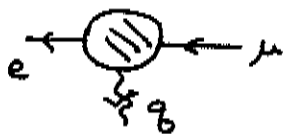


Physics 330 - Final Exam

Solutions

a) $i\mathcal{M} =$ 

$$= -ie \bar{u}(p') i\sigma^{\mu\nu} (-q_\nu) \bar{F}_2(q^2) u(p) \Big|_{q^2=0}$$

$$\sum_{\text{spin}} |\mathcal{M}|^2 = e^2 |\bar{F}_2|^2 \text{tr} \left\{ \underbrace{(\not{p}' + m_e)}_{\text{ignore}} [-i\sigma^{\mu\nu} q_\nu] \underbrace{(\not{p} + m_\mu)}_{\text{ignore}} [i\sigma^{\lambda\sigma} q_\sigma] \right\} (-g_{\mu\lambda})$$

$$= e^2 |\bar{F}_2(q^2)|^2 \text{tr} \left(\not{p}' \frac{1}{2} [\gamma^\mu, \not{q}] (\not{p} + m_\mu) \frac{1}{2} [\gamma^\lambda, \not{q}] \right) (-g_{\mu\lambda})$$

the traces with m_μ are 0. The other terms are:

$$= e^2 \frac{|\bar{F}_2(q^2)|^2}{4} \text{tr} \left[\not{p}' (\gamma^\mu \not{q} - \not{q} \gamma^\mu) \not{p} (\gamma_\lambda \not{q} - \not{q} \gamma_\lambda) \right] (+1)$$

$$= \frac{e^2 |\bar{F}_2(q^2)|^2}{4} \text{tr} \left\{ \not{p}' \gamma^\mu \not{q} \not{p} \gamma_\lambda \not{q} - \not{p}' \gamma^\mu \not{q} \not{p} \not{q} \gamma_\lambda \right. \\ \left. - \not{p}' \not{q} \gamma^\mu \not{p} \gamma_\lambda \not{q} + \not{p}' \not{q} \gamma^\mu \not{p} \not{q} \gamma_\lambda \right\}$$

$$= \frac{e^4 |\bar{F}_2(q^2)|^2}{4} \text{tr} \left\{ 4 \not{p}' \not{q} \not{p} \not{q} + 2 \not{p}' \not{q} \not{p} \not{q} \right. \\ \left. + 2 \not{p}' \not{q} \not{p} \not{q} + 4 \not{p}' \not{q} \not{p} \not{q} \right\}$$

$$= \frac{e^2 |\bar{F}(\omega)|^2}{4} \cdot 4 \cdot \left\{ 4 \cdot 2 \cdot \underline{p' \cdot q} \cdot \underline{p \cdot q} + 2 \cdot 2 \cdot \left(2 \cdot \underline{p' \cdot q} \cdot \underline{p \cdot q} - \underline{p' \cdot p} \cdot \underline{q^2} \right) \right\} \\ = 0$$

$$= e^2 |\bar{F}(\omega)|^2 16 \underline{p' \cdot q} \underline{p \cdot q}$$

$$\text{low } m_\mu^2 = p^2 = (p' + q)^2 = 2 \underline{p' \cdot q} = 2 \underline{(p + q) \cdot q} = 2 \underline{p \cdot q}$$

$$\text{so } 2 \underline{p' \cdot q} = 2 \underline{p \cdot q} = m_\mu^2$$

$$\sum |M|^2 = 4 e^2 |\bar{F}_2(\omega)|^2 m_\mu^4$$

$$I(\mu \rightarrow e\gamma) = \frac{1}{2m_\mu} \cdot \underbrace{\frac{1}{8\pi}}_{\text{phase space}} \cdot \frac{1}{2} \sum_{\text{spin}} |M|^2$$

$$= \frac{1}{32\pi m_\mu} \cdot 4 e^2 |\bar{F}_2(\omega)|^2 m_\mu^4$$

$$I(\mu \rightarrow e\gamma) = \frac{\alpha}{2} |\bar{F}_2(\omega)|^2 m_\mu^3$$

b.) with $\bar{F}_2(0) = G \frac{\alpha}{4\pi} \frac{m_\mu}{M^2}$

$$I(\mu \rightarrow e\gamma) = \frac{\alpha^3}{32\pi^2} G^2 \frac{m_\mu^5}{M^4}$$

compare this to

$$I(\mu) = \frac{\alpha_w^2 m_\mu^5}{384\pi m_w^4} \quad 384 = 32 \cdot 12$$

$$BR(\mu \rightarrow e\gamma) = \frac{12\alpha^3}{40\alpha_w^2} G^2 \left(\frac{m_w}{M}\right)^4$$

