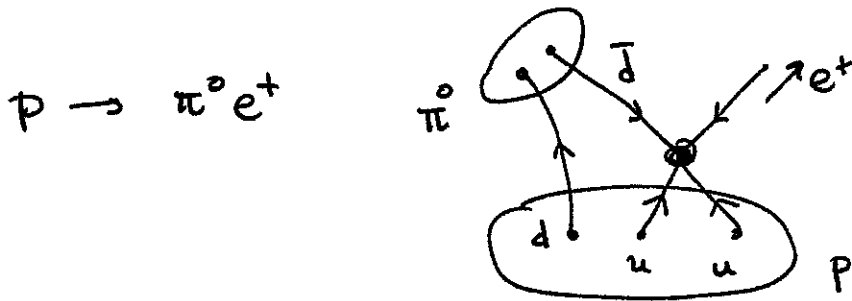


Physics 332 - Problem Set #7

Solutions

1.) a.) Consider for definiteness



for a $1 \rightarrow 2$ deg, iM has dimensions of $(\text{mass})^4$

so
$$iM \sim \frac{2}{m_{\Sigma}^2} \cdot m_p^3$$

we might want to add a factor $\partial U \sim 4\pi \left(\frac{g^2}{4\pi}\right) \sim \frac{1}{2}$

It makes little difference. $\hookrightarrow \sim \frac{1}{24}$

$$I \sim \frac{1}{2m_p} \frac{1}{8\pi} \cdot \left| \frac{2m_p^3}{m_{\Sigma}^2} \right|^2$$

$$\sim \frac{1}{4\pi} \frac{m_p^5}{m_{\Sigma}^4} \sim \frac{1}{4\pi} \frac{(1 \text{ GeV})^5}{(10^{16} \text{ GeV})^4}$$

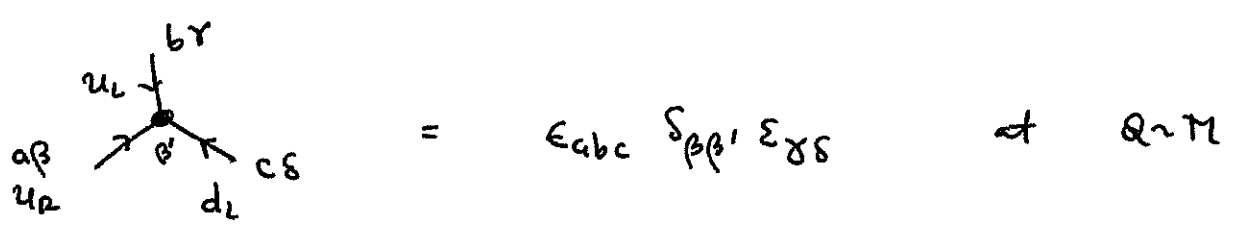
so

$$\Gamma \sim 10^{-65} \text{ GeV}^{+1} \cdot (6.58 \times 10^{-25} \text{ GeV-sec})^{-1}$$

$$\sim (10^{41} \text{ sec})^{-1} \sim (10^{34} \text{ yr})^{-1}$$

b) for $[O_\beta]_M = \epsilon_{abc} \epsilon^{\gamma\delta} u_{R\alpha\beta} u_{Lb\gamma} d_{Lc\delta}$

a renormalized vertex is



this vertex is corrected by QCD diagrams:



to compute these diagrams, we need the Feynman identities for

$$\sigma^\mu = (1, \vec{\sigma})^\mu \quad \bar{\sigma}^\mu = (1, -\vec{\sigma})^\mu \quad \gamma^\mu = \begin{pmatrix} 0 & \sigma^\mu \\ \bar{\sigma}^\mu & 0 \end{pmatrix}$$

$$(\sigma^\mu)_{\alpha\beta} (\sigma^\nu)_{\gamma\delta} = -2 \epsilon_{\alpha\gamma} \epsilon_{\beta\delta}$$

