

Physics 330 - Problem Set 1

Solutions

1.)

I set $c=1$, so

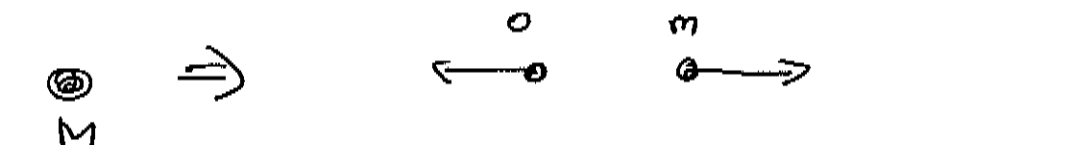
for a massless particle moving $\parallel \hat{z}$:

$$p^\mu = (p, 0, 0, p)$$

for a massive particle moving $\parallel \hat{z}$

$$p^\mu = (E_p, 0, 0, p) \quad E_p = [|\vec{p}|^2 + m^2]^{1/2}$$

a.)



$\text{at rest} \quad p^\mu = (p, 0, 0, -p) \quad p^\mu = (E_p, 0, 0, p)$

$$p + E_p = M \quad \text{conserv. of energy.}$$

then

$$[p^2 + m^2]^{1/2} = M - p$$

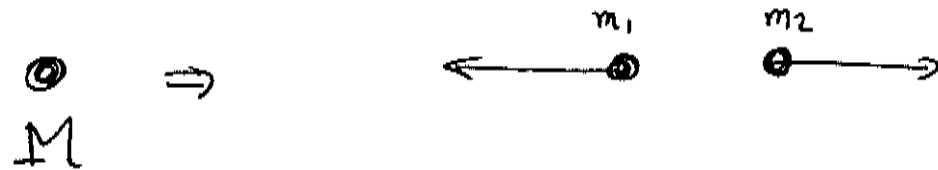
$$p^2 + m^2 = M^2 - 2Mp + p^2$$

$$p = \frac{M^2 - m^2}{2M}$$

↑
energy + mom of massless particle

$$E_p = \frac{M^2 + m^2}{2M}$$

energy of massive particle

b.) 

$$p_1 = (E_1, 0, 0, p) \quad p_2 = (E_2, 0, 0, p)$$

$$E_1 = (p^2 + m_1^2)^{1/2} \quad E_2 = (p^2 + m_2^2)^{1/2}$$

$$E_1 + E_2 = M$$

$$E_1^2 + E_2^2 + 2E_1E_2 = M^2$$

$$4E_1^2E_2^2 = (M^2 - E_1^2 - E_2^2)^2$$

$$4E_1^2E_2^2 = M^4 - 2M^2(E_1^2 + E_2^2) + (E_1^2 + E_2^2)^2$$

$$0 = M^4 - 2M^2(E_1^2 + E_2^2) + (E_1^2 - E_2^2)^2$$

$$0 = M^4 - 2M^2(2p^2 + m_1^2 + m_2^2) + (p^2 + m_1^2 - p^2 + m_2^2)^2$$

so

$$4p^2M^2 = M^4 - 2M^2(m_1^2 + m_2^2) + (m_1^2 - m_2^2)$$

$$p = \frac{[M^4 - 2M^2(m_1^2 + m_2^2) + (m_1^2 - m_2^2)]^{1/2}}{2M}$$

$$E_1^2 = p^2 + m_1^2 = \frac{M^4 - 2M^2(m_1^2 + m_2^2) + (m_1^2 - m_2^2)^2 + 4M^2m_1^2}{4M^2}$$

$$= \frac{M^4 + 2M^2(m_1^2 - m_2^2) + (m_1^2 - m_2^2)^2}{4M^2} = \left(\frac{M^2 + m_1^2 - m_2^2}{2M} \right)^2$$

so

$$p = \frac{[M^4 - 2M^2(m_1^2 + m_2^2) + (m_1^2 - m_2^2)^2]^{1/2}}{2M}$$

$$E_1 = \frac{M^2 + m_1^2 - m_2^2}{2M} \quad E_2 = \frac{M^2 + m_2^2 - m_1^2}{2M}$$