$$\text{for } \frac{m_w}{M} = \frac{80 \text{ GeV}}{200 \text{ GeV}} = G^2 \cdot (3.3 \times 10^{-5})$$

so $BR(\tau \rightarrow e/\mu \gamma) = G^2 \cdot (6.3 \times 10^{-6})$ (divide by 5.3)

$$BR(\mu \rightarrow e\gamma) < 1.2 \times 10^{-11} \Rightarrow G_{\mu e}^2 < 3.6 \times 10^{-7} \text{ or } G_{\mu e} < 6 \times 10^{-4}$$

$$BR(\tau \rightarrow e\gamma) < 2.7 \times 10^{-6} \Rightarrow G_{\tau e}^2 < 0.43 \text{ or } G_{\tau e} < 0.66$$

$$BR(\tau \rightarrow \mu\gamma) < 3.2 \times 10^{-7} \Rightarrow G_{\tau \mu}^2 < 5.1 \times 10^{-2} \text{ or } G_{\tau \mu} < 0.23$$

In a theory in which $G_{\mu e} \sim G_{\tau \mu} \sim G_{\tau e}$, the $\mu \rightarrow e\gamma$ constraint is obviously the strongest.

Since $(m_\tau/m_\mu)^2 = 280$, the $\mu \rightarrow e\gamma$ constraint is stronger even if $G_{\mu \rightarrow \mu e} \sim m_e^2$.

c.) As we learned,

$$\sigma(e^+e^- \rightarrow \tau^+\tau^-) \underset{s \gg m_\tau^2}{\sim} \frac{4\pi\alpha^2}{3s} = \frac{86.8 \text{ nb}}{[E_{\text{cm}} (\text{GeV})]^2}$$

$$= 0.78 \text{ nb} \quad \text{at } E_{\text{cm}} = 10.58 \text{ GeV}$$

then 86.3 fb^{-1} gives

$$\# \text{ of events} = (0.78 \times 10^{-9}) (86.3 \times 10^{15})$$

$$= 7.0 \times 10^7 \quad \tau^+\tau^- \text{ pairs}$$

If we could record all of these events and saw no $\tau \rightarrow \mu \gamma$ decays. (< 2.3 decays at the 90% CL) we would have a limit of 1.5×10^{-8} (ideal limit)

In practice, the BELLE result is quite an efficient use of the data.

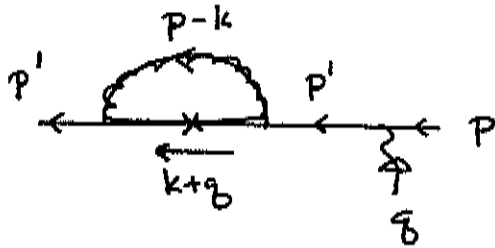
BABAR also has a huge and growing data sample of τ decays
— this could be your thesis!

d) The propagators are

$$e: \quad \overleftarrow{p} \quad \frac{i}{\not{p} - m_e} \quad \mu: \quad \overleftarrow{p} \quad \frac{i}{\not{p} - m_\mu}$$

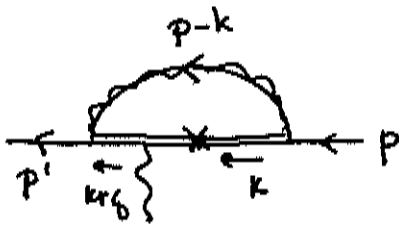
$$\phi_e: \quad \overleftarrow{p} \quad \frac{i}{p^2 - m^2} \quad \phi_\mu: \quad \overleftarrow{p} \quad \frac{i}{p^2 - m^2}$$

$$b: \quad \overleftarrow{p} \quad \frac{i}{\not{p} - m_b}$$



$$= \bar{u}(p') \int \frac{d^4 k}{(2\pi)^4} (ig) \frac{i(\not{p}-k+m_b)}{(p-k)^2 - m_b^2} (ig) \frac{i(\not{p}'+m_\mu)}{(p')^2 - m_\mu^2} (-ie\gamma^\mu) u(p) \\ \cdot \frac{i}{(k+q)^2 - m^2} (-i\delta m^2) \frac{i}{(k+q)^2 - m^2}$$