check: $(\sigma^\mu \bar{\sigma}^\nu \sigma^\mu)_{\alpha\delta} = -2 (\sigma^\nu)_{\alpha\delta} \stackrel{?}{=} -2 (i\sigma^2)_{\alpha\gamma} (\bar{\sigma}^\nu)^T_{\gamma\beta} (i\sigma^2)_{\beta\delta}$

$$= +2 [\sigma^2 (\bar{\sigma}^\nu)^T \sigma^2]_{\alpha\delta}$$

$$= 2 (\sigma^\nu)_{\alpha\delta}$$

✓

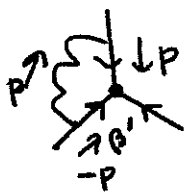
$$\underline{\text{all}}$$

$$(\bar{\sigma}^\mu)_{\alpha\beta} (\sigma_\mu)_{\gamma\delta} = 2 \delta_{\alpha\delta} \delta_{\gamma\beta}$$

$$\text{check: } (\bar{\sigma}^\mu \cdot 1 \cdot \sigma_\mu)_{\alpha\beta} \stackrel{?}{=} 2 \delta_{\alpha\beta} \delta_{\gamma\gamma}$$

$$= 4 \delta_{\alpha\beta} \stackrel{?}{=} 4 \delta_{\alpha\beta}$$

now



$$= (ig)^2 \int \frac{d^4 p}{(2\pi)^4} \frac{-i}{p^2} \epsilon_{a'b'c} (t^A)_{a'a} (t^A)_{b'b}$$

$$\epsilon_{\gamma'\delta'} \delta_{c's'} \cdot \left(\frac{-i p \cdot \bar{\sigma}}{p^2} \right)_{\beta'\alpha'} (\sigma^\mu)_{\beta'\alpha'}$$

$$\cdot \left(\frac{i p \cdot \sigma}{p^2} \sigma_\mu \right)_{\gamma'\delta'}$$

$$= ig^2 \epsilon_{\gamma'\delta'} \epsilon_{a'b'c} (t^A)_{a'a} (t^A)_{b'b}$$

$$\cdot \int \frac{d^4 p}{(2\pi)^4} \frac{1}{(p^2)^3} (p \cdot \bar{\sigma})_{\beta'\alpha'} (p \cdot \sigma)_{\gamma'\delta'} \cdot 2 \delta_{\eta\theta} \delta_{\theta\eta}$$

now

$$\epsilon_{a'b'c} (t^A)_{a'a} (t^A)_{b'b} = t^A t^A \text{ for } 3 \times 3 \rightarrow \bar{3}$$

$$= \frac{1}{2} [C_2(\bar{3}) - C_2(3) - C(3)]$$

$$= \frac{1}{2} \left[\frac{4}{3} - \frac{4}{3} - \frac{4}{3} \right]$$

$$= -\frac{2}{3}$$

$$\begin{aligned}
&= -\frac{2}{3} i g^2 \varepsilon_{\gamma\delta} \varepsilon_{abc} \cdot \left(\frac{i}{4\pi}\right)^2 \frac{\Gamma(2-d/2)}{2} \cdot \frac{1}{2} g^{\alpha\beta} \\
&\quad \cdot 2 (\bar{\sigma}_\alpha)_{\beta'\gamma} (\sigma_\beta)_{\delta'\delta} \\
&= -\frac{2}{3} i g^2 \cdot \left(\frac{i}{4\pi}\right)^2 \frac{2}{\epsilon} \cdot \frac{1}{4} \cdot 2 \cdot 2 \varepsilon_{abc} \varepsilon_{\gamma\delta} \delta_{\beta\beta'} \\
&= \frac{2}{3} g^2 \frac{1}{(4\pi)^2} \frac{2}{\epsilon} \delta_{\beta\beta'} \varepsilon_{\gamma\delta} \varepsilon_{abc}
\end{aligned}$$



gives the same result.