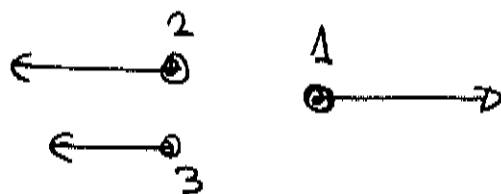
One sometimes sees the notation:

$$\Delta(a, b, c) = a^2 + b^2 + c^2 - 2ab - 2ac - 2bc$$

$$p = \frac{\sqrt{\Delta(M^2, m_1^2, m_2^2)}}{2M}$$

c) To answer this question, we should ask: What configurations allow the energies E_1 , E_2 , E_3 to take their maximum values?

For E_1 , this configuration is



1 recoils against 2 + 3: $p_1 = p_2 + p_3$

For all massless particles $p_1 = E_1$ $p_2 = E_2$ $p_3 = E_3$

$$E_1 + E_2 + E_3 = M \quad \text{so} \quad p_1 = M/2$$

In terms of the x_i $x_i = 2E_i/M$

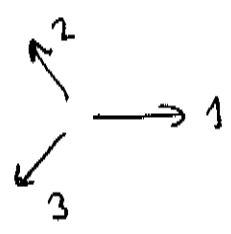
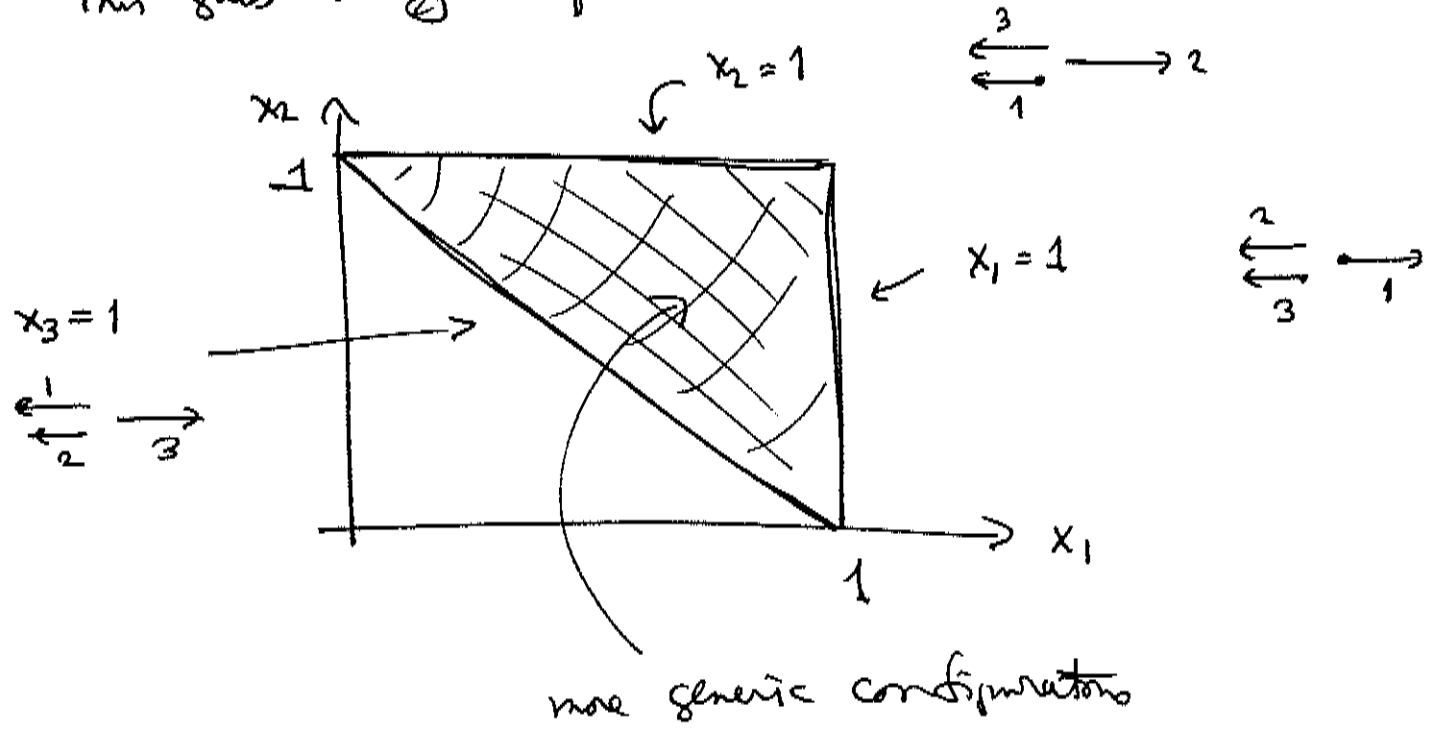
$x_1 = 1$ $x_2 + x_3 = 1$

