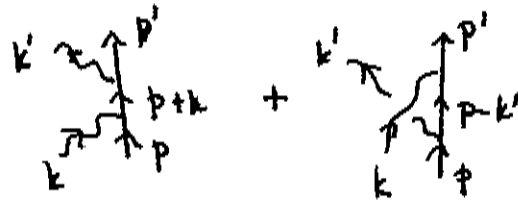


# Physics 330 - Problem Set # 7

## Solutions

1.)



$$iM = (-ie)^2 \bar{u}(p') \left\{ \gamma^\mu \Sigma_\mu^*(k') \frac{i(\not{p} + \not{k} + m)}{(p+k)^2 - m^2} \gamma^\nu \Sigma_\nu(k) \right. \\ \left. + \gamma^\nu \Sigma_\nu(k) \frac{i(\not{p} - \not{k}' + m)}{(p-k')^2 - m^2} \gamma^\mu \Sigma_\mu^*(k') \right\} u(p)$$

$$\text{now } (\not{p} + m) \gamma^\nu u(p) = [2p^\nu - \gamma^\nu (\not{p} - m)] u(p) \\ = 2p^\nu u(p)$$

$$(p+k)^2 - m^2 = m^2 + 2p \cdot k + k^2 - m^2 = 2p \cdot k$$

$$(p-k')^2 - m^2 = m^2 - 2p \cdot k' + k'^2 - m^2 = -2p \cdot k'$$

kinematic relations are:

$$s = (p+k)^2 = (p+k')^2 \Rightarrow p^2 + 2p \cdot k = p^2 + 2p \cdot k'$$

$$t = (p-p')^2 = (k-k')^2 \Rightarrow 2m^2 - 2p \cdot p' = -2kk'$$

$$u = (p-k')^2 = (k-p')^2 \Rightarrow m^2 - 2p \cdot k' = m^2 - 2k \cdot p$$

$$\text{so } p \cdot k = p' \cdot k' \quad p \cdot k' = p' \cdot k \quad p \cdot p' = m^2 + k \cdot k'$$

$$p \cdot (k' + p' - k) = p^2 = m^2 \Rightarrow p \cdot p' = m^2 + p \cdot k - p \cdot k'$$

I will try to systematically eliminate  $\not{p}'k'$ ,  $p'k$ ,  $k \cdot k'$   
in favor of  $p \cdot k$ ,  $p \cdot k'$ ,  $p \cdot p'$

2

$$M = -ie^2 \sum_{\lambda} \epsilon_{\lambda}^{\dagger}(k) \epsilon_{\lambda}(k) \cdot \bar{u}(p') \left[ \frac{2\gamma^{\mu} \not{p}^{\nu} + \gamma^{\mu} \not{k} \gamma^{\nu}}{2p \cdot k} - \frac{2\gamma^{\nu} \not{p}^{\mu} - \gamma^{\nu} \not{k}' \gamma^{\mu}}{2p \cdot k'} \right] u(p)$$

then

$$\begin{aligned} \frac{1}{4} \sum_{\text{Spin}} |M|^2 &= \frac{e^4}{4} \text{tr} \left[ (\not{p}' + m) \left\{ \frac{2\gamma^{\mu} \not{p}^{\nu} + \gamma^{\mu} \not{k} \gamma^{\nu}}{2p \cdot k} + \frac{-2\gamma^{\nu} \not{p}^{\mu} + \gamma^{\nu} \not{k}' \gamma^{\mu}}{2p \cdot k'} \right\} \right. \\ &\quad \left. \cdot (\not{p} + m) \left\{ 2\frac{\gamma_{\nu} \not{p}^{\nu} + \gamma_{\nu} \not{k} \gamma_{\nu}}{2p \cdot k} + \frac{-2\gamma_{\nu} \not{p}^{\mu} + \gamma_{\nu} \not{k}' \gamma_{\nu}}{2p \cdot k'} \right\} \right] \\ &= \frac{e^4}{4} \left[ \frac{\text{I}}{(2p \cdot k)^2} + \frac{\text{II}}{2p \cdot k \cdot 2p \cdot k'} + \frac{\text{III}}{2p \cdot k' \cdot 2p \cdot k} + \frac{\text{IV}}{(2p \cdot k')^2} \right] \end{aligned}$$

now compute I, II, III, IV in turn.

$$\begin{aligned} \text{I} &= \text{tr} (\not{p}' + m) (2\gamma^{\mu} \not{p}^{\nu} + \gamma^{\mu} \not{k} \gamma^{\nu}) (\not{p} + m) (2\gamma_{\nu} \not{p}^{\nu} + \gamma_{\nu} \not{k} \gamma_{\nu}) \\ &= \text{tr} (\not{p}' + m) 4\gamma^{\mu} (\not{p} + m) \gamma_{\mu} p^2 \\ &\quad + \text{tr} (\not{p}' + m) \cdot 2\gamma^{\mu} (\not{p} + m) \not{p} \not{k} \gamma_{\nu} \\ &\quad + \text{tr} (\not{p}' + m) \gamma^{\mu} \not{k} \not{p} (\not{p} + m) \gamma_{\nu} \cdot 2 \\ &\quad + \text{tr} (\not{p}' + m) \gamma^{\mu} \not{k} \gamma^{\nu} (\not{p} + m) \gamma_{\nu} \not{k} \gamma_{\mu} \end{aligned}$$

Here

$$\begin{aligned}
 \text{line 1} &= 4m^2 \text{tr} (\not{p} + m) \gamma^\mu (\not{p} + m) \gamma_\mu \\
 &= 4m^2 \text{tr} (\not{p} + m) (-2\not{p} + 4m) \\
 &= 4 \cdot (-2) \cdot 4 m^2 p \cdot p + 4 \cdot 4 \cdot 4 m^4 \\
 &= -32 p \cdot p m^2 + 64 m^4
 \end{aligned}$$

$$\begin{aligned}
 \text{line 2} &= 2 \text{tr} (\not{p} + m) \gamma^\mu (\not{p} + m) \not{k} \gamma_\mu \\
 &= 2 \text{tr} (\not{p} + m) \gamma^\mu (m^2 + m\not{p}) \not{k} \gamma_\mu \\
 &= 2 \text{tr} (\not{p} + m) (-2\not{k} m^2 + 4 p \cdot k m) \\
 &= -4 \cdot 4 m^2 (p \cdot k) + 2 \cdot 4 \cdot 4 m^2 p \cdot k \\
 &= -16 m^2 (p \cdot k) + 32 m^2 p \cdot k
 \end{aligned}$$

$$\begin{aligned}
 \text{line 3} &= 2 \text{tr} (\not{p} + m) \gamma^\mu \not{k} \not{p} (\not{p} + m) \gamma_\mu \\
 &= 2 \text{tr} (\not{p} + m) \gamma^\mu \not{k} (m^2 + m\not{p}) \gamma_\mu \\
 &= 2 \text{tr} (\not{p} + m) (-2\not{k} m^2 + 4 p \cdot k m) \\
 &= -4 \cdot 4 (p \cdot k) m^2 + 2 \cdot 4 \cdot 4 m^2 p \cdot k \\
 &= -16 p \cdot k m^2 + 32 m^2 p \cdot k
 \end{aligned}$$