$$\begin{aligned}
\text{triangle}^p &= \delta_{\beta\beta'} (ig)^2 \int \frac{d^d p}{(2\pi)^d} \frac{-i}{p^2} \varepsilon_{ab'c'} (t^A)_{b'b} (t^A)_{c'c} \\
&\quad \varepsilon_{\gamma'\delta'} \left(\frac{-i p \cdot \sigma}{p^2} \bar{\sigma}^\gamma \right)_{\beta'\gamma} \left(\frac{i p \cdot \sigma}{p^2} \sigma^\delta \right)_{\delta'\delta} \\
&= \left(-\frac{2}{3}\right) i g^2 \delta_{\beta\beta'} \int \frac{d^d p}{(2\pi)^d} \frac{1}{(p^2)^3} \varepsilon_{abc} \varepsilon_{\gamma'\delta'} (p \cdot \sigma)_{\gamma'\eta} (p \cdot \sigma)_{\delta'\eta} (-2 \varepsilon_{\eta\gamma} \varepsilon_{\delta\delta'}) \\
&= -i \frac{2}{3} g^2 \delta_{\beta\beta'} \varepsilon_{abc} \varepsilon_{\gamma'\delta'} \frac{i}{(4\pi)^2} \frac{\Gamma(2-d/2)}{2} \frac{1}{2} g^{\alpha\beta} (\sigma^\alpha)_{\gamma'\eta} (\sigma^\alpha)_{\delta'\eta} \varepsilon_{\eta\gamma} \varepsilon_{\delta\delta'} \\
&\quad (-2) \\
&= \frac{2}{3} g^2 \delta_{\beta\beta'} \varepsilon_{abc} \varepsilon_{\gamma'\delta'} \frac{1}{(4\pi)^2} \frac{2}{\epsilon} \frac{1}{4} (-2 \varepsilon_{\gamma'\delta'} \varepsilon_{\eta\delta}) \varepsilon_{\eta\gamma} \varepsilon_{\delta\delta'} (-2) \\
&= \frac{2}{3} g^2 \delta_{\beta\beta'} \varepsilon_{abc} \varepsilon_{\gamma\delta} \frac{1}{(4\pi)^2} \frac{2}{\epsilon} \frac{1}{4} \cdot 4 \cdot 4
\end{aligned}$$

5.

$$\text{triangle} = \text{triangle} = \frac{2}{3} \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \delta_{\beta\beta'} \delta_{\gamma\gamma'} \epsilon_{abc}$$

$$\text{triangle} = \frac{8}{3} \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \delta_{\beta\beta'} \delta_{\gamma\gamma'} \epsilon_{abc}$$

$$\delta_0 = -\frac{12}{3} \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} = -4 \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon}$$

recall $\delta_2 = -\frac{g^2}{(4\pi)^2} \cdot \frac{4}{3} \cdot \frac{2}{\epsilon}$

with $\mu_0 \rightarrow \mu_0 \Lambda^2 / M^2$ $M \frac{\partial}{\partial M} (\frac{2}{\epsilon}) = -2$

$$\begin{aligned} \gamma_0 &= M \frac{\partial}{\partial M} (-\delta_0 + \frac{3}{2} \delta_2) \\ &= (-8 + \frac{3}{2} \cdot \frac{4}{3} \cdot 2) \frac{g^2}{(4\pi)^2} \end{aligned}$$

$$\gamma_0 = -4 \frac{g^2}{(4\pi)^2}$$

Why $\gamma_0 = \frac{a_0}{(4\pi)^2} g^2$, this operator is enhanced at

large distances \hookrightarrow

$$\left[\frac{\mu_0 (M_{\text{pl}}^2 / \Lambda^2)}{\mu_0 (M_{\text{p}}^2 / \Lambda^2)} \right]^{a_0/2b_0}$$

$$b_0 = 11 - \frac{2}{3} n_f = 7 \quad \text{for } n_f = 6$$

$$m_p = 1 \text{ GeV} \quad \Lambda = 200 \text{ MeV}$$

$$M_X = 10^{16} \text{ GeV}$$

The enhancement factor is:

$$\left(\frac{\log M_X^2/\Lambda^2}{\log M_p^2/\Lambda^2} \right)^{2/7} = 2.5$$

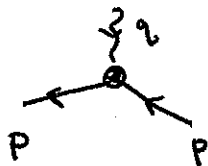
The proton decay rate is increased by the square of this factor

