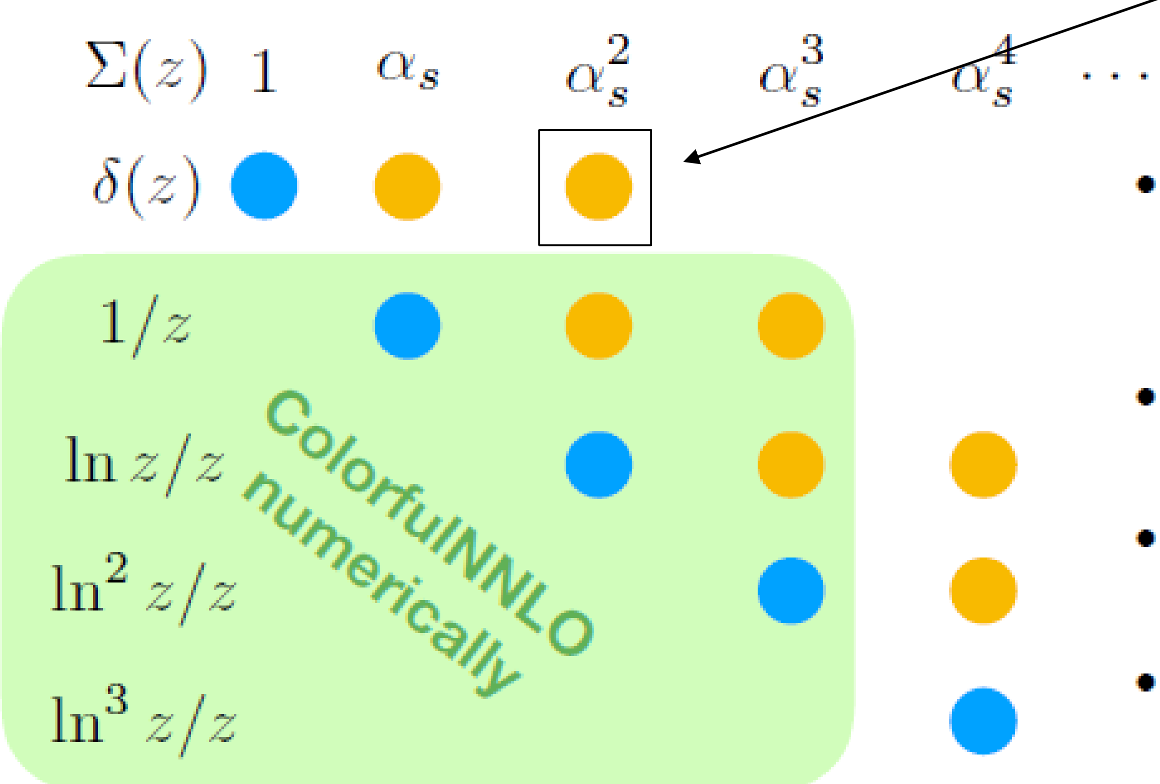
$$\frac{d\vec{J}\left(\ln\frac{zQ^2}{\mu^2}, \alpha_s\right)}{d\ln\mu^2} = \int_0^1 dy\, y^2\, \vec{J}\left(\ln\frac{zy^2Q^2}{\mu^2}\right) \cdot \widehat{P}_T(y,\mu)$$

- LL evolution only uses $N=3$ time-like moments (y^2)
- Beyond LL, need "nearby" moments, $\ln y \leftrightarrow \frac{\partial}{\partial N}$

Counting the order

- ● LL — Konishi, Ukawa, Veneziano, 1979
- ● NLL + NNLL — Dixon, Moult, HXZ, 2019, This talk

Get these jet function constants indirectly using sum rules



- To get to NNLL require:

- NNLO splitting kernel
  Moch, Vermaseren, Vogt
- NNLO hard function
  Mitov, Moch, 2006;
  Almasy, Moch, Vogt, 2011
- NNLO jet function

Very challenging!

Sum rules

- Energy conservation, $Q^2 = (\sum E_i)^2 = \sum E_i E_j$ implies a sum rule,

$$\int_0^1 dz \frac{d\sigma}{dz} = \sigma_{\text{tot}}$$

- Momentum conservation involving

$$p_i \cdot p_j = E_i E_j (1 - \cos\chi)$$

- implies a second sum rule Korchemsky, 1905.01444, Kologlu et al., 1905.01311,

$$\int_0^1 dz\, z \frac{d\sigma}{dz} = \int_0^1 dz\, (1 - z) \frac{d\sigma}{dz} = \frac{1}{2} \sigma_{\text{tot}}$$

Using sum rule(s) to get $\alpha_s^2\ \delta(z)$

- Total cross section σ_{tot} known, for e^+e^- and Higgs, to very high order, e.g. Herzog, Ruijl, Ueda, Vermaseren, Vogt, 1707.01044
- First sum rule requires both $\delta(z)$ and $\delta(1-z)$ terms
- Second sum rule decouples them.
- But one can also get $\alpha_s^2\,\delta(1-z)$ in course of resumming $z \rightarrow 1$ Zhu, et al. (2019)
- Know α_s^2 distribution for e^+e^- and Higgs analytically for $0 < z < 1$, so we can integrate it, up to delta functions.
- $\delta(z)$ coefficients involve sum of H and J. Get H from Almasy, Mitov, Moch, Vogt
- → Use the two $\delta(z)$ coefficients to fix 2-loop J_q, J_g