I have used my tricks:  $\text{tr} [\text{odd} \# \gamma^a] = 0$

$$\begin{aligned}
\text{line 4} &= \text{tr}(\not{p}' + m) \gamma^\mu \not{k} \gamma^\nu (\not{p} + m) \gamma_\nu \not{k} \gamma_\mu \\
&= \text{tr} \gamma_\mu (\not{p}' + m) \gamma^\mu \not{k} \gamma^\nu (\not{p} + m) \gamma_\nu \not{k} \\
&= \text{tr}(-2\not{p}' + 4m) \not{k} (-2\not{p} + 4m) \not{k} \\
&= 4 \cdot (-2)^2 [p'_i k_i p \cdot k \cdot 2 - \underbrace{p' \cdot p}_{=0} \underbrace{k^2}_{=0}] + 4 \cdot 4 \cdot 4m^2 \underbrace{k^2}_{=0} \\
&= 32 p \cdot k' p \cdot k
\end{aligned}$$

in all

$$\begin{aligned}
\text{I} &= -32 p' p m^2 + 64m^4 + 2(-16m^2 p \cdot k' + 32m^2 p \cdot k) \\
&\quad + 32 p \cdot k p \cdot k' \\
&= -32m^2 (m^2 + p \cdot k - p \cdot k') + 64m^4 - 32 p \cdot k' m^2 + 64m^2 p \cdot k \\
&\quad + 32 p \cdot k p \cdot k' \\
&= 32m^4 + 32m^2 p \cdot k + 32 p \cdot k p \cdot k'
\end{aligned}$$

In eq. (5.84) of P+S it is claimed:

$$\begin{aligned}
\text{I} &= 16(2m^4 + m^2(s - m^2) - \frac{1}{2}(s - m^2)(u - m^2)) \\
&= 32m^4 + 16m^2 \cdot 2p \cdot k - \frac{1}{2} \cdot 16 \cdot (2p \cdot k)(-2p \cdot k') \\
&= 32m^4 + 32m^2 p \cdot k + 32 p \cdot k p \cdot k' \quad \checkmark
\end{aligned}$$

$$\underline{IV} = \text{tr}(\not{p}'+m)[2\gamma^\nu \not{p}' - \gamma^\nu \not{k}' \gamma^\mu](\not{p}'+m)[2\gamma_\nu \not{p}_\mu - \gamma_\nu \not{k}' \gamma_\nu]$$

which is I with  $\mu \leftrightarrow \nu$ ,  $k \leftrightarrow -k'$   
so we can write the answer.

$$\underline{IV} = 32 [m^4 - m^2 p \cdot k' + p \cdot k p \cdot k']$$

$$\underline{II} = \text{tr}(\not{p}'+m)(2\gamma^\mu \not{p}' + \gamma^\mu \not{k}' \gamma^\nu)(\not{p}'+m)(-2\gamma_\nu \not{p}_\mu + \gamma_\nu \not{k}' \gamma_\nu)$$

$$\underline{III} = \text{tr}(\not{p}'+m)(-2\gamma_\nu \not{p}_\mu + \gamma_\nu \not{k}' \gamma_\nu)(\not{p}'+m)(2\gamma^\mu \not{p}' + \gamma^\nu \not{k}' \gamma^\mu)$$

$$= \text{tr}(-2\gamma_\nu \not{p}_\mu + \gamma_\nu \not{k}' \gamma_\nu)(\not{p}'+m)(-2\gamma_\nu \not{p}_\mu + \gamma_\nu \not{k}' \gamma_\nu)(\not{p}'+m)$$

= II w. order of  $\gamma$ 's in the trace reversed.

$$= \underline{II} \quad \text{by eq. (5.7)}$$

$$\underline{II} = \text{tr}(\not{p}'+m) 2\gamma^\mu \not{p}' (\not{p}'+m)(-2\gamma_\nu \not{p}_\mu)$$

$$+ \text{tr}(\not{p}'+m) 2\gamma^\mu \not{p}' (\not{p}'+m)(\gamma_\mu \not{k}' \gamma_\nu)$$

$$+ \text{tr}(\not{p}'+m) \gamma^\mu \not{k}' \gamma^\nu (\not{p}'+m)(-2\gamma_\nu \not{p}_\mu)$$

$$+ \text{tr}(\not{p}'+m) \gamma^\mu \not{k}' \gamma^\nu (\not{p}'+m) \gamma_\mu \not{k}' \gamma_\nu$$

$$\text{line 1} = -4 \text{tr}(\not{p}'+m) \not{p}' (\not{p}'+m) \not{p}'$$

$$= -4 \text{tr}(\not{p}'+m) (m^2 \not{p}' + m^3)$$

$$= -4 \cdot 4 (p \cdot p' m^2 + m^4)$$

$$= -16 (p \cdot p' m^2 + m^4)$$

$$\begin{aligned}
 \text{line 2} &= 2 \text{tr}(\not{p}'+m)\gamma^\mu(\not{p}+m)\gamma_\mu k' \not{p} \\
 &= 2 \text{tr}(\not{p}'+m)(-2\not{p}+4m) \not{k}' \not{p} \\
 &= -4 \cdot 4 [\text{tr}(\not{p}'\not{p} \not{k}'\not{p} \cdot 2 - \text{tr}(\not{p}'\not{k}'m^2)] + 2 \cdot 4 \cdot 4m^2 \text{tr}(\not{k}'\not{p}) \\
 &= -32 \text{tr}(\not{p}'\not{p} \not{k}'\not{p}) + 16m^2 \text{tr}(\not{p}'\not{k}') + 32m^2 \text{tr}(\not{p}'\not{k}')
 \end{aligned}$$

$$\begin{aligned}
 \text{line 3} &= -2 \text{tr}(\not{p}'+m) \not{p} \not{k}' \gamma^\nu(\not{p}+m)\gamma_\nu \\
 &= -2 \text{tr}(\not{p}'+m) \not{p} \not{k}' (-2\not{p}+4m) \\
 &= 4 \cdot 4 [\text{tr}(\not{p}'\not{p} \not{k}'\not{p} \cdot 2 - \text{tr}(\not{p}'\not{k}'m^2)] - 2 \cdot 4 \cdot 4m^2 \text{tr}(\not{p}'\not{k}') \\
 &= 32 \text{tr}(\not{p}'\not{p} \not{k}'\not{p}) - 16m^2 \text{tr}(\not{p}'\not{k}') - 32m^2 \text{tr}(\not{p}'\not{k}')
 \end{aligned}$$

$$\begin{aligned}
 \text{line 4} &= \text{tr}(\not{p}'+m)\gamma^\mu \not{k}' \gamma^\nu(\not{p}+m)\gamma_\mu k' \gamma_\nu \\
 &= \text{tr}(\not{p}'+m)\gamma^\mu \not{k}' (-2k' \gamma_\mu \not{p} + 4m k' \gamma_\mu) \\
 &= -2 \cdot 4 k \cdot k' \text{tr}(\not{p}'\not{p}) + 4m^2 \text{tr}(\not{k}'\not{k}') \\
 &= -8 k \cdot k' \cdot 4 \text{tr}(\not{p}'\not{p}) + 4 \cdot 4m^2 k \cdot k' \\
 &= -32 k \cdot k' \text{tr}(\not{p}'\not{p}) + 16m^2 k \cdot k'
 \end{aligned}$$