$$\underline{\times 6.}$$

2.) Consider the operator

$$\mathcal{O}_f^{\mu_1 \rightarrow \mu_2} = [\bar{\Psi} \gamma^{\mu_1} (iD)^{\mu_2} - (iD)^{\mu_2} \Psi]_f$$

at the scale $Q \sim m$, it has matrix element



$$= \gamma^{\mu_1} p^{\mu_2} \dots p^{\mu_2}$$

$$\text{for } g \rightarrow 0 \quad \text{or } |q^j| \ll |p^j|$$

the vertex obtains QCD correction from



$$\begin{array}{c} k \\ \swarrow \\ \text{---} \circ \text{---} \\ \searrow \\ k \\ \text{---} \circ \text{---} \\ \swarrow \\ k \\ \searrow \\ k-p \\ \text{---} \circ \text{---} \\ \swarrow \\ p \\ \searrow \\ p \end{array} = (ig)^2 \int \frac{d^4 k}{(2\pi)^4} \frac{-i}{(k-p)^2} \gamma^\lambda t^a \frac{i \not{k}}{k^2} (\gamma^{\mu_1} k^{\mu_2} - k^{\mu_1} \gamma^{\mu_2}) \frac{i \not{k}}{k^2} \gamma_\lambda t^a$$

combine denominators using

$$\left(\frac{1}{k^2}\right)^2 \frac{1}{(k-p)^2} = \int_0^1 dx \frac{2(1-x)}{[k^2 + x(1-x)p^2]^3}$$

with

$$k = k - xp$$

$$k = k + xp$$

We need to find a divergent term with the structure $\gamma^{\mu_1} p^{\mu_2} - p^{\mu_1} \gamma^{\mu_2}$.

To do this, set

$$k \rightarrow k$$

$$k^{\mu_i} \rightarrow x p^{\mu_i}$$

then

$$= -i g^2 \underbrace{(t^a t^a)}_{4/3} \int_0^1 dx \frac{2(1-x)}{[k^2 + x(1-x)p^2]^3} \int \frac{d^4 k}{(2\pi)^4} \frac{1}{[k^2 + x(1-x)p^2]^3}$$

$$\cdot \gamma^\lambda t^a \not{k} \gamma^{\mu_1} \not{k} \gamma^{\mu_2} \not{k} \gamma_\lambda t^a \cdot x^{n-1} p^{\mu_2} - p^{\mu_1}$$

$$= -i g^2 \cdot \frac{4}{3} \int_0^1 dx \frac{2(1-x) x^{n-1}}{(4\pi)^{d/2}} \Gamma\left(\frac{2-d/2}{2}\right) \cdot \frac{1}{2} \gamma^\lambda \gamma^{\mu_1} \gamma^{\mu_2} \gamma_\lambda$$


$$\cdot p^{\mu_2} - p^{\mu_1}$$

Evaluate the residue at $\epsilon=0 \approx d=4$

$$\gamma^\lambda \gamma^\alpha \gamma^{\mu_1} \gamma_\alpha \gamma_\lambda = (-2)^2 \gamma^{\mu_1}$$

$$= \frac{4}{3} \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \gamma^{\mu_1} p^{\mu_2} \dots p^{\mu_n} \cdot \int_0^1 dx \frac{2(1-x) x^{n-1}}{\downarrow}$$

$$2 \cdot \left(\frac{1}{n} - \frac{1}{n+1} \right) = \frac{2}{n(n+1)}$$



$$= \frac{8}{3} \frac{1}{n(n+1)} \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \{\gamma^{\mu_1} p^{\mu_2} \dots p^{\mu_n}\}$$