Two loop jet constants in QCD

$$j_2^q = C_F n_f \left(\frac{9}{5}\zeta_2 + \frac{703847}{24000}\right) + C_F C_A \left(-76\zeta_4 + 280\zeta_3 + \frac{1063}{15}\zeta_2 - \frac{164883727}{324000}\right)$$
$$+ C_F^2 \left(152\zeta_4 - 478\zeta_3 - 106\zeta_2 + \frac{3498505}{5184}\right),$$
$$j_2^g = n_f^2 \left(-\frac{8}{15}\zeta_2 + \frac{2344}{1125}\right) + C_F n_f \left(4\zeta_3 + \frac{14}{5}\zeta_2 - \frac{1528667}{108000}\right)$$
$$+ C_A n_f \left(\frac{44}{5}\zeta_3 - \frac{127}{25}\zeta_2 + \frac{68111303}{1620000}\right) + C_A^2 \left(76\zeta_4 - \frac{1054}{5}\zeta_3 - \frac{2159}{75}\zeta_2 + \frac{133639871}{810000}\right)$$

- In pure N=1 SYM ($C_F \to C_A$, $n_f \to C_A$):

$$j_2^{q,\ \mathcal{N}=1} = C_A^2\left(76\,\zeta_4 - 198\,\zeta_3 - \frac{100}{3}\zeta_2 + \frac{78117}{400}\right)$$

$$j_2^{g,\ \mathcal{N}=1} = C_A^2\left(76\,\zeta_4 - 198\,\zeta_3 - \frac{158}{5}\zeta_2 + \frac{263197}{1350}\right)$$

Solve jet evolution, first for N=4 SYM

- Scale invariance → solution for N=4 SYM a pure power law:

$$J\Big(\frac{zQ^2}{\mu^2},\alpha_s\Big) = C_J(\alpha_s)\Big(\frac{zQ^2}{\mu^2}\Big)^{\gamma_J^{\mathcal{N}=4}(\alpha_s)}$$

- Insert into evolution equation, find that

$$\begin{aligned} 2\gamma_J^{\mathcal{N}=4} &= -2\int_0^1 dy\, y^{2+2\gamma_J^{\mathcal{N}=4}} P_{T,\text{uni.}}(y) \\ &= 2\gamma_T^{\mathcal{N}=4}(N = 1 + 2\gamma_J^{\mathcal{N}=4}) \end{aligned}$$

- Using "time-like space-like reciprocity relation" Drell, Levy, Yan (1969), Gribov, Lipatov (1972), …, Basso, Korchemsky, hep-th/0612247
this is actually the space-like N=3 moment
(N=1 for "universal" N=4 SYM anomalous dimension):

$$\gamma_J^{\mathcal{N}=4} = \gamma_S^{\mathcal{N}=4}(N=1)$$

N=4 SYM result

- Jet function solution leads to:

$$\Sigma(z) = \frac{1}{2} C(\alpha_s)\, z^{\gamma_J^{\mathcal{N}=4}(\alpha_s)}$$

where

$$\gamma_J^{\mathcal{N}=4}(\alpha_s) = \gamma_S^{\mathcal{N}=4}(1,\alpha_s) = \frac{C_A\alpha_s}{\pi} + \left(-\frac{\zeta_3}{2} + \zeta_2 - 2\right)\left(\frac{C_A\alpha_s}{\pi}\right)^2 + \left(\frac{3}{2}\zeta_5 + \frac{3}{8}\zeta_4 - \frac{3}{2}\zeta_3 - 4\zeta_2 + 8\right)\left(\frac{C_A\alpha_s}{\pi}\right)^3 + \left(-\frac{69}{16}\zeta_7 + \frac{1}{2}\zeta_2\zeta_5 - \frac{5}{16}\zeta_3\zeta_4 + \frac{9}{4}\zeta_3^2 - \frac{107}{32}\zeta_6 + 8\zeta_5 - \frac{13}{2}\zeta_2\zeta_3 - \frac{23}{8}\zeta_4 + 7\zeta_3 + 24\zeta_2 - 40\right)\left(\frac{C_A\alpha_s}{\pi}\right)^4 + \mathcal{O}(\alpha_s^5)\,.$$

Kotikov, Lipatov, Rej, Staudacher, Velizhanin, 0704.3586;
Kotikov, Velizhanin, hep-ph/0501274

- Sum rule gives:

$$C(\alpha_s) = 1 - \frac{C_A\alpha_s}{\pi} + \left(\frac{11}{4}\zeta_4 - 3\zeta_2 + 7\right)\left(\frac{C_A\alpha_s}{\pi}\right)^2 + \mathcal{O}(\alpha_s^3)$$

- Agrees with recent “space-like” analysis
  Kologlu et al., 1905.01311; Korchemsky, 1905.01444