~ all

$$\begin{aligned}
 \text{III} &= -16m^4 - 16 \text{tr}(\not{p}'\not{p})m^2 - 32 \text{tr}(\not{p}'\not{p} \not{p}'\not{k}') + 16m^2 \text{tr}(\not{p}'\not{k}') + 32m^2 \text{tr}(\not{p}'\not{k}') \\
 &\quad + 32 \text{tr}(\not{p}'\not{p} \not{k}'\not{p}) - 16m^2 \text{tr}(\not{p}'\not{k}') - 32m^2 \text{tr}(\not{p}'\not{k}') \\
 &\quad - 32 \text{tr}(\not{p}'\not{p}' \not{k}'\not{k}') + 16m^2 k \cdot k' \\
 &= 16 \text{tr}(\not{p}'\not{p}) \{ -m^2 - 2 \text{tr}(\not{p}'\not{k}') + 2 \text{tr}(\not{p}'\not{k}') - 2 k \cdot k' \} \\
 &\quad + 16 \{ -m^4 + m^2 \text{tr}(\not{p}'\not{k}') + m^2 \text{tr}(\not{p}'\not{k}') + m^2 k \cdot k' \}
 \end{aligned}$$

$$\begin{aligned}
 &= 16 p \cdot p' \left\{ -m^2 - 2 \overbrace{(p+k)}^{p+k'} \cdot k' + 2 p \cdot k \right\} \\
 &\quad + 16 \left\{ -m^4 - m^2 p \cdot k + m^2 p \cdot k' + m^2 (p \cdot p' - m^2) \right\} \\
 &= 16 p \cdot p' \left\{ -m^2 - 2 \cancel{p} \cdot \cancel{k}' + 2 \cancel{p} \cdot k \right\} + 16 m^2 p \cdot p' \\
 &\quad - 32 m^4 - 16 m^2 p \cdot k + 16 m^2 p \cdot k'
 \end{aligned}$$

$$\text{II} = -32 m^4 - 16 m^2 p \cdot k + 16 m^2 p \cdot k'$$

this does agree with  $\text{III} = -8(4m^4 + m^2(s-m^2) + m^2(u-m^2))$

in all

$$\frac{1}{4} \sum_{\text{spins}} |M|^2 = \frac{e^4}{4} \left[ \frac{32}{4} \right]$$

$$\cdot \left\{ \frac{1}{(p \cdot k)^2} (p \cdot k p \cdot k' + m^2 p \cdot k + m^4) \right.$$

$$\left. + \frac{1}{(p \cdot k')^2} (p \cdot k p \cdot k' - m^2 p \cdot k' + m^4) \right.$$

$$\left. + \frac{1}{(p \cdot k)(p \cdot k')} (-2m^4 - m^2 p \cdot k + m^2 p \cdot k') \right\}$$

$$= 2e^4 \left\{ \frac{p \cdot k'}{p \cdot k} + \frac{p \cdot k}{p \cdot k'} + 2m^2 \left( \frac{1}{p \cdot k} - \frac{1}{p \cdot k'} \right) \right.$$

$$\left. + m^4 \left( \frac{1}{(p \cdot k)} - \frac{1}{(p \cdot k')} \right)^2 \right\}$$

= the CM frame

8

$$p = (E, 0, 0, p)$$

$$k = (p, 0, 0, -p)$$

$$p' = (E, p \sin \theta, 0, p \cos \theta)$$

$$k' = (p, -p \sin \theta, 0, -p \cos \theta)$$

$$p \cdot k = p(E+p)$$

$$E+p = E_{cm}$$

$$p \cdot k' = p(E+p \cos \theta)$$

$$p = \frac{s - m_e^2}{2\sqrt{s}}$$

$$\frac{1}{p \cdot k} - \frac{1}{p \cdot k'} = \frac{1}{p} \left( \frac{1}{E+p} - \frac{1}{E+p \cos \theta} \right)$$

$$= -\frac{1}{p} \frac{p \cos \theta}{(E+p)(E+p \cos \theta)}$$

$$= -\frac{\cos \theta}{(E+p)(E+p \cos \theta)}$$

so

$$\frac{1}{4} \sum |M|^2 = 2e^4 \left\{ \left( \frac{E+p \cos \theta}{E+p} + \frac{E+p}{E+p \cos \theta} \right) - 2m^2 \frac{\cos \theta}{(E+p)(E+p \cos \theta)} + \frac{m^4 \cos^2 \theta}{(E+p)^2 (E+p \cos \theta)^2} \right\}$$

$$d\pi = \frac{1}{16\pi} d\cos \theta \cdot \left( \frac{2p}{E+p} \right)$$

then

$$\frac{d\sigma}{d\cos\theta} = \frac{1}{4E_p \left(\frac{p}{E} + 1\right)} \frac{1}{16\pi} \frac{2p}{(E+p)} 2e^4 \{ \dots \}$$

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{E_{cm}^2} \left\{ \frac{2E^2 + 2Ep(1+\cos\theta) + p^2(1+\cos^2\theta)}{(E+p)(E+p\cos\theta)} - \frac{2m^2 \cos\theta}{(E+p)(E+p\cos\theta)} + \frac{m^4 \cos^2\theta}{(E+p)^2 (E+p\cos\theta)^2} \right\}$$

for  $E \gg m$ 

$$\frac{d\sigma}{d\cos\theta} \sim \frac{\pi\alpha^2}{4E^2} \cdot \frac{5 + 2\cos\theta + \cos^2\theta}{2(1+\cos\theta)}$$

2.) a) In the limit of  $m_e \rightarrow 0$

$$iM = -ie^2 \bar{u}(p') \left\{ \gamma \cdot \epsilon^+(k') \frac{\not{p} + \not{k}}{2pk} \gamma \cdot \epsilon(k) \right. \\ \left. + \gamma \cdot \epsilon(k) \frac{\not{p} - \not{k}'}{-2pk'} \gamma \cdot \epsilon^+(k') \right\} u(p)$$