The other boundaries correspond to

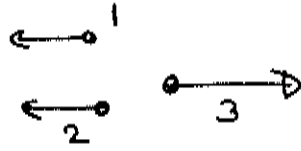
$x_2 = 1$ $x_1 + x_3 = 1$

$x_3 = 1$ $x_1 + x_2 = 1$

This gives a region of the $x_1 x_2$ plane:



d.) If particle 3 has mass m :



how has $p_3 = p_1 + p_2$ $E_3 = (p_3^2 + m^2)^{1/2}$

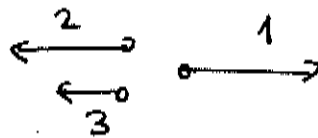
1+2 is a massless 4-vector $(p_1 + p_2)^2 = 0$ or

$(p_1 + p_2)^\mu = (p_3, 0, 0, -p_3)$. Then this is the problem we solved in part (a).

$$p_3 = \frac{M^2 - m^2}{2M} \quad \Rightarrow \quad x_1 + x_2 = \left(1 - \frac{m^2}{M^2}\right)$$

$$E_3 = \frac{M^2 + m^2}{2M}$$

For the configuration



$$p_1 = p_2 + p_3$$

$$M = p_1 + p_2 + (p_3^2 + m^2)^{1/2}$$

$$M = p_1 - p_2 = ((p_1 - p_2)^2 + m^2)^{1/2}$$

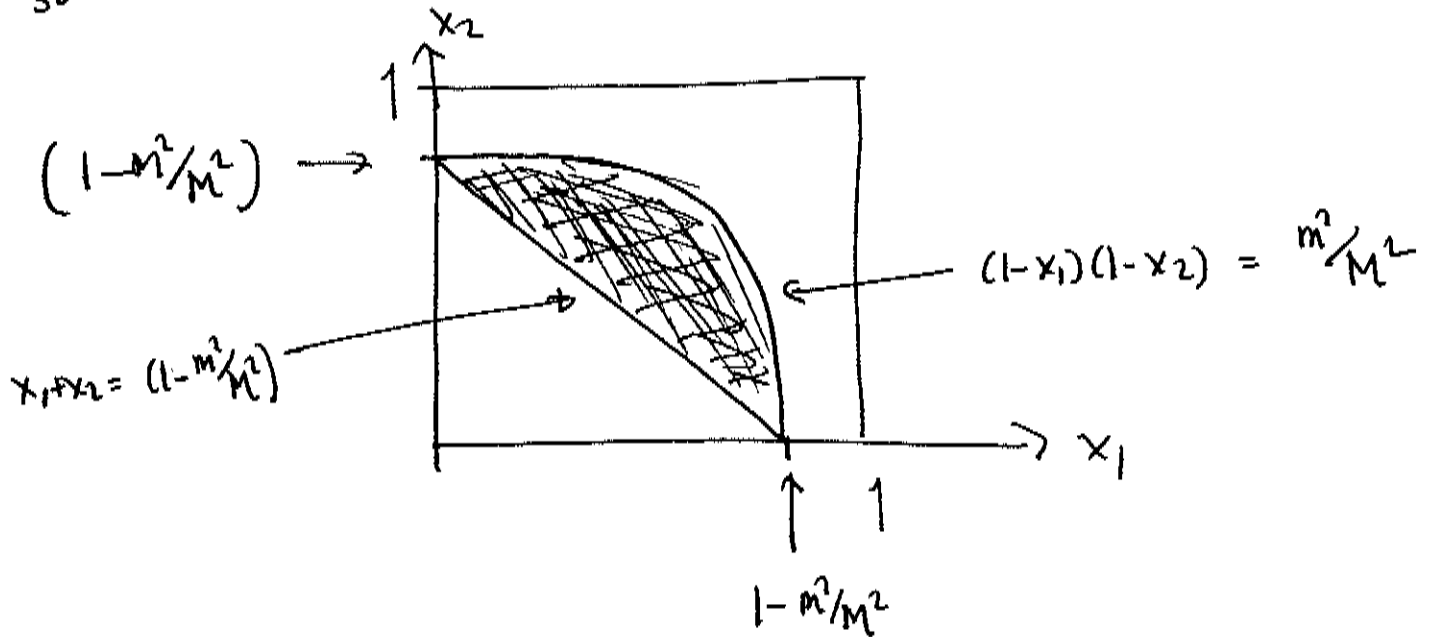
$$M^2 - 2M(p_1 + p_2) + (p_1 + p_2)^2 = (p_1 - p_2)^2 + m^2$$