Solving jet evolution for N=1 SYM

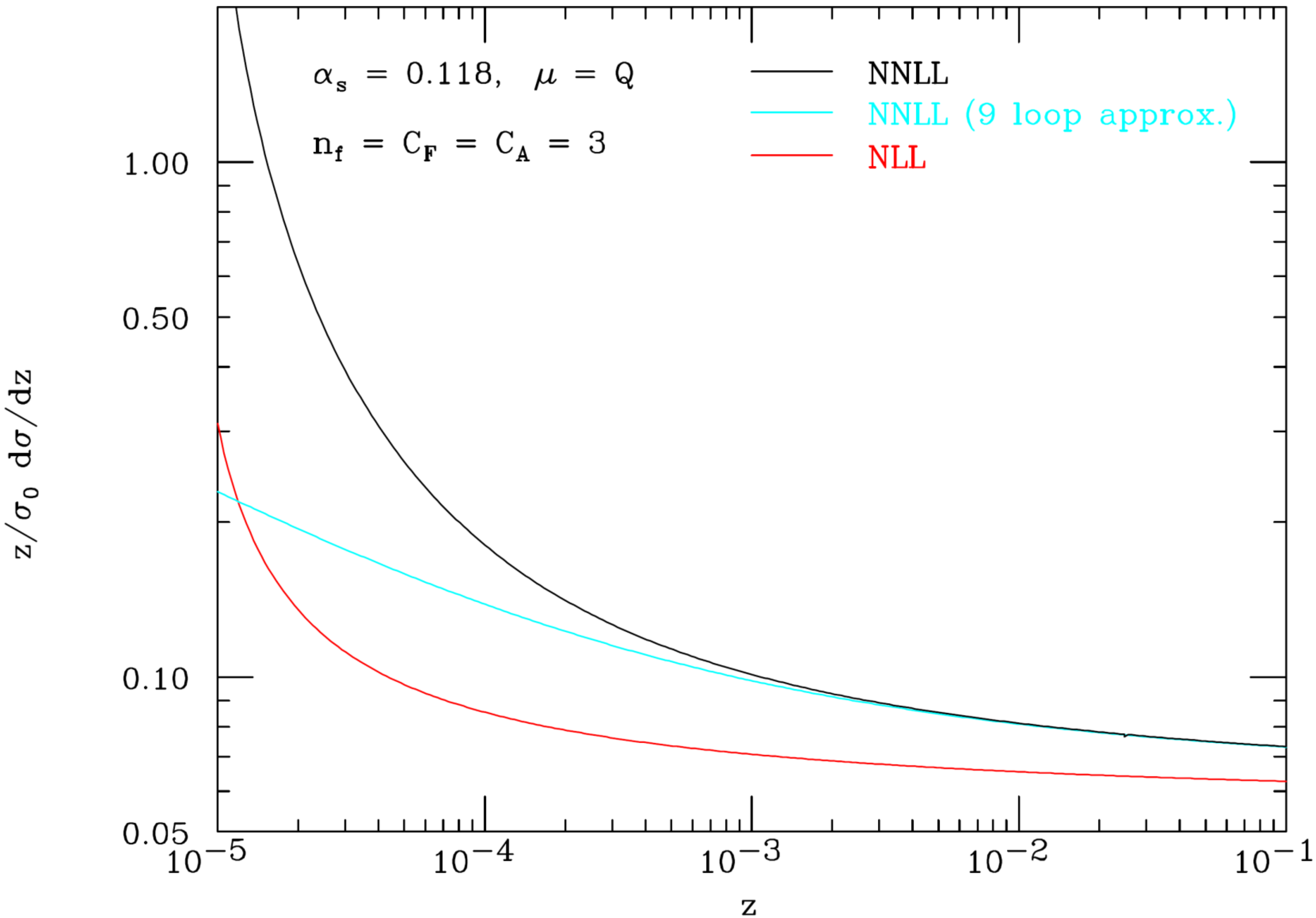
- We can do it exactly at NNLL, because 2 x 2 matrix equation is effectively 1 x 1. Solution is:

$$\Sigma^{\mathcal{N}=1}_{\rm NNLL}(z) = c_1^S(\alpha_s) + c_2^S(\alpha_s)\ln z + c_3^S(\alpha_s)\frac{\ln z}{1+\beta_0 a_s \ln z}$$
$$+ c_4^S(\alpha_s)\ln[1+\beta_0 a_s \ln z]$$
$$+ c_5^S(\alpha_s)\ln\left(1 - 2C_A a_s \frac{\ln[1+\beta_0 a_s \ln z]}{1+\beta_0 a_s \ln z}\right)$$

$$c_1^H = \frac{1}{2} + \frac{69}{8}a + a^2\left(22\zeta_4 - 66\zeta_3 - \frac{95}{3}\zeta_2 + \frac{81949}{432}\right),$$
$$c_1^\gamma = \frac{1}{2} + \frac{13}{24}a + a^2\left(22\zeta_4 - 44\zeta_3 + \frac{22}{9}\zeta_2 + \frac{2911}{162}\right),$$