Note that

$$\bar{u}(p') \gamma^\mu \gamma^\nu \gamma^\lambda u(p)$$

$$= u^\dagger(p') \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^x \\ \sigma^x & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^y \\ \sigma^y & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^z \\ \sigma^z & 0 \end{pmatrix} u(p)$$

$$= u^\dagger(p') \begin{bmatrix} \sigma^x \sigma^y \sigma^z & 0 \\ 0 & \sigma^x \sigma^y \sigma^z \end{bmatrix} u(p)$$

so left-handed spinors  $u_L(p) = \begin{pmatrix} \mathbb{X} \\ 0 \end{pmatrix}$  and right-handed

spinors  $u_R(p) = \begin{pmatrix} 0 \\ \mathbb{X} \end{pmatrix}$  do not mix.

so the possible nonzero helicity amplitudes are

$$e^- \gamma_L \rightarrow e^- \gamma_L$$

$$e^- \gamma_L \rightarrow e^- \gamma_R$$

$$e^- \gamma_L \rightarrow e^- \gamma_R$$

$$e^- \gamma_L \rightarrow e^- \gamma_L$$

$$e^- \gamma_R \rightarrow e^- \gamma_L$$

$$e^- \gamma_R \rightarrow e^- \gamma_R$$

$$e^- \gamma_R \rightarrow e^- \gamma_R$$

$$e^- \gamma_R \rightarrow e^- \gamma_L$$

first compute  $e_L \gamma \rightarrow e_L \gamma$  amplitudes. Then

$$\mathcal{M} = -ie^2 \sqrt{2E} (-s_2 c_2) \bar{v} \cdot \epsilon^*(k') \frac{\sigma(p+k)}{2p \cdot k} \bar{v} \cdot \epsilon(k) \begin{pmatrix} 0 \\ 1 \end{pmatrix} \sqrt{2E}$$

$$\mathcal{M} = +ie^2 \sqrt{2E} (-s_2 c_2) \bar{v} \cdot \epsilon(k) \frac{\sigma(p-k')}{2p \cdot k'} \bar{v} \cdot \epsilon^*(k') \begin{pmatrix} 0 \\ 1 \end{pmatrix} \sqrt{2E}$$

$s_2 = s \sin \Theta/2 \quad c_2 = c \cos \Theta/2$

now  $2p \cdot k = 4E^2$        $2p \cdot k' = 2E^2(1 + \cos \Theta)$   
 so  $s = \sin \Theta, \quad c = \cos \Theta$

$$\mathcal{M} = -ie^2 \frac{1}{2E} (-s_2 c_2) \bar{v} \cdot \epsilon^*(k') \sigma(p+k) \bar{v} \cdot \epsilon(k) \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\mathcal{M} = +ie^2 \frac{1}{E(1+c)} (-s_2 c_2) \bar{v} \cdot \epsilon(k) \sigma(p-k') \bar{v} \cdot \epsilon^*(k') \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$p+k = (2E, \vec{0})$  so  $\sigma(p+k) = 2E \cdot \underline{1}$

$p-k' = (E \ 0 \ 0 \ E) - (E, -Es, 0, -Ec) = (0, +Es, 0, E(1+c))$

so  $\sigma(p-k') = -E \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix}$

so

$$\mathcal{M} = -ie^2 (-s_2 c_2) \bar{v} \cdot \epsilon^*(k') \underline{1} \bar{v} \cdot \epsilon(k) \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\mathcal{M} = \frac{-ie^2}{(1+c)} (-s_2 c_2) \bar{v} \cdot \epsilon(k) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \bar{v} \cdot \epsilon^*(k') \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

for all we need

$$\bar{\sigma} \cdot \epsilon(k) = L: \frac{1}{\sqrt{2}} (-\sigma^1 - i\sigma^2) = \frac{1}{\sqrt{2}} \left[ \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \right] = -\sqrt{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$R: \frac{1}{\sqrt{2}} (-\sigma^1 + i\sigma^2) = \frac{1}{\sqrt{2}} \left[ \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} + i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \right] = -\sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$\bar{\sigma} \cdot \epsilon^*(k) = L: \frac{1}{\sqrt{2}} \left( \begin{pmatrix} 0 & -c \\ -c & 0 \end{pmatrix} + i \begin{pmatrix} 0 & -c \\ i & 0 \end{pmatrix} + s \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \right) = \frac{1}{\sqrt{2}} \begin{pmatrix} s & 1-c \\ -(1+c) & -s \end{pmatrix}$$

$$R: \frac{1}{\sqrt{2}} \left( \begin{pmatrix} 0 & -c \\ -c & 0 \end{pmatrix} - i \begin{pmatrix} 0 & -c \\ i & 0 \end{pmatrix} + s \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \right) = \frac{1}{\sqrt{2}} \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix}$$

again:

$\bar{\sigma} \cdot \epsilon_L(k)$	$\bar{\sigma} \cdot \epsilon_R(k)$	$\bar{\sigma} \cdot \epsilon_L^*(k')$	$\bar{\sigma} \cdot \epsilon_R^*(k')$
$-\sqrt{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$	$-\sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$	$\frac{1}{\sqrt{2}} \begin{pmatrix} s & 1-c \\ -(1+c) & -s \end{pmatrix}$	$\frac{1}{\sqrt{2}} \begin{pmatrix} s & -(1+c) \\ 1-c & -s \end{pmatrix}$

Begin with  $\}$  note that  $\bar{\sigma} \cdot \epsilon_R(k) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 0$  so

$$\mathcal{M}(\bar{e}_L \gamma_R \rightarrow \bar{e}_L \gamma_{L \text{ or } R}) = 0$$

$$\begin{aligned} i\mathcal{M}(\bar{e}_L \gamma_L \rightarrow \bar{e}_L \gamma_L) &= -ie^2 (-s_2 c_2) \frac{1}{\sqrt{2}} \begin{pmatrix} s & (1-c) \\ -(1+c) & -s \end{pmatrix} (-\sqrt{2}) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ &= ie^2 (-s_2 c_2) \begin{pmatrix} s & (1-c) \\ -(1+c) & -s \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ &= ie^2 [-s_2 s - c_2 (1+c)] \\ &= +ie^2 [-\cos\theta_L - (\cos\theta \cos\theta/2 + \sin\theta \sin\theta/2)] \\ &= -2ie^2 \cos\theta/2 \end{aligned}$$

$$\begin{aligned}
 iM(e_L^- \gamma_L \rightarrow e_L^- \gamma_R) &= -ie^2 (-s_2 c_2) \frac{1}{\sqrt{2}} \begin{pmatrix} s & -(1+c) \\ 1-c & -s \end{pmatrix} (-\sqrt{2}) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
 &= ie^2 (-s_2 c_2) \begin{pmatrix} s & -(1+c) \\ 1-c & -s \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
 &= ie^2 (-s_2 s + c_2 - c_2 c) \\
 &= ie^2 (c_2 - (c_2 c + s_2 s)) = 0
 \end{aligned}$$

so for

$$\left. \begin{array}{l} \uparrow \\ \downarrow \end{array} \right\} iM(e_L^- \gamma_L \rightarrow e_L^- \gamma_L) = -2ie^2 c_2$$

all other helicity amplitudes are 0.