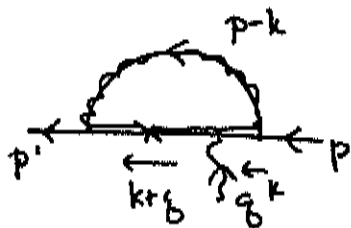
$$= eg^2 \delta m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{\not{p}-k+m_b}{(p-k)^2 - m_b^2} \frac{1}{[(k+q)^2 - m^2]^2} \frac{\not{p}'+m_\mu}{(p')^2 - m_\mu^2} \gamma^\mu u(p)$$



$$= \bar{u}(p') \int \frac{d^4 k}{(2\pi)^4} (ig) \frac{i[\not{p}-k+m_b]}{(p-k)^2 - m_b^2} ig u(p)$$

$$\cdot \frac{i}{(k+q)^2 - m^2} [-ie(2k+q)^\mu] \frac{i}{k^2 - m^2} (-i\delta m^2) \frac{i}{k^2 - m^2}$$

$$= eg^2 \delta m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{\not{p}-k+m_b}{(p-k)^2 - m_b^2} \frac{(2k+q)^\mu}{(k+q)^2 - m^2 (k^2 - m^2)^2} u(p)$$



$$\begin{aligned}
 &= \bar{u}(p') \int \frac{d^4 k}{(2\pi)^4} (ig) \frac{i \not{p} - \not{k} + m_b}{(p-k)^2 - m_b^2} (ig) u(p) \\
 &\quad \cdot \frac{i}{(k+q)^2 - m^2} (-i \delta m^2) \frac{i}{(k+q)^2 - m^2} - ie (2k+q)^\mu \frac{i}{k^2 - m^2} \\
 &= eg^2 \delta m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{\not{p} - \not{k} + m_b}{(p-k)^2 - m_b^2} \frac{(2k+q)^\mu}{[(k+q)^2 - m^2]^2 (k^2 - m^2)} u(p)
 \end{aligned}$$

f.) In the first diagram, we: $q = p' - p$

$$\begin{aligned}
 \not{q} &= \not{p}' - \not{p} \\
 \not{q} \cdot \frac{\not{p} + m_e}{p^2 - m_e^2} &= (p' - p) \frac{\not{p} + m_e}{p^2 - m_e^2} \\
 &= [(\not{p}' - m_e) - (\not{p} - m_e)] \frac{1}{p - m_e}
 \end{aligned}$$

now $\bar{u}(p') (\not{p}' - m_e) = 0$

so

$$\not{q} \cdot \frac{\not{p} + m_e}{p^2 - m_e^2} = -eg^2 \delta m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{\not{p} - \not{k} + m_b}{(p-k)^2 - m_b^2} \frac{1}{(k^2 - m^2)^2} u(p)$$

In the second diagram in (e), use

$$\frac{(p'+m_\mu)}{(p')^2 - m_\mu^2} \gamma^\mu \not{q} u(p)$$

$$= \frac{1}{p' - m_\mu} [(p' - m_\mu) - (p' - m_\mu)] u(p) = u(p)$$

$$\text{with } (p' - m_\mu) u(p) = 0$$

then

$$\not{q} \cdot \frac{\text{diagram}}{\text{denominator}} = + e g^2 S m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{p - k + m_b}{(p - k)^2 - m_b^2} \frac{1}{[(k + q)^2 - m^2]^2} u(p)$$

In the last two diagrams, we:

$$\begin{aligned} \not{q}^\mu (2k + q)_\mu &= 2k \cdot q + q^2 = (k + q)^2 - k^2 \\ &= [(k + q)^2 - m^2] - [k^2 - m^2] \end{aligned}$$

so

$$\not{q}^\mu \left(\frac{\text{diagram}}{\text{denominator}} + \frac{\text{diagram}}{\text{denominator}} \right)$$

$$= e g^2 S m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{p - k + m_b}{(p - k)^2 - m_b^2} [(k + q)^2 - m^2 - (k^2 - m^2)] u(p)$$

$$\left(\frac{1}{[(k + q)^2 - m^2][k^2 - m^2]^2} + \frac{1}{[(k + q)^2 - m^2]^2[k^2 - m^2]} \right)$$

$$= e g^2 S m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{p - k + m_b}{(p - k)^2 - m_b^2} u(p)$$

$$\left\{ \frac{1}{[k^2 - m^2]^2} - \frac{1}{[(k + q)^2 - m^2][k^2 - m^2]} + \frac{1}{[(k + q)^2 - m^2][k^2 - m^2]} - \frac{1}{[(k + q)^2 - m^2]^2} \right\}$$

$$g_{\mu}^{\alpha} \left(\text{diagram 1} + \text{diagram 2} \right)$$

$$= e g^2 \delta m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{\not{p} - \not{k} + m_b}{(p-k)^2 - m_b^2} \left(\frac{1}{[k^2 - m_a^2]^2} - \frac{1}{(k+q)^2 - m_a^2} \right) u(p)$$

which cancels the contribution from the last two diagrams.