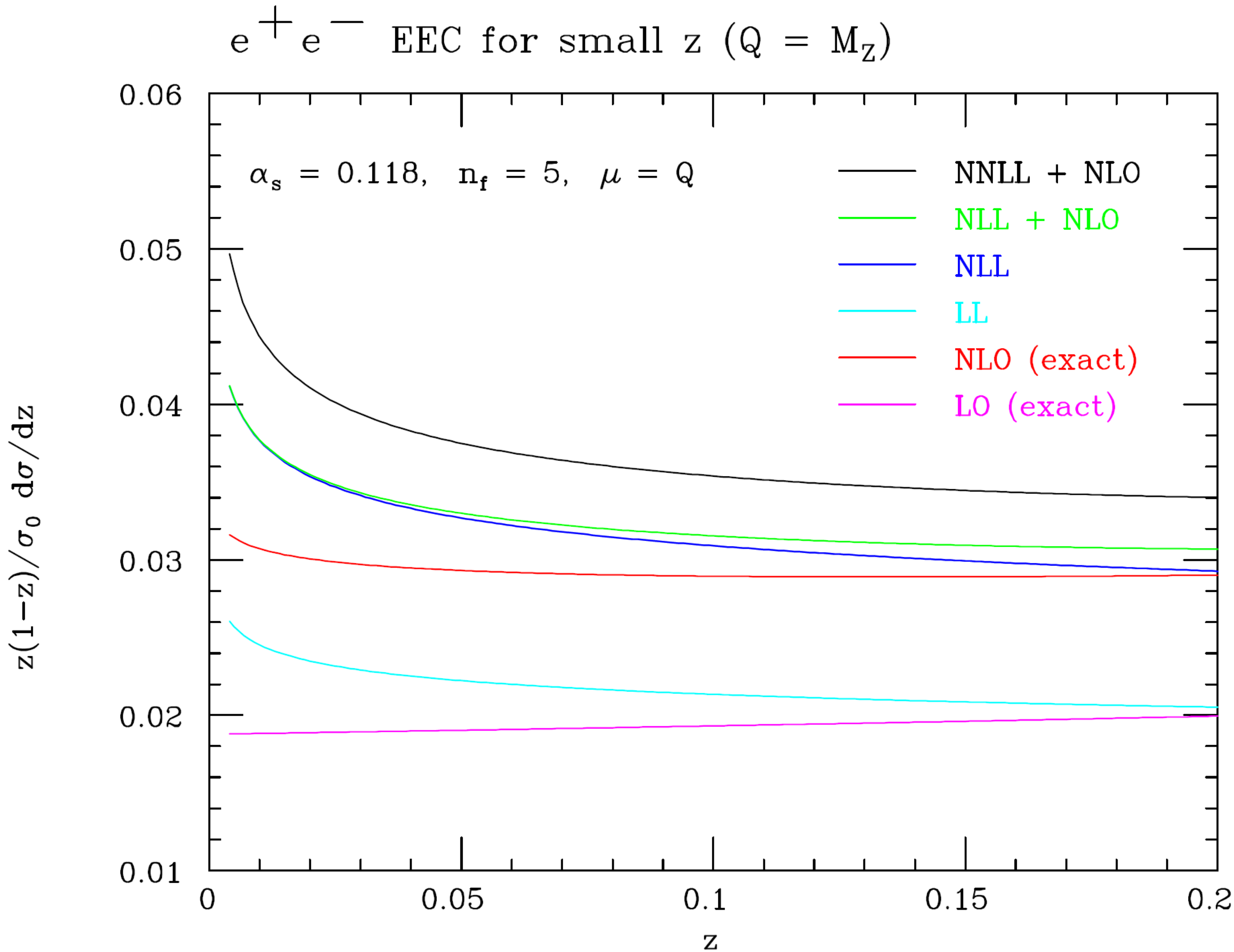
etc., $a = \frac{C_A \alpha_s}{4\pi}$

e^+e^- EEC in N=1 SYM for small z ($Q = M_Z$)