for  $\left. \begin{array}{l} \uparrow \\ \downarrow \end{array} \right\}$ :

$$\begin{aligned}
 e_L^- \gamma_L \rightarrow e_L^- \gamma_L : &= \frac{-ie^2}{1+c} (-s_2 c_2) (-\sqrt{2}) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} s & 1-c \\ -(1+c) & -s \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
 &= \frac{+ie^2}{1+c} (0, -s_2) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} 1-c \\ -s \end{pmatrix} \\
 &= \frac{+ie^2}{1+c} [-s_2 s (1-c) + s_2 s (1+c)] \\
 &\rightarrow \frac{-2ie^2}{1+c} s_2 s
 \end{aligned}$$

$$\begin{aligned}
e_L^- \gamma_L &\rightarrow e_L^- \gamma_R : \quad \frac{-ie^2}{1+c} (-s_2 c_2) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} (-\sqrt{2}) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
&= \frac{ie^2}{1+c} (0 -s_2) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} -(1+c) \\ -s \end{pmatrix} \\
&= \frac{ie^2}{1+c} [ +s_2 s(1+c) - s_2 s(1+c) ] = 0
\end{aligned}$$

$$\begin{aligned}
e_L^- \gamma_R &\rightarrow e_L^- \gamma_L \quad \frac{-ie^2}{1+c} (-s_2 c_2) -\sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} s & 1+c \\ -(1+c) & -s \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
&= \frac{+ie^2}{1+c} (c_2, 0) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} 1+c \\ -s \end{pmatrix} \\
&= \frac{+ie^2}{1+c} [ c_2(1+c)(1-c) - c_2 s^2 ] = 0
\end{aligned}$$

$\underbrace{(1+c)(1-c)}_{\sin^2 \theta}$

$$\begin{aligned}
e_L^- \gamma_R &\rightarrow e_L^- \gamma_R \quad \frac{-ie^2}{1+c} (-s_2 c_2) (-\sqrt{2}) \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
&= \frac{+ie^2}{1+c} (+s_2, 0) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} -(1+c) \\ -s \end{pmatrix} \\
&= \frac{ie^2}{1+c} [-c_2(1+c)^2 - c_2 s^2] \\
&= \frac{ie^2}{1+c} [-c_2(1+c) - c_2(1-c)] (1/c) \\
&= -2ie^2 c_2
\end{aligned}$$

now add the contributions

$$\mathcal{H} + \mathcal{H} =$$

$$e_L \gamma_L \rightarrow e_L \gamma_L \quad -2ie^2 \left[ c_2 + \frac{ss_2}{(1+c_2)} \right]$$

$$= -2ie^2 \frac{c_2 + (c_2 c + s_2 s)}{(1+c)}$$

$$= -4ie^2 \frac{c_2}{1+c}$$

$$= -2ie^2 \frac{1}{c_2}$$

$$c c_2 + s s_2 = \cos^2 \theta = c_2$$

$$c_2^2 = \frac{1+c}{2}$$

$$e_L \gamma_L \rightarrow e_L \gamma_R = 0$$

$$e_L \gamma_R \rightarrow e_R \gamma_L = 0$$

$$e_L \gamma_R \rightarrow e_L \gamma_R = 0 - 2ie^2 c_2$$

$$= -2ie^2 c_2$$

The amplitudes for  $e_R \gamma \rightarrow e_R \gamma$  should follow the same pattern. The basic ingredients are:

$$\mathcal{H} = -ie^2 \sqrt{2E} (c_2 s_2) \sigma \cdot \epsilon^*(k') \frac{\bar{\sigma} \cdot p + k}{2p \cdot k} \sigma \cdot \epsilon(k) \begin{pmatrix} 1 \\ 0 \end{pmatrix} (\sqrt{2E})^*$$

$$\mathcal{H} = +ie^2 \sqrt{2E} (c_2 s_2) \sigma \cdot \epsilon(k) \frac{\bar{\sigma} \cdot p - k'}{2p \cdot k'} \sigma \cdot \epsilon^*(k') \begin{pmatrix} 1 \\ 0 \end{pmatrix} \sqrt{2E}$$

now  $\bar{\sigma} \cdot (p+k) = 2E \cdot \underline{1}$        $\bar{\sigma} \cdot (p-k') = +E \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix}$

so

$$\chi = -ie^2 (c_2 s_2) \sigma \cdot \epsilon^*(k') \cdot 1 \cdot \sigma \cdot \epsilon(k) \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\chi = \frac{+ie^2}{1+c} (c_2 s_2) \sigma \cdot \epsilon(k) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \sigma \cdot \epsilon^*(k') \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

Evaluate these using:

$\sigma \cdot \epsilon_L(k)$	$\sigma \cdot \epsilon_R(k)$	$\sigma \cdot \epsilon_L^*(k')$	$\sigma \cdot \epsilon_R^*(k')$
$+\sqrt{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$	$+\sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$	$-\frac{1}{\sqrt{2}} \begin{pmatrix} s & 1-c \\ -(1+c) & -s \end{pmatrix}$	$-\frac{1}{\sqrt{2}} \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix}$

for  $\chi$   $\sigma \cdot \epsilon(k) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0$ , so the possibly nonzero amplitudes are:

$$\begin{aligned} e_R^- \gamma_R \rightarrow e_R^- \gamma_L &: -ie^2 (c_2 s_2) \left(-\frac{1}{\sqrt{2}}\right) \begin{pmatrix} s & 1-c \\ -(1+c) & -s \end{pmatrix} \sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ &= ie^2 (c_2 s_2) \begin{pmatrix} s & 1-c \\ -(1+c) & -s \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ &= ie^2 [c_2(1-c) - s_2 s] = ie^2 [c_2 - (c c_2 + s s_2)] \\ &= 0 \end{aligned}$$

$$\begin{aligned} e_R^- \gamma_R \rightarrow e_R^- \gamma_R &: -ie^2 (c_2 s_2) \left(-\frac{1}{\sqrt{2}}\right) \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix} \sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ &= ie^2 (c_2 s_2) \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ &= ie^2 [-c_2(1+c) - s s_2] = -2ie^2 c_2 \end{aligned}$$

for  $\chi_L$

$$\begin{aligned}
 \bar{e}_R \chi_L &\rightarrow \bar{e}_R \gamma_L : \frac{+ie^2}{1+c} (c_2 s_2) \sqrt{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} -1 \\ \sqrt{2} \end{pmatrix} \begin{pmatrix} s(1-c) \\ -(1+c)s \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
 &= \frac{-ie^2}{1+c} (0, c_2) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} s \\ -(1+c) \end{pmatrix} \\
 &= \frac{ie^2}{1+c} [c_2 s^2 + c_2 (1+c)^2] \\
 &= \frac{-ie^2}{1+c} (1+c) [c_2(1-c) + c_2(1+c)] = 2ie^2 c_2
 \end{aligned}$$