For the other diagram, we

$$iD^\mu = i\partial^\mu + g A^{\mu a} t^a$$



$$= \sum_{j=2}^n (ig) \cdot g \int \frac{d^d k}{(2\pi)^d} \gamma^\lambda t^a \frac{i k}{k^2} \gamma^{\mu_1}$$

$$\cdot k^{\mu_2} \dots k^{\mu_{j-1}} t^a g_a^{\mu_j} p^{\mu_{j+1}} \dots p^{\mu_n}$$

$$\cdot \frac{-i}{(k-p)^2}$$

$$= \sum_{j=2}^n ig^2 \cdot (t^a t^a) \int_0^1 dx \int \frac{d^d k}{(2\pi)^d} \gamma^{\mu_j} \cancel{x p} \gamma^{\mu_1} \frac{1}{(k^2 - \Delta)^2}$$

$$\cdot \underbrace{(x p)^{\mu_2} \dots (x p)^{\mu_{j-1}} p^{\mu_{j+1}} \dots p^{\mu_n}}_{(j-2 \text{ x's.})}$$


(j-2 x's.)

now

$$\begin{aligned} \gamma^{\mu_j} \not{p} \gamma^{\mu_1} &= \frac{1}{2} [\gamma^{\mu_j} \not{p} \gamma^{\mu_1} + (\mu_1 \leftrightarrow \mu_j)] \\ &= \frac{1}{2} [2 p^{\mu_j} \gamma^{\mu_1} - \gamma^{\mu_j} \gamma^{\mu_1} + \mu_1 \leftrightarrow \mu_j] \\ &= 2 \gamma^{\mu_1} p^{\mu_j} - (\text{term} \propto g^{\mu_1 \mu_j}) \end{aligned}$$

$$= \sum_{j=2}^n i g^2 \cdot \frac{4}{3} \int_0^1 dx \frac{i}{(4\pi)^{d/2}} \Gamma(2-d/2) x^{j-1} 2 \gamma^{\mu_1} p^{\mu_2} - p^{\mu_n}$$

$$= - \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \cdot \sum_{j=2}^n \frac{2}{d} \cdot \gamma^{\mu_1} p^{\mu_2} - p^{\mu_n}$$



$$= \sum_{j=2}^n (i g) \cdot g \cdot \int \frac{d^d k}{(2\pi)^d} \gamma^{\mu_1} p^{\mu_2} \dots p^{\mu_{j-1}} g^{\mu_j \lambda} t^a$$

$$\cdot k^{\mu_{j+1}} - k^{\mu_n} \frac{i k}{k^2} \gamma_{\lambda} t^a \frac{-i}{(k-p)^2}$$

$$= \sum_{j=2}^n i g^2 \cdot \frac{4}{3} \int_0^1 dx \int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 - \Delta)^2}$$

$$\cdot \gamma^{\mu_1} p^{\mu_2} - p^{\mu_{j-1}} p^{\mu_{j+1}} - p^{\mu_n} \not{p} \gamma^{\mu_j} x^{n-j+1}$$

$$l = n - j + 2$$

$$= \sum_{l=2}^2 i g^2 \cdot \frac{4}{3} \left(\frac{i}{(4\pi)^2} \frac{2}{\epsilon} \right) \int_0^1 dx x^{l-1}$$

$$\cdot 2 \gamma^{\mu_1} p^{\mu_2} - p^{\mu_n}$$

$$= (\text{same result as above})$$

so

$$= \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \cdot \left(-\frac{4}{3}\right) \cdot \left[\sum_{j=2}^n \frac{4}{j} - \frac{2}{n(n+1)} \right]$$

$$\delta_{\mathcal{O}_f}^n = \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \frac{4}{3} \sum_{j=2}^n \frac{4}{j} - \frac{2}{n(n+1)}$$

$$\gamma_f^m = M \frac{\partial}{\partial M} (-\delta_{\mathcal{O}_f}^n + \delta_2) \quad \delta_2 = -\frac{4}{3} \frac{g^2}{(4\pi)^2} \frac{2}{\epsilon}$$

$$= M \frac{\partial}{\partial M} \left(-\frac{g^2}{(4\pi)^2} \frac{2}{\epsilon} \cdot \frac{4}{3} \left[\left(\sum_{j=2}^n \frac{4}{j} - \frac{2}{n(n+1)} \right) + 1 \right] \right)$$

$$M \frac{\partial}{\partial M} \left(\frac{2}{\epsilon} \right) = -2$$

$$\gamma_f^m = \frac{8}{3} \frac{g^2}{(4\pi)^2} \left[1 + \sum_{j=2}^n \frac{4}{j} - \frac{2}{n(n+1)} \right]$$