g.) The Gordon identity is: $(\not{p} u(p) = m_{\mu} u(p) \quad \bar{u}(p') \not{p}' = \bar{u}(p') m_e)$

$$\bar{u}(p') [\not{p}' \gamma^{\mu} + \gamma^{\mu} \not{p}] u(p) = (m_e + m_{\mu}) \bar{u}(p') \gamma^{\mu} u(p)$$

$$= \bar{u}(p') \left(\frac{1}{2} \{ \not{p}', \gamma^{\mu} \} + \frac{1}{2} \{ \gamma^{\mu}, \not{p} \} - \frac{1}{2} [\gamma^{\mu}, \not{p}' - \not{p}] \right) u(p)$$

$$= \bar{u}(p') \left((p+p')^{\mu} + i \sigma^{\mu\nu} q_{\nu} \right) u(p)$$

later we'll use this in the form

$$\bar{u}(p') (p+p')^{\mu} u(p) = (m_e + m_{\mu}) \bar{u}(p') \gamma^{\mu} u(p) - \bar{u}(p') i \sigma^{\mu\nu} q_{\nu} u(p)$$

this will allow us to put the sum of diagrams into the form:

$$\text{-ie } \bar{u}(p') \left[\gamma^{\mu} F_1(q^2) + \not{q}^{\mu} F_3(q^2) + i \sigma^{\mu\nu} q_{\nu} F_2(q^2) \right] u(p)$$

$$g^{\mu} \text{ (this)} = 0 \Rightarrow \bar{u}(p') \left[\not{q}^{\mu} F_1(q^2) + \not{q}^2 F_3(q^2) + 0 \right] u(p) = 0$$

$$\bar{u}(p') \left[(q' - q) F_1(q^2) + \not{q}^2 F_3(q^2) \right] u(p) = 0$$

$$(m_e - m_{\mu}) F_1(q^2) + \not{q}^2 F_3(q^2) = 0$$

2

$$F_1(q^2) = \frac{q^2}{m_\mu - m_e} F_3(q^2)$$

if $F_3(q^2)$ has a smooth limit as $q^2 \rightarrow 0$, $F_1(q^2) \propto q^2 \rightarrow 0$
as $q^2 \rightarrow 0$.

h.) Now we have to get to work computing diagrams. A little table of integrals will be useful:

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{[k^2 - \Delta]^3} = \frac{-i}{(4\pi)^2} \cdot \frac{1}{2} \frac{1}{\Delta}$$

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{[k^2 - \Delta]^4} = \frac{i}{(4\pi)^2} \cdot \frac{1}{3 \cdot 2} \frac{1}{\Delta^2}$$

$$\int \frac{d^4 k}{(2\pi)^4} \frac{k^2}{[k^2 - \Delta]^4} = \frac{-i}{(4\pi)^2} \cdot \frac{2}{3 \cdot 2} \frac{1}{\Delta}$$

$$\int \frac{d^4 k}{(2\pi)^4} \frac{k^\mu k^\nu}{[k^2 - \Delta]^4} = \frac{-i}{(4\pi)^2} \frac{1}{3 \cdot 2} \frac{1}{\Delta} \cdot \frac{1}{2} g^{\mu\nu}$$

then

$$\int \frac{d^4k}{(2\pi)^4} \frac{\not{p} + m_e}{p^2 - m_e^2} \frac{\not{p} - \not{k} + m_b}{[(p-k)^2 - m_b^2] (k^2 - m^2)^2} u(p)$$

combine denominators using:

$$\frac{1}{[(p-k)^2 - m_b^2] (k^2 - m^2)^2} = \int_0^1 dz \frac{2z(1-z)}{[z(p-k)^2 - m_b^2 + (1-z)(k^2 - m^2)]^3}$$