$$\begin{aligned}
 \bar{e}_R \chi_L &\rightarrow \bar{e}_R \gamma_R : \frac{+ie^2}{1+c} (c_2 s_2) \sqrt{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} -1 \\ \sqrt{2} \end{pmatrix} \begin{pmatrix} s & -(1+c) \\ (1-c) & -s \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
 &= \frac{-ie^2}{1+c} (c_2 s_2) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} s \\ (1-c) \end{pmatrix} \\
 &= \frac{-ie^2}{1+c} [c_2 s^2 - c_2 \underbrace{(1+c)(1-c)}_{s^2}] = 0
 \end{aligned}$$

$$\begin{aligned}
 \bar{e}_R \gamma_R &\rightarrow \bar{e}_R \chi_L : \frac{+ie^2}{1+c} (c_2 s_2) \sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} -1 \\ \sqrt{2} \end{pmatrix} \begin{pmatrix} s(1-c) \\ -(1+c)s \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
 &= \frac{-ie^2}{1+c} (s_2, 0) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} s \\ -(1+c) \end{pmatrix} \\
 &= \frac{-ie^2}{1+c} [s_2 s(1+c) - s_2 s(1+c)] = 0
 \end{aligned}$$

$$\begin{aligned}
 \bar{e}_R \gamma_R \rightarrow \bar{e}_R \gamma_R : & \quad \frac{+ie^2}{1+c} (c_2 s_2) \sqrt{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} -1 \\ \sqrt{2} \end{pmatrix} \begin{pmatrix} s-(1+c) \\ (1-c)-s \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
 & = \frac{-ie^2}{1+c} (s_2) \begin{pmatrix} 1+c & s \\ s & -(1+c) \end{pmatrix} \begin{pmatrix} s \\ (1-c) \end{pmatrix} \\
 & = \frac{-ie^2}{1+c} [s_2(1+c)s + s_2 s(1-c)] \\
 & = \frac{-2ie^2}{1+c} s s_2
 \end{aligned}$$

the complete amplitude for  $\bar{e}_R \gamma_R \rightarrow \bar{e}_R \gamma_R$  is:

$$\begin{aligned}
 -2ie^2 c_2 - \frac{2ie^2}{1+c} s s_2 & = \frac{-2ie^2}{1+c} [c_2 + c c_2 + s s_2] \\
 & = -2ie^2 \frac{2c_2}{1+c} = -2ie^2 \frac{1}{c_2}
 \end{aligned}$$

so!

$$i\mathcal{M}(\bar{e}_L \gamma_L \rightarrow \bar{e}_L \gamma_L) = i\mathcal{M}(\bar{e}_R \gamma_R \rightarrow \bar{e}_R \gamma_R) = -2ie^2 \frac{1}{c_2}$$

$$i\mathcal{M}(\bar{e}_L \gamma_R \rightarrow \bar{e}_L \gamma_R) = i\mathcal{M}(\bar{e}_R \gamma_L \rightarrow \bar{e}_R \gamma_L) = -2ie^2 c_2$$

all other helicity amplitudes vanish.

then the polarization summed/averaged cross section is

19

$$\begin{aligned}
 \frac{ds}{d\cos\Theta} &= \frac{1}{2s} \cdot \frac{1}{16\pi} \cdot \frac{1}{4} \sum |M|^2 \\
 &= \frac{1}{2s} \frac{1}{16\pi} \frac{1}{4} \cdot 4e^4 \cdot 2 \cdot \left[ \frac{1}{c_2^2} + c_2^2 \right] \\
 &= \frac{\pi\alpha^2}{E_{cm}^2} \left[ \frac{2}{1+\cos\Theta} + \frac{1+\cos\Theta}{2} \right] \\
 &= \frac{\pi\alpha^2}{4E^2} \frac{4 + (1+\cos\Theta)^2}{2(1+\cos\Theta)}
 \end{aligned}$$

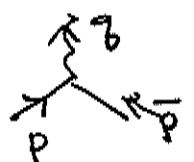
which reproduces the result on p. 9.

3.) The vertex for this problem is

$$i\Gamma = -ig \gamma^\mu$$

I will ignore the electron mass except where it is essential in part (c).

a)  $\sigma(e^+e^- \rightarrow B)$  is given by



The diagram shows an incoming electron line with momentum p and an incoming positron line with momentum p-bar. They meet at a vertex, from which a single outgoing line with momentum q (representing particle B) emerges.

$$= -ig \bar{v}(p) \gamma^\mu u(p) \cdot \Sigma_\mu^*(q)$$

$$\begin{aligned}
\frac{1}{4} \sum |M|^2 &= \frac{g^2}{4} \text{tr} [\bar{\not{p}} \gamma^\mu \not{p} \gamma^\nu] (-g_{\mu\nu}) \\
&= -\frac{g^2}{4} \text{tr} \not{p} (-2\not{p}) \\
&= \frac{g^2}{4} \cdot 2 \cdot 4 \bar{p} \cdot p = \frac{1}{2} 2 \bar{p} \cdot p = g^2 (\bar{p} + p)^2 \\
&= \frac{1}{2} g^2 = M^2 \cdot g^2
\end{aligned}$$

$$\sigma(e^+e^- \rightarrow B) = \frac{1}{2S} \int \frac{d^3q}{(2\pi)^3} \frac{1}{2E_q} (2\pi)^4 \delta(p + \bar{p} - q) \left( \frac{1}{4} |M|^2 \right)$$

$$= \frac{1}{2S} 2\pi \frac{1}{2M} \delta(p^0 + \bar{p}^0 - M) \cdot \frac{1}{2} M^2$$

CM frame.

$$= \frac{1}{2S} \cdot 2\pi \delta((p + \bar{p})^2 - M^2) \frac{1}{2} M^2$$

$$s = M^2 \quad = \quad \pi g^2 \delta(s - M^2)$$

$$\Gamma(B \rightarrow e^+e^-) = \frac{1}{2M} \int d\pi_2 \left( \frac{1}{3} \sum |M|^2 \right)$$