3.) In the problem set, I asked you to work out the Altarelli-Parisi evolution results from

$$\frac{d}{d\ln Q} f(x, Q) = \frac{\alpha_s(Q)}{\pi} \int_x^1 \frac{dz}{z} P_{\frac{q}{q}}(z) f\left(\frac{x}{z}\right)$$

This is the evolution of the non-singlet part of the parton distribution function. For example $(f_u - f_d)$ or $(f_u - f_{\bar{u}})$ obey this equation.

Integrate $\int_0^1 dx x^{n-1}$. Then

$$\frac{d}{d\ln Q} M_n = \frac{\alpha_s(Q)}{\pi} \int_0^1 dx x^{n-1} \int_x^1 \frac{dz}{z} P_{\frac{q}{q}}(z) f\left(\frac{x}{z}\right)$$

$$\text{let } y = \frac{x}{z} \quad \frac{dx}{z} = dy \quad x = yz$$

$$= \frac{\alpha_s(Q)}{\pi} \int \int dy dz (yz)^{n-1} P_{\frac{q}{q}}(z) f(y)$$

The region of integral is

$$0 < x < 1$$

$$x < z < 1 \quad \text{or} \quad y < 1$$

so the integral covers

$$0 < y < 1$$

$$0 < z < 1$$

then

$$\begin{aligned} \frac{d}{d\log Q} M_n &= \frac{\alpha_5(Q)}{\pi} \left[\int_0^1 dz z^{n-1} P_{\frac{2}{3} \log Q}(z) \right] \left[\int_0^1 dy y^{n-1} f(y) \right] \\ &= \frac{\alpha_5(Q)}{\pi} \cdot P_n \cdot M_n \end{aligned}$$

where

$$\begin{aligned} P_n &= \int_0^1 dz z^{n-1} P_{\frac{2}{3} \log Q}(z) \\ &= \int_0^1 dz z^{n-1} \frac{4}{3} \left[\frac{1+z^2}{(1-z)_+} + \frac{3}{2} \delta(z-1) \right] \\ &= \frac{4}{3} \cdot \left\{ \frac{3}{2} + \int_0^1 dz \frac{z^{n-1}(1+z^2) - 2}{(1-z)_+} \right\} \\ &= \frac{4}{3} \cdot \left\{ \frac{3}{2} + \int_0^1 dz \frac{(z^{n-1}-1) + (z^{n+1}-1)}{(1-z)} \right\} \\ &= \frac{4}{3} \cdot \left\{ \frac{3}{2} + \int_0^1 dz (-1) \left[z^{n-2} + z^{n-3} + \dots + 1 \right. \right. \\ &\quad \left. \left. + z^n + z^{n-1} + z^{n-2} + \dots + 1 \right] \right\} \\ &= \frac{4}{3} \cdot \left\{ \frac{3}{2} - \frac{1}{n+1} - \frac{1}{n} - 2 \left(\frac{1}{n-1} + \dots + 1 \right) \right\} \\ &= \frac{4}{3} \cdot \left\{ -\frac{1}{2} - 2 \sum_2^n \frac{1}{j} + \frac{1}{n} - \frac{1}{n+1} \right\} \\ &= -\frac{2}{3} \cdot \left\{ 1 + 4 \sum_2^n \frac{1}{j} - \frac{2}{n(n+1)} \right\} \end{aligned}$$

so

$$\begin{aligned} \frac{d}{dy} M_n &= - \frac{g_s(\alpha)}{\pi} \cdot \frac{2}{3} \cdot \left\{ 1 + 4 \sum_2^n \frac{1}{j} - \frac{2}{n(n+1)} \right\} M_n \\ &= - \frac{8}{3} \frac{g^2}{(4\pi)^2} \left[1 + 4 \sum_2^n \frac{1}{j} - \frac{2}{n(n+1)} \right] M_n \end{aligned}$$

or

$$\frac{d}{dy} M_n = - \gamma_f^n M_n$$