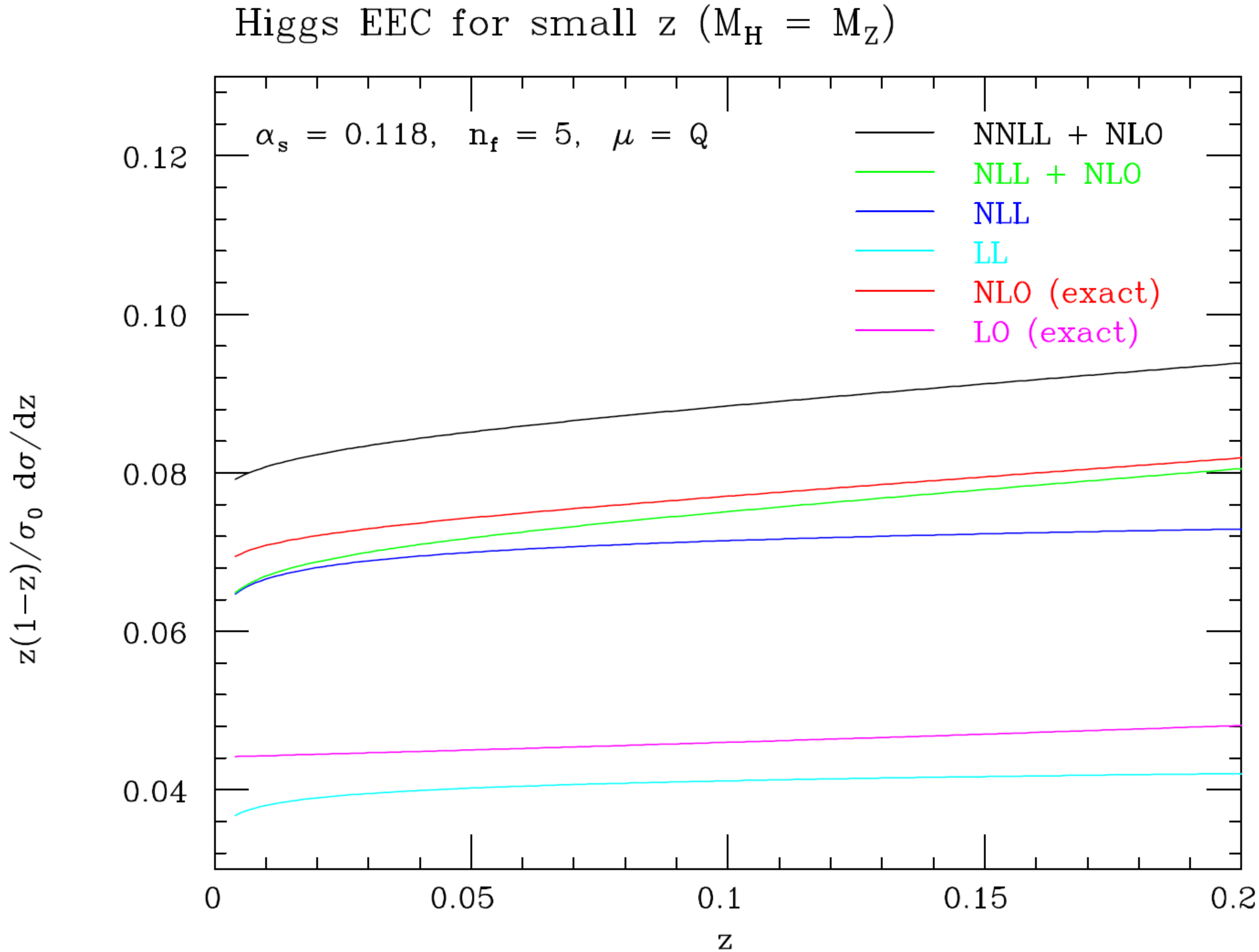


Solving jet evolution for QCD

- Equation is 2 x 2. For $z > 0.004$, we can get away with solving it order by order through 9 loops
- Do it for both e^+e^- and Higgs to illustrate different behavior of quark and gluon jets.
- Competition between β function and splitting plays out very differently for quarks and gluons, $C_F = \frac{4}{3}$ versus $C_A = 3$.

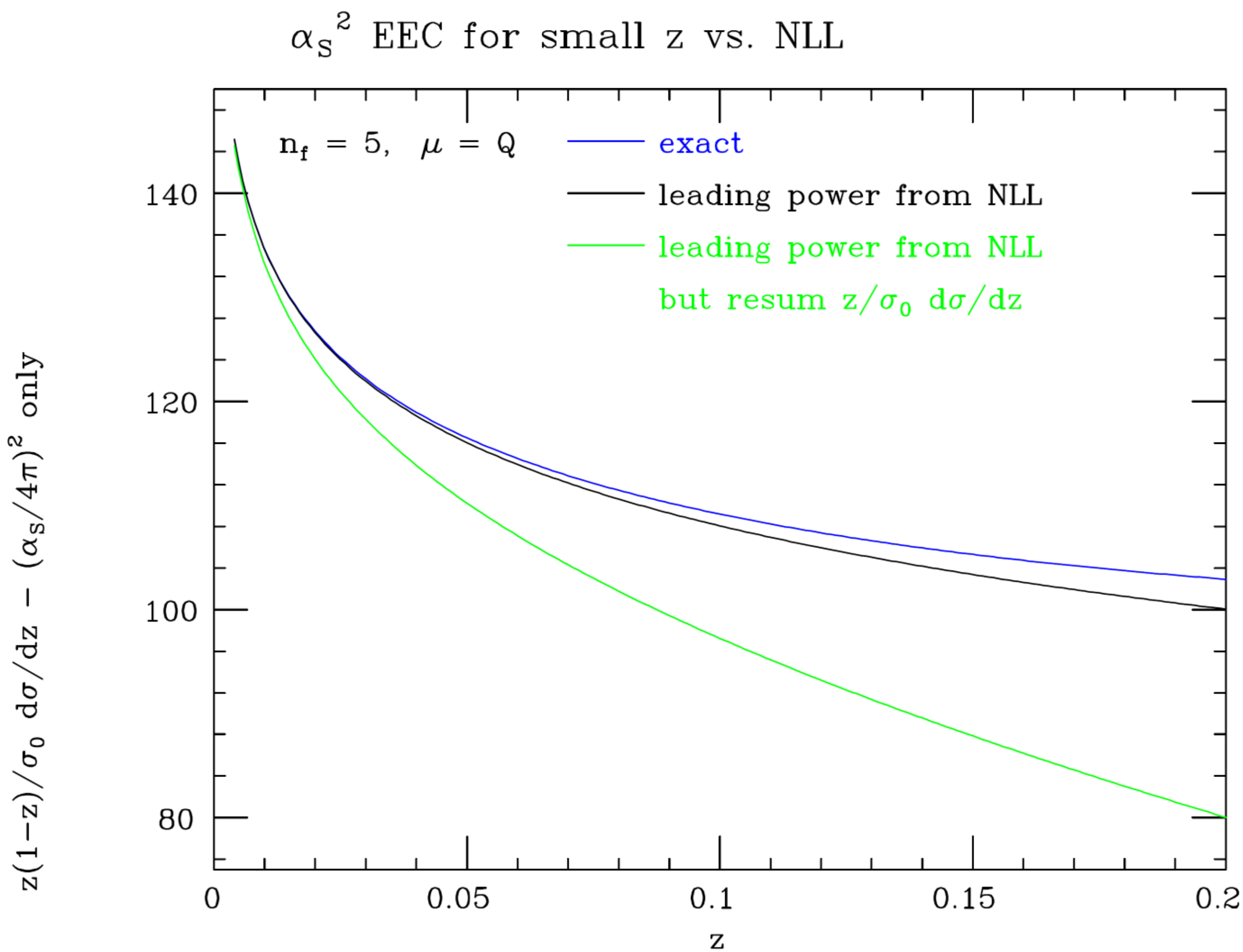


e+e− EEC for small z (Q = M_Z)
α_s = 0.118, n_f = 5, μ = Q
NNLL + NLO
NLL + NLO
NLL
LL
NLO (exact)
LO (exact)
z(1−z)/σ_0 dσ/dz
z
0.01
0.02
0.03
0.04
0.05
0.06
0
0.05
0.1
0.15
0.2