$$M^2 - 2M(p_1 + p_2) + 4p_1 p_2 = m^2$$

$$(M - 2p_1)(M - 2p_2) = m^2$$

$$(1 - x_1)(1 - x_2) = m^2/M^2$$

So the allowed kinematic region is now:



2.) a)
$$R = \begin{pmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{pmatrix}$$

$$R \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} \sin \theta \\ 0 \\ \cos \theta \end{pmatrix}$$

$$R \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} \cos \theta \\ 0 \\ -\sin \theta \end{pmatrix}$$

$$R \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad \text{this is the axis of the rotation}$$

$$\begin{aligned}
 \text{b.) } \frac{d}{d\theta} R &= \begin{pmatrix} -\sin\theta & 0 & \cos\theta \\ 0 & 0 & 0 \\ -\cos\theta & 0 & -\sin\theta \end{pmatrix} \\
 &= \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \cos\theta & 0 & \sin\theta \\ 0 & 1 & 0 \\ -\sin\theta & 0 & \cos\theta \end{pmatrix}
 \end{aligned}$$

so $R(\theta)$ satisfies

$$\frac{d}{d\theta} R(\theta) = -i \begin{bmatrix} & & \\ & J^2 & \\ & & \end{bmatrix} R(\theta)$$

where $J^2 = \begin{bmatrix} 0 & 0 & i \\ 0 & 0 & 0 \\ -i & 0 & 0 \end{bmatrix}$

the solution of this equation with initial condition $R(\theta=0) = 1$

is $R(\theta) = \exp[-i J^2 \theta]$

c.) In the σ_y - $\frac{1}{2}$ representation:

$$e^{-i J^2 \theta} = e^{-i \sigma_y^2 \frac{\theta}{2}} \quad \sigma_y^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$(\sigma_y^2)^2 = 1$$

$$= \cos \frac{\theta}{2} - i \sigma_y^2 \sin \frac{\theta}{2} = \begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$$

d) In the $Spin-1$ representation:

$$\text{rotate by } \phi \text{ about } \hat{3} = R_1 = \begin{pmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\text{rotate by } \theta \text{ about } \hat{2} = R_2 = \begin{pmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{pmatrix}$$

$$\text{rotate by } -\phi \text{ about } \hat{3} : R_3 = \begin{pmatrix} \cos \phi & \sin \phi & 0 \\ -\sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