↑ spin average for spin-1

$$\sum |M|^2 \text{ is the same as above } = 4g^2 M^2$$

$$= \frac{1}{2M} \cdot \frac{1}{8\pi} \cdot \frac{4}{3} \cdot g^2 M^2 = \frac{g^2 M}{12\pi}$$

so

$$\sigma(e^+e^- \rightarrow B) = \pi g^2 \delta(s - M^2)$$

$$\Gamma_B = \frac{g^2 M}{12\pi}$$

$$\left. \begin{array}{l} \sigma(e^+e^- \rightarrow B) = \pi g^2 \delta(s - M^2) \\ \Gamma_B = \frac{g^2 M}{12\pi} \end{array} \right\} \sigma(e^+e^- \rightarrow B) = \frac{12\pi^2}{M} \Gamma_B \delta(s - M^2)$$

b.) für  $e^+e^- \rightarrow \gamma + B$ :



$$= (-ig)(-ie)\bar{u}(\bar{p}) \left\{ \gamma \cdot \epsilon^*(q) \frac{i(\not{p}-\not{k})}{(p-k)^2} \gamma \cdot \epsilon^*(k) \right. \\ \left. + \gamma \cdot \epsilon^*(k) \frac{i(\not{k}-\not{p})}{(p-k)^2} \gamma \cdot \epsilon^*(q) \right\} u(p)$$

$$(p-k)^2 = p^2 - 2p \cdot k + k^2 = -2p \cdot k$$

$$(\bar{p}-k)^2 = -2\bar{p} \cdot k$$

$$\frac{1}{4} \sum |M|^2 = \frac{e^2 g^2}{4} \text{tr} \left\{ \not{p} \left[ \frac{\gamma^\mu (\not{p}-\not{k}) \gamma^\nu}{2p \cdot k} + \frac{\gamma^\nu (\not{k}-\not{p}) \gamma^\mu}{2\bar{p} \cdot k} \right] \right. \\ \left. \not{\bar{p}} \left[ \frac{\gamma_\mu (\not{p}-\not{k}) \gamma_\nu}{2p \cdot k} + \frac{\gamma_\nu (\not{k}-\not{p}) \gamma_\mu}{2\bar{p} \cdot k} \right] \right\}$$

$$= \frac{e^2 g^2}{4} \left\{ \frac{\text{I}}{(2p \cdot k)^2} + \frac{\text{II}}{2p \cdot k \cdot 2\bar{p} \cdot k} + \frac{\text{III}}{2\bar{p} \cdot k \cdot 2p \cdot k} + \frac{\text{IV}}{(2\bar{p} \cdot k)^2} \right\}$$

$$\begin{aligned}
 I &= \text{tr } \bar{p} \gamma^\mu (\not{p}-\not{k}) \gamma^\nu \not{k} \gamma_\nu (\not{p}-\not{k}) \gamma_\mu \\
 &= \text{tr } \gamma_\mu \bar{p} \gamma^\nu (\not{p}-\not{k}) \gamma^\mu \not{k} \gamma_\nu (\not{p}-\not{k}) \\
 &= (-2)^2 \text{tr } \bar{p} (\not{p}-\not{k}) \not{k} (\not{p}-\not{k})
 \end{aligned}$$

$$\begin{aligned}
 p^2 = \bar{p}^2 = k^2 = 0 \\
 g^2 = M^2
 \end{aligned}$$

$$\begin{aligned}
 &= 4 \cdot 4 \cdot [ \bar{p} \cdot (\not{p}-\not{k}) \not{k} (\not{p}-\not{k}) \cdot 2 - \bar{p} \cdot \bar{p} (\not{p}-\not{k})^2 ] \\
 &= 16 \cdot [ -2 \bar{p} \cdot p \not{k} + 2 \bar{p} \cdot k \not{k} + 2 \bar{p} \cdot \bar{p} \not{k} ] \\
 &= 32 [ \not{k} ] [ -\cancel{\bar{p} \cdot p} + \bar{p} \cdot k + \cancel{\bar{p} \cdot p} ] \\
 &= 32 \not{k} \bar{p} \cdot k
 \end{aligned}$$

$$\text{IV} = \text{tr } \bar{p} \gamma^\nu (\not{\bar{p}}-\not{k}) \gamma^\mu \not{k} \gamma_\mu (\not{\bar{p}}-\not{k}) \gamma_\nu$$

$$\begin{aligned}
 &= (-2)^2 \text{tr } \bar{p} (\not{\bar{p}}-\not{k}) \not{k} (\not{\bar{p}}-\not{k}) \\
 &= 4 \cdot 4 \cdot \{ \bar{p} \cdot (\not{\bar{p}}-\not{k}) \not{k} (\not{\bar{p}}-\not{k}) \cdot 2 - \bar{p} \cdot \bar{p} (\not{\bar{p}}-\not{k})^2 \} \\
 &= 16 \{ -\bar{p} \cdot k (\cancel{2\bar{p} \cdot \bar{p}} - 2 \not{k}) - (\cancel{\bar{p} \cdot \bar{p}} - 2 \bar{p} \cdot k) \} \\
 &= 32 \bar{p} \cdot k \not{k}
 \end{aligned}$$

$$\text{II} = \text{tr } \bar{p} \gamma^\mu (\not{p}-\not{k}) \gamma^\nu \not{k} \gamma_\mu (\not{\bar{p}}-\not{k}) \cdot (-1) \gamma_\nu$$

$$\begin{aligned}
 &= (-1) \text{tr } \bar{p} \gamma^\mu (\not{p}-\not{k}) (-2) (\not{\bar{p}}-\not{k}) \gamma_\mu \not{k} \\
 &= 2 \text{tr } \bar{p} \not{k} (\not{p}-\not{k}) (\not{\bar{p}}-\not{k}) \not{k} \\
 &= 2 \cdot 4 \cdot 4 (\bar{p} \cdot \bar{p}) (\bar{p} \cdot \bar{p} - \bar{p} \cdot k - \bar{p} \cdot k)
 \end{aligned}$$

$$\text{now } s = 2\bar{p} \cdot \bar{p} \quad t = -2\bar{p} \cdot k \quad u = -2\bar{p} \cdot k$$

$$s+t+u = M^2 \quad \text{so} \quad \bar{p} \cdot \bar{p} - \bar{p} \cdot k - \bar{p} \cdot k = M^2/2$$

$$\text{II} = 16 \bar{p} \cdot \bar{p} M^2$$

$$\begin{aligned} \text{III} &= \text{tr} \not{\bar{p}} \gamma^\nu (\not{\bar{p}} - \not{k}) (-i \gamma^\mu \not{\not{p}} \gamma^\nu (\not{p} - \not{k}) \gamma_\mu) \\ &= \text{tr} (-i)(-i) \not{\bar{p}} \gamma^\nu (\not{\bar{p}} - \not{k}) (\not{p} - \not{k}) \gamma_\nu \not{\not{p}} \\ &= 2 \cdot 4 \text{tr} \not{\bar{p}} (\not{\bar{p}} - \not{k}) (\not{p} - \not{k}) \not{\not{p}} \\ &= 2 \cdot 4 \cdot 4 \bar{p} \cdot \bar{p} (\bar{p} - k) \cdot (p - k) = 16 \bar{p} \cdot \bar{p} M^2 \end{aligned}$$