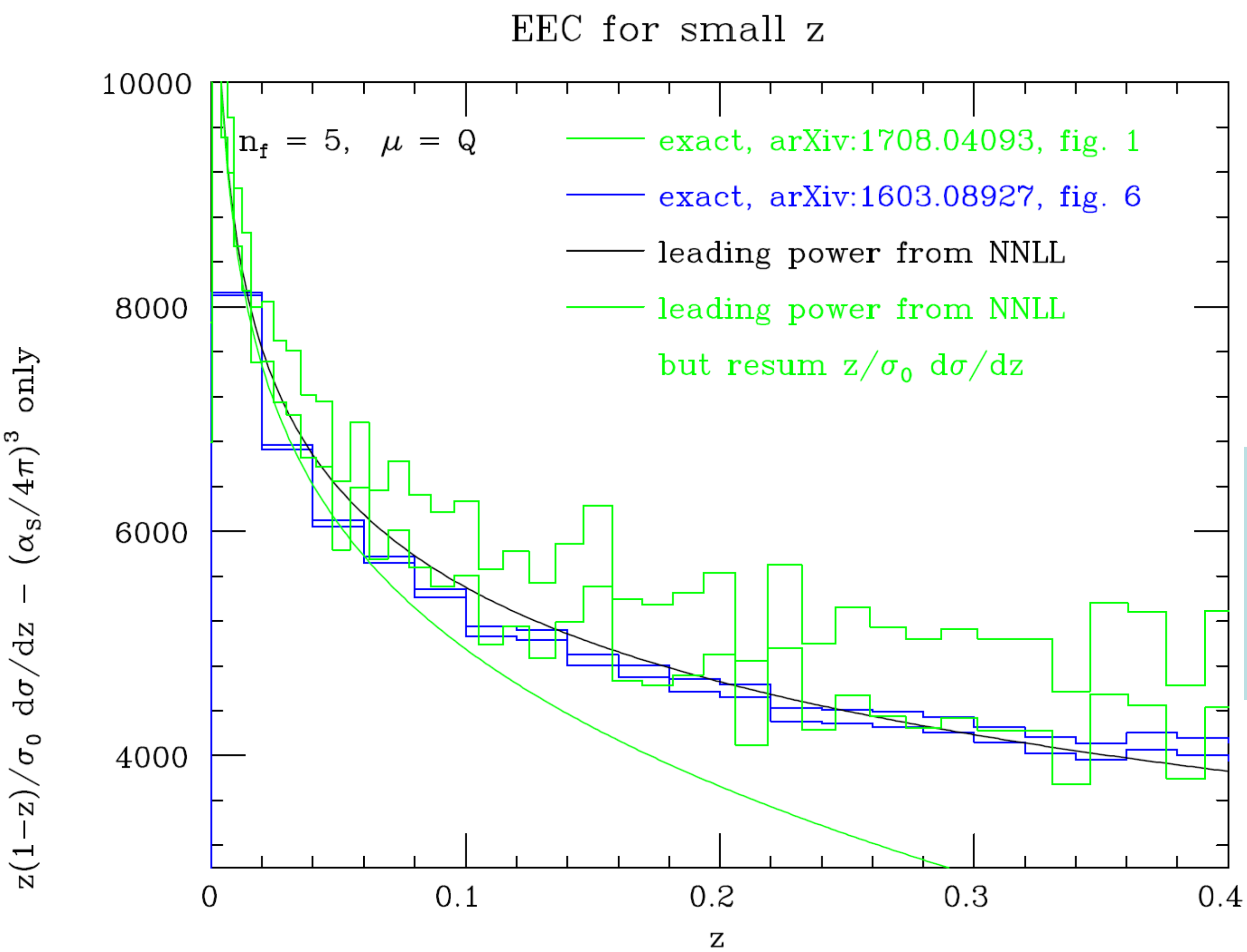
Higgs EEC for small z ($M_H = M_Z$)
$\alpha_s = 0.118$, $n_f = 5$, $\mu = Q$
NNLL + NLO
NLL + NLO
NLL
LL
NLO (exact)
LO (exact)
0.12
0.10
0.08
0.06
0.04
0
0.05
0.1
0.15
0.2
z
$z(1-z)/\sigma_0\ d\sigma/dz$

Comparison with fixed order at α_s^2



Power corrections small for $z(1\text{-}z)$ normalization

Comparison with fixed order at α_s^3



Thanks to
Del Duca, Duhr,
Kardos, Somogyi,
Trocsanyi, Tulipant
for numbers

Leading power remarkably accurate all the way out to $z = 0.4$

Conclusions

- All order time-like factorization formula for small angle EEC for generic theories.
- Explicitly computed through NNLL, two orders more accurate than previous jet calculus approach.
- Quite different behavior for quarks (e^+e^-) vs. gluons (Higgs)
- Good agreement with ColorfulNNLO results at $O(\alpha_s^3)$
- Also can see Landau pole behavior explicitly in N=1 SYM
- Time-like/space-like reciprocity of twist 2 anomalous dimensions relates formula to other N=4 SYM approaches

Outlook

- May be possible to go to NNNLL in QCD, at least approximately. Also approximate NNNLO?
- Same jet functions also apply to suitable "jet substructure" observables at LHC, could use to discriminate quark/gluon jets at LHC.
- May eventually lead to more precise value of α_s, as well as more precise jet substructure understanding at LHC

Extra Slides

Reverse unitarity

Anastasiou, Melnikov, hep-ph/0207004;
Anastasiou, LD, Melnikov, Petriello, hep-ph/0312266

- Phase space integral over final-state partons is like a loop integral with $\delta(p_i^2)$ factor for every propagator crossing the cut, and with one extra delta function, which can be turned into a fake propagator:

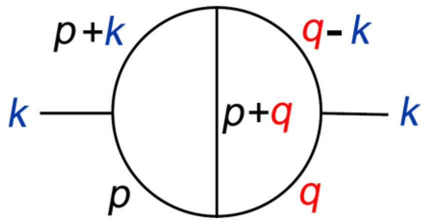
$$\delta[\mathcal{M}_{ij}(\chi)] = \frac{1}{2\pi i}\left[\frac{1}{\mathcal{M}_{ij}(\chi) - i\varepsilon} - \frac{1}{\mathcal{M}_{ij}(\chi) + i\varepsilon}\right]$$

where
$$\begin{aligned}\mathcal{M}_{ij}(\chi) &= (p_i \cdot Q\, p_j \cdot Q)(\vec{n}_i \cdot \vec{n}_j - \cos\chi) \\ &= (p_i \cdot Q\, p_j \cdot Q)(1 - \cos\chi) - p_i \cdot p_j\end{aligned}$$

- Nonlinear in parton momenta p_i , p_j
- Sum over i,j

Integration by parts (IBP)

- Multi-loop integration technology



Chetyrkin, Tkachov (1981)

$$0 = \int d^D p d^D q \ldots \frac{\partial}{\partial q^\mu} \frac{k^\mu}{p^2 q^2 (p+q)^2 \ldots}$$

- Reduces problem to system of linear equations, initially solved recursively by MINCER, now by Laporta algorithm, in terms of "master integrals"

Gorishnii, Larin, Surguladze, Tkachov (1989)

Laporta, hep-ph/0102033

No-scale problem $\Rightarrow$
like total hadronic cross section
maximal analytic simplicity:
pure numbers, Riemann zeta values

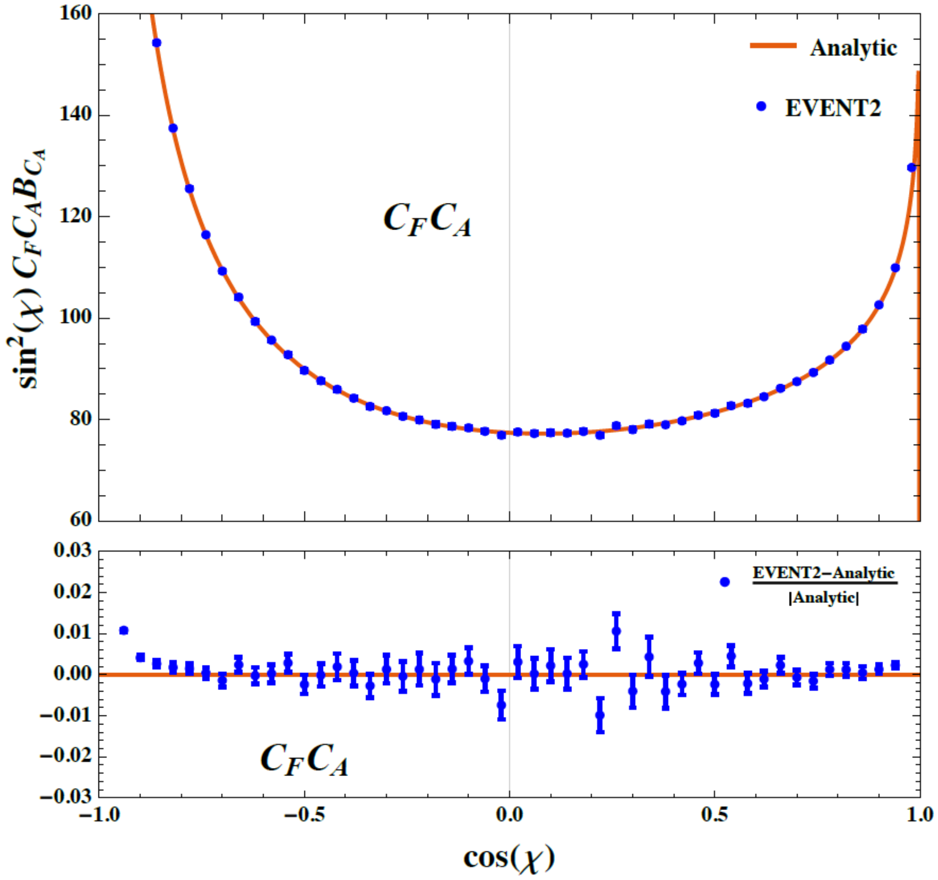
$$\zeta(n) = \sum_{k=1}^{\infty} \frac{1}{k^n}$$