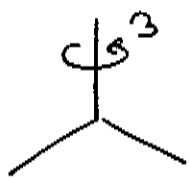
$$R_3 R_2 R_1 = \begin{pmatrix} \cos \phi & \sin \phi & & \\ -\sin \phi & \cos \phi & & \\ & & 1 & \\ & & & 1 \end{pmatrix} \begin{pmatrix} \cos \theta & & \sin \theta & \\ & 1 & & \\ -\sin \theta & & \cos \theta & \\ & & & 1 \end{pmatrix} \begin{pmatrix} \cos \phi & -\sin \phi & & \\ \sin \phi & \cos \phi & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$$

$$= \begin{pmatrix} \cos \phi \cos \theta & \sin \phi & \cos \phi \sin \theta & \\ -\sin \phi \cos \theta & \cos \phi & -\sin \phi \sin \theta & \\ & & \cos \theta & \\ -\sin \theta & 0 & & 1 \end{pmatrix}$$

$$= \begin{pmatrix} \cos^2 \phi \cos \theta + \sin^2 \phi & -\sin \phi \cos \phi \cos \theta + \sin \phi \cos \phi & \cos \phi \sin \theta \\ \sin \phi \cos \phi \cos \theta + \sin \phi \cos \phi & \sin^2 \phi \cos \theta + \cos^2 \phi & -\sin \phi \sin \theta \\ -\sin \theta \cos \phi & \sin \theta \sin \phi & \cos \theta \end{pmatrix}$$

what is this?

we have done rotations



It is obvious that the \hat{z} axis should go into

$$\text{rotate of } (\sin \theta, 0, \cos \theta)$$

$$= (\cos \phi \sin \theta, -\sin \phi \sin \theta, \cos \theta)$$

but the rest is less obvious.

Try the $\sigma^y - \frac{1}{2}$ case:

$$R_1 = e^{-i\sigma^z \frac{\phi}{2}} = \begin{pmatrix} e^{-i\phi/2} & 0 \\ 0 & e^{+i\phi/2} \end{pmatrix}$$

$$R_2 = \begin{pmatrix} \cos \theta/2 & -\sin \theta/2 \\ \sin \theta/2 & +\cos \theta/2 \end{pmatrix}$$

$$R_3 = \begin{pmatrix} e^{+i\phi/2} & 0 \\ 0 & e^{-i\phi/2} \end{pmatrix}$$

so

$$R_3 R_2 R_1 = \begin{pmatrix} \cos \theta/2 e^{+i\phi/2} & -\sin \theta/2 e^{+i\phi/2} \\ \sin \theta/2 e^{-i\phi/2} & \cos \theta/2 e^{-i\phi/2} \end{pmatrix} \begin{pmatrix} e^{-i\phi/2} & 0 \\ 0 & e^{+i\phi/2} \end{pmatrix}$$

$$= \begin{pmatrix} \cos \theta/2 & -\sin \theta/2 e^{+i\phi} \\ \sin \theta/2 e^{-i\phi} & \cos \theta/2 \end{pmatrix}$$

now

$$e^{-i \vec{\theta} \cdot \vec{\sigma} / 2} = \cos \frac{\theta}{2} - i \hat{\theta} \cdot \vec{\sigma} \sin \frac{\theta}{2} \quad \theta = |\vec{\theta}|$$

the above is:

$$\begin{aligned} R_3 R_2 R_1 &= \cos \theta/2 \cdot 1 - \cos \phi \sin \theta/2 \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} - i \sin \phi \sin \theta/2 \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ &= \cos \theta/2 - i [\sin \phi \sigma^1 + \cos \phi \sigma^2] \sin \theta/2 \end{aligned}$$

so this is a rotation about the axis

$$\hat{n} = \begin{pmatrix} \sin \phi \\ \cos \phi \\ 0 \end{pmatrix}$$

by an angle θ .

There are some relatively easy ways to check that p. 8
is also a rotation of θ about the axis \hat{n}

For any rotation by θ : $\text{tr } R = \cos \theta + \cos \theta + 1 = 1 + 2\cos \theta$

For any rotation about \hat{n} : $R \hat{n} = \hat{n}$

Let's check

$$\text{tr}(R_3 R_2 R_1) = \cos^2 \phi \cos \Theta + \sin^2 \phi + \sin^2 \phi \cos \Theta + \cos^2 \phi + \cos \Theta$$

$$= 2 \cos \Theta + 1 \quad \checkmark$$

$$(R_3 R_2 R_1) \begin{pmatrix} \sin \phi \\ \cos \phi \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} (\cos^2 \phi \cos \Theta + \sin^2 \phi) \sin \phi + (-\sin \phi \cos \phi \cos \Theta + \sin \phi \cos \phi) \cos \phi \\ (-\sin \phi \cos \phi \cos \Theta + \sin \phi \cos \phi) \sin \phi + (\sin^2 \phi \cos \Theta + \cos^2 \phi) \cos \phi \\ (-\sin \Theta \cos \phi) \sin \phi + (\sin \Theta \sin \phi) \cos \phi \end{pmatrix}$$