$\therefore$  all:

$$\begin{aligned} \frac{1}{4} \sum |M|^2 &= \frac{e^2 g^2}{4 \cdot 4} \left\{ \frac{32 \bar{p} \cdot k \bar{p} \cdot k}{(\bar{p} \cdot k)^2} + \frac{32 \bar{p} \cdot k \bar{p} \cdot k}{(\bar{p} \cdot k)^2} \right. \\ &\quad \left. + 2 \frac{16 \bar{p} \cdot \bar{p} M^2}{(\bar{p} \cdot k)(\bar{p} \cdot k)} \right\} \\ &= 2 g^2 e^2 \left\{ \frac{\bar{p} \cdot k}{\bar{p} \cdot k} + \frac{\bar{p} \cdot k}{\bar{p} \cdot k} + \frac{\bar{p} \cdot \bar{p} M^2}{\bar{p} \cdot k \bar{p} \cdot k} \right\} \end{aligned}$$

In the CM system:

$$p = (E, 0, 0, E)$$

$$\bar{p} = (E, 0, 0, -E)$$

$$k = (p, p \sin \theta, 0, p \cos \theta)$$

$$q = (E, -p \sin \theta, 0, -p \cos \theta)$$

$$\text{st. } \quad \varepsilon^2 - p^2 = M^2 \quad \varepsilon + p = 2E$$

$$p = \frac{4E^2 - M^2}{2E} \quad \varepsilon = \frac{4E^2 + M^2}{2E}$$

$$p \cdot k = E p (1 - \cos \theta) \quad \bar{p} \cdot k = E p (1 + \cos \theta)$$

$$p \cdot \bar{p} = 2E^2$$

$$\begin{aligned} \frac{1}{4} \sum |M|^2 &= 2e^2 g^2 \left\{ \frac{1 + \cos \theta}{1 - \cos \theta} + \frac{1 - \cos \theta}{1 + \cos \theta} + \frac{2E^2 M^2}{E^2 p^2 (1 - \cos^2 \theta)} \right\} \\ &= 4e^2 g^2 \left\{ \frac{1 + \cos^2 \theta}{1 - \cos^2 \theta} + \frac{(M^2/p^2)}{(1 - \cos^2 \theta)} \right\} \end{aligned}$$

$$\frac{d\sigma}{d\cos \theta} = \frac{1}{8E^2} \frac{1}{16\pi} \frac{2p}{2E} 4e^2 g^2 \left\{ \frac{M^2/p^2 + (1 + \cos^2 \theta)}{1 - \cos^2 \theta} \right\}$$

$$\frac{d\sigma}{d\cos \theta} = \frac{1}{2E^3 p} \pi \alpha \cdot \frac{g^2}{4\pi} \cdot \left( \frac{M^2 + p^2 (1 + \cos^2 \theta)}{1 - \cos^2 \theta} \right)$$

where, again  $P = \frac{4E^2 - M^2}{2E}$

c.) Restoring the electron mass in the t-channel propagator, the denominator is

$$(p-k)^2 - m_e^2 = p^2 - 2p \cdot k - m_e^2 = -2p \cdot k$$

but not  $p = (E, 0, 0, p_e)$   $p_e = (E^2 - m_e^2)^{1/2}$   
 $\approx E - \frac{m_e^2}{2E} + \dots$

$$p \cdot k = (E - p_e \cos \theta) P$$

Using  $p_e \cong E - \frac{m_e^2}{2E} + \dots$

$$p \cdot k = pE \left\{ (1 - \cos \theta) + \frac{m_e^2}{2E^2} \cos \theta + \dots \right\}$$

$$\cong pE \left( 1 - \cos \theta + \frac{m_e^2}{2E^2} \right) \quad \text{as } \cos \theta \rightarrow 1$$

then

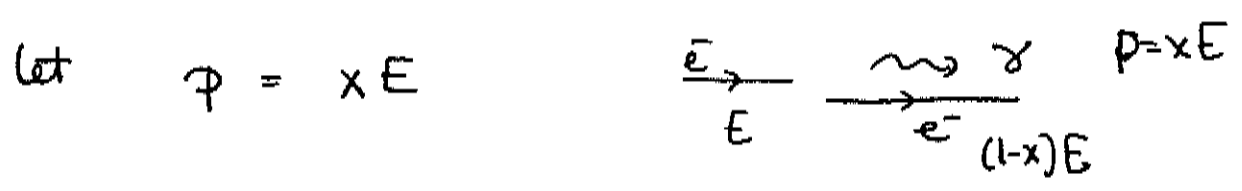
$$\int_0^{\pi} d\cos \theta \frac{1}{1 - \cos \theta} \rightarrow \int_0^{\pi} d\cos \theta \frac{1}{(1 - \cos \theta + \frac{m_e^2}{2E^2})}$$

$$= \log \frac{2E^2}{m_e^2}$$

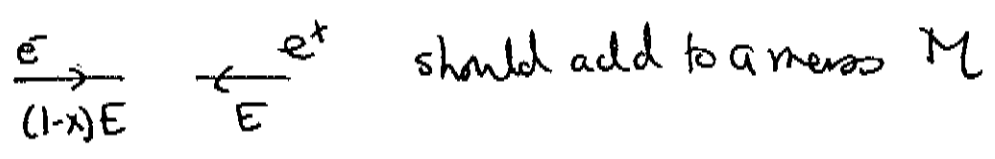
~~integrate~~  $\frac{d\sigma}{d\cos \theta}$  over forward  $\gamma_0$ , we find

$$\sigma(e^+e^- \rightarrow \text{forward } \gamma + B)$$

$$\cong \frac{\pi \alpha}{2E^3 p} \frac{q^2}{4\pi} \frac{M^2 + p^2 \cdot 2}{2} \log \left( \frac{2E^2}{m_e^2} \right)$$



$x$  is fixed by the condition that



$$M^2 = ((1-x)P + \bar{P})^2 = 2(1-x)P \cdot \bar{P}$$

$$\text{or } (1-x) = M^2/S$$

$$P^2 = x^2 E^2 = \frac{x^2 S}{4}$$

can replace  $\ln \frac{2E^2}{m_e^2} \rightarrow \ln \frac{S}{m_e^2}$   
with small error.

then

$$\sigma(e^+e^- \rightarrow \text{femal } \gamma + B)$$

$$= \frac{\pi \alpha}{2 E^4 x} \frac{g^2}{4\pi} \frac{1}{2} [(1-x) + x^2/2] \cdot S \cdot \ln \frac{S}{m_e^2}$$

$$= \int_0^1 dx \cdot S \cdot \pi g^2 \delta(M^2 - (1-x)S)$$

$$\cdot \frac{\alpha}{S^2 x} \frac{1}{2\pi} (2 - 2x + x^2) \cdot S \cdot \ln \frac{S}{m_e^2}$$

$$= \int_0^1 dx \pi g^2 \delta(M^2 - (1-x)S)$$

$$\cdot \frac{\alpha}{2\pi} \frac{1}{x} [1 + (1-x)^2] \ln \frac{S}{m_e^2}$$

$$= \int_0^1 dx \left\{ \frac{\alpha}{2\pi} \frac{1}{x} [1 + (1-x)^2] \ln \frac{S}{m_e^2} \right\} [\sigma(e^+e^- \rightarrow B)]$$

$E_m^2 = (1-x)S$