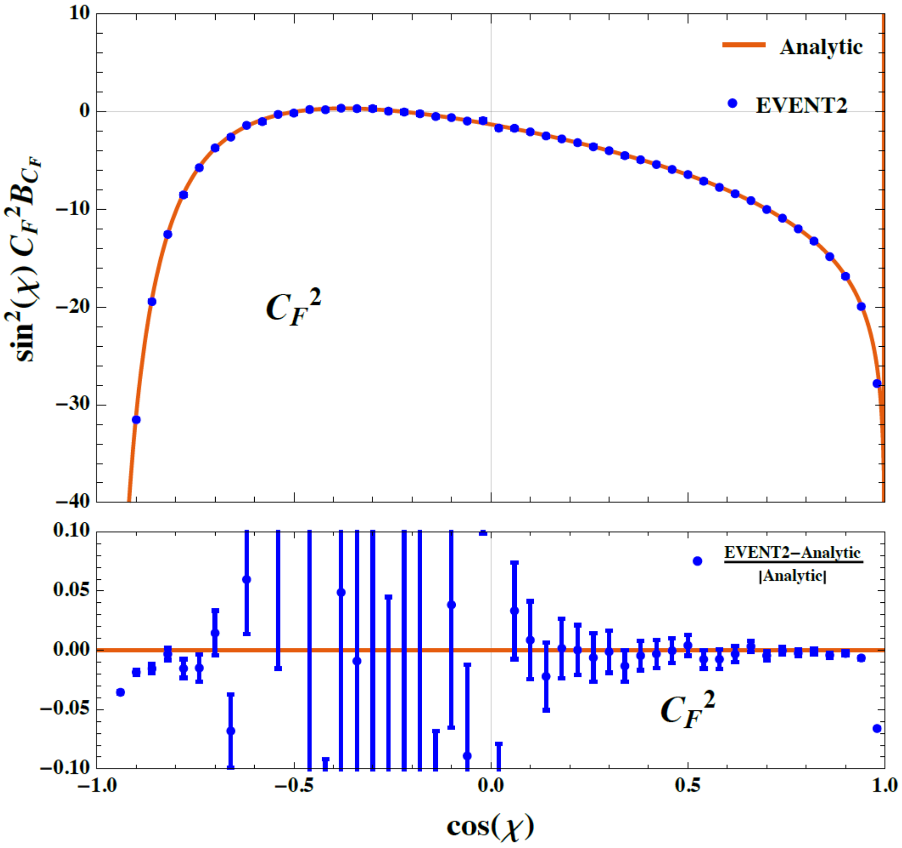
$$\begin{aligned}\frac{R_{e^+e^-}}{R^{(0)}} &= 1 + \frac{\alpha_s}{\pi} + \left(\frac{\alpha_s}{\pi}\right)^2 \left[-11\zeta(3) + \frac{365}{24} + n_f\left(\frac{2}{3}\zeta(3) - \frac{11}{12}\right)\right] \\ &+ \left(\frac{\alpha_s}{\pi}\right)^3 \Big[\frac{275}{6}\zeta(5) - \frac{1103}{4}\zeta(3) - \frac{121}{8}\zeta(2) + \frac{87029}{288} \\ &+ n_f\left(-\frac{25}{9}\zeta(5) + \frac{262}{9}\zeta(3) + \frac{11}{6}\zeta(2) - \frac{7847}{216}\right) \\ &+ n_f^2\left(-\frac{19}{27}\zeta(3) - \frac{1}{18}\zeta(2) + \frac{151}{162}\right)\Big]\end{aligned}$$

EEC is "next-to-simplest case"

QCD: NLO program EVENT2 validated

M. Seymour





Analytic properties of QCD moments

- With analytic formulae, compute the integrals

$$B_N = \int_0^1 dz\ z^N B(z)$$

numerically to high accuracy, for each color coefficient

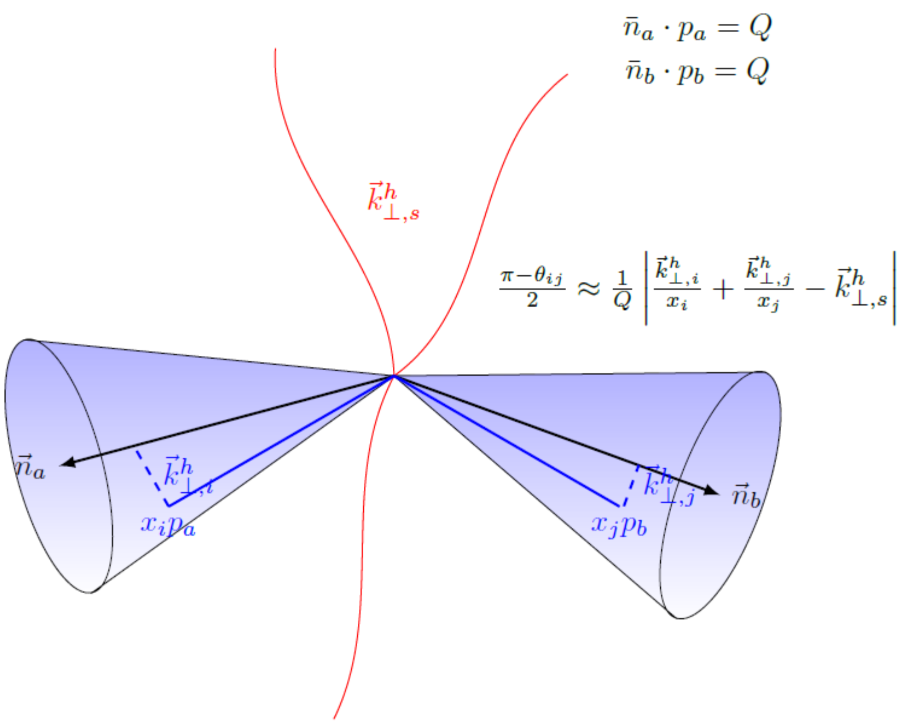
- Using PSLQ, it is always of the form

$$B_N = r_N^{(4)}\zeta(4) \ + \ r_N^{(3)}\zeta(3) \ + \ r_N^{(2)}\zeta(2) \ + \ r_N^{(0)}$$

where the $r_N^{(w)}$ are rational numbers.

- E.g. $B_3(C_A) = -\frac{207}{2}\zeta(4) + \frac{14902}{35}\zeta(3) - \frac{553}{450}\zeta(2) - \frac{2369041}{5040}$
- Could they be zeta values at higher loop orders too?
- Expression for general N in terms of $\psi(N)$ functions?

$z \rightarrow 1$ (cont.)



Moult, Zhu, 1801.02627

Soft gluons contribute, but only via recoil, by deflecting the hard quark jet

- Factorization theorem recently proved: Relate EEC to back-to-back production of identified hadrons Collins, Soper 1981-1982
- Should allow NNNLL resummation soon