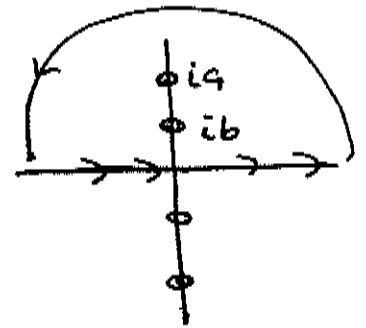
$$= \begin{pmatrix} \sin^2 \phi + \sin \phi \cos^2 \phi \\ \sin^2 \phi \cos \phi + \cos^3 \phi \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} \sin \phi \\ \cos \phi \\ 0 \end{pmatrix} = \hat{n} \quad \checkmark$$

∞ $R_3 R_2 R_1$ is the spin $\frac{1}{2}$ of spin 1 cases are rotations about the same axis by the same angle.

$$3) a) \quad I_1 = \int_0^{\infty} dx \frac{x^2}{(a^2+x^2)(b^2+x^2)}$$

$$= \frac{1}{2} \int_{-\infty}^{\infty} dx \frac{x^2}{(a^2+x^2)(b^2+x^2)}$$



$$= 2\pi i \cdot \frac{1}{2} \cdot \left\{ \begin{array}{l} \text{(Residue at } x=ia) \\ + \text{(Residue at } x=ib) \end{array} \right\}$$

$$= \pi i \left\{ \frac{(ia)^2}{2ia(b^2-a^2)} + \frac{(ib)^2}{(a^2-b^2)(2ib)} \right\}$$

$$= \frac{\pi}{2(a^2-b^2)} \{ a-b \}$$

$$\text{so } I_1 = \frac{\pi}{2} \frac{1}{(a+b)}$$

a more direct method:
$$I_1 = \frac{1}{2} \int_{-\infty}^{\infty} dx \left(\frac{a^2}{a^2+x^2} - \frac{b^2}{b^2+x^2} \right) \frac{1}{(a^2-b^2)}$$

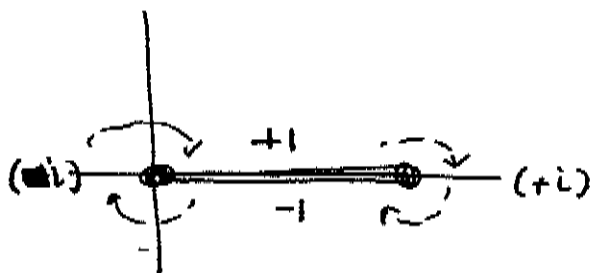
$$\text{now } \int_{-\infty}^{\infty} dx \frac{1}{a^2+x^2} = \frac{\pi}{a}$$

$$\text{so } I_1 = \frac{\pi}{2} (a-b) \frac{1}{a^2-b^2} = \frac{\pi}{2} \frac{1}{(a+b)} \quad \checkmark$$

b) To do this exercise, first look at the analytic structure of $\frac{1}{\sqrt{x(1-x)}}$: Define this so that

$$\frac{1}{\sqrt{x(1-x)}} = + \frac{1}{\sqrt{|x(1-x)|}} \quad (\text{real + positive})$$

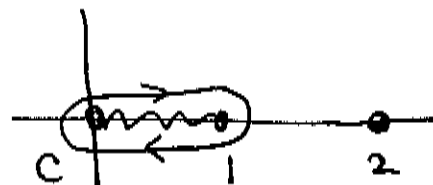
just above the real axis in the x plane. Then the phase of $\frac{1}{\sqrt{x(1-x)}}$ at other points is:



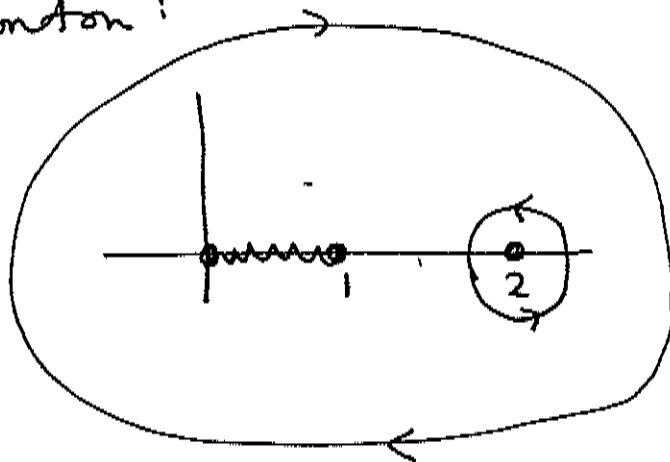
Notice that

$$I_2 = \int_0^1 dx \frac{1}{\sqrt{x(1-x)}} \frac{1}{x-2}$$

$$= \frac{1}{2} \int_C dx \frac{1}{\sqrt{x(1-x)}} \frac{1}{x-2}$$



Now deform the contour:



The big contour at ∞ gives $\int dx \frac{1}{x^2} \rightarrow 0$

The pole at $x=2$ gives

$$\begin{aligned} I_2 &= \frac{1}{2} 2\pi i \left[\frac{1}{\sqrt{x(1-x)}} \right]_{x=2} \\ &= \frac{1}{2} 2\pi i \frac{(+i)}{\sqrt{|2 \cdot 1|}} \\ &\quad \uparrow \\ &\quad \text{[phase from p.13]} \end{aligned}$$

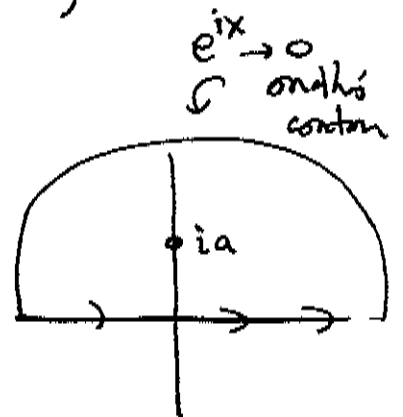
$$I_2 = -\frac{\pi}{\sqrt{2}}$$

$$\begin{aligned} \text{e.) } I_3 &= \int_0^{\infty} dx \frac{x \sin x}{a^2 + x^2} \\ &= \frac{1}{2} \int_{-\infty}^{\infty} dx \frac{x}{a^2 + x^2} \left(\frac{e^{ix} - e^{-ix}}{2i} \right) \end{aligned}$$

$$= \int_{-\infty}^{\infty} dx \frac{x}{a^2 + x^2} \frac{e^{ix}}{2i}$$

$$= \frac{2\pi i}{2i} \frac{ia}{2ia} e^{i(ia)}$$

$$I_3 = \frac{\pi}{2} e^{-a}$$



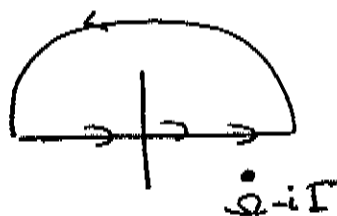
d.) The important point for this integral is that

$$e^{-i\omega t} \rightarrow 0 \text{ for } \text{Im } \omega \rightarrow \infty \text{ when } t < 0$$

$$\text{for } \text{Im } \omega \rightarrow -\infty \text{ when } t > 0$$

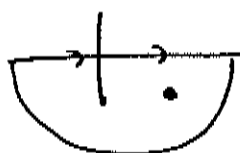
so

$t < 0$



$$I_4 = 0$$

$t > 0$



$$I_4 = \int \frac{d\omega}{2\pi} \frac{e^{-i\omega t}}{\omega - (\Omega - i\Gamma)}$$

$$= -2\pi i \frac{1}{2\pi} e^{-it(\Omega - i\Gamma)}$$

so

$$I_4 = \begin{cases} 0 & t < 0 \\ -i e^{-i\Omega t} e^{-\Gamma t} & t > 0 \end{cases}$$