The denominator is:

$$\begin{aligned} & z(k-p)^2 - z m_b^2 + (1-z)k^2 - (1-z)m^2 \\ &= k^2 - 2k \cdot zp + zp^2 - (z m_b^2 + (1-z)m^2) \\ &= k^2 + \underbrace{z(1-z)p^2 - (z m_b^2 + (1-z)m^2)}_{-\Delta} \end{aligned}$$

$$k = k - zp$$

$$k = k + zp$$

$$k-p = k - (1-z)p$$

integrates to 0

so

$$= eg^2 S m^2 \bar{u}(p') \gamma^\mu \frac{\not{p} + m_e}{p^2 - m_e^2} \int_0^1 dz \frac{2z(1-z)}{(2\pi)^4} \frac{((1-z)\not{p} - \not{k} + m_b)}{[k^2 - \Delta]^3} u(p)$$

$$= \frac{-ie g^2}{(4\pi)^2} S m^2 \bar{u}(p') \gamma^\mu \frac{\not{p} + m_e}{p^2 - m_e^2} \int_0^1 dz \frac{(1-z) [(1-z)\not{p} + m_b]}{[z m_b^2 + (1-z)m^2 - z(1-z)p^2]} u(p)$$

now use $\not{p} u(p) = m_\lambda u(p) \quad p^2 = m_\lambda^2$

$$\frac{\not{p} + m_e}{p^2 - m_e^2} = \frac{1}{m_\lambda - m_e}$$

so

$$\int \frac{d^4 k}{(2\pi)^4} = (-ie) \frac{g^2 \Delta m^2}{(4\pi)^2} \frac{1}{m_\lambda - m_e} \bar{u} \gamma^\mu u$$

$$\cdot \int_0^1 dz \frac{(1-z) ((1-z)m_\lambda + m_b)}{[zm_b^2 + (1-z)m^2 - z(1-z)m_\lambda^2]}$$

$$\int \frac{d^4 k}{(2\pi)^4} = e^2 g^2 \Delta m^2 \int \frac{d^4 k}{(2\pi)^4} \bar{u}(p') \frac{(\not{p} - \not{k} + m_b)}{(\not{p} - \not{k})^2 - m_b^2} \frac{1}{[(\not{k} + \not{q})^2 - m^2]^2}$$

$$\cdot \frac{\not{p}' + m_\lambda}{\not{p}'^2 - m_\lambda^2} \gamma^\mu u(p)$$

combine denominators with:

$$\frac{1}{[(\not{p} - \not{k})^2 - m_b^2][(\not{k} + \not{q})^2 - m^2]^2} = \int_0^1 dz \frac{2(1-z)}{[z(\not{p} - \not{k})^2 - m_b^2 + (1-z)((\not{k} + \not{q})^2 - m^2)]^3}$$

the denominator is:

$$k^2 + 2k \cdot (-zp + (1-z)q) + zp^2 + (1-z)q^2 - (zm_b^2 + (1-z)m^2)$$

$$= k^2 - (zp - (1-z)q)^2 + zp^2 + (1-z)q^2 - (zm_b^2 + (1-z)m^2)$$

$$= k^2 + (z - z^2)p^2 + 2pqz(1-z) + (1-z)zq^2 - (zm_b^2 + (1-z)m^2)$$

$$= k^2 + z(1-z)(p+q)^2 - zm_b^2 + (1-z)m^2$$

$$k = k - zp + (1-z)q$$

$$k = k + zp - (1-z)q$$

$$k-p = k - (1-z)(p+q) = k - (1-z)p'$$

So

$$\int \frac{d^4 k}{(2\pi)^4} \bar{u} \left[\frac{p'(1-z) - k + m_b}{[k^2 - D]^3} \right] = e g^2 \Sigma^2 \int_0^1 dz \, 2(1-z) \frac{p' + m_\mu}{p'^2 - m_\mu^2} u(p)$$

again the k then integrates away. Then we can set $\bar{u} p' = \bar{u} m_e$

$$\text{so} \quad \frac{p' + m_\mu}{p'^2 - m_\mu^2} = \frac{1}{m_e - m_\mu}$$

$$\int \frac{d^4 k}{(2\pi)^4} = (-ie) \frac{g^2 \Sigma^2}{(4\pi)^2} \left(\frac{-1}{m_\mu - m_e} \right) \bar{u} \gamma^4 u \int_0^1 dz \, (1-z) \frac{((1-z)m_e + m_b)}{[z m_b^2 + (1-z)m^2 - 2(1-z)m_e^2]}$$

For the remaining two integrals, we need to combine denominators accordingly to:

$$\frac{1}{[(p-k)^2 - m_b^2] [(k+z)^2 - m^2] [k^2 - m^2]} = \int dx dy dz \delta(x+y+z-1) \frac{2}{[z(p+z)^2 - m_b^2 + x((k+z)^2 - m^2) + y(k^2 - m^2)]^3}$$

$$\frac{1}{[(p-k)^2 - m_b^2] ((k+z)^2 - m^2) (k^2 - m^2)^2} = \int dx dy dz \delta(x+y+z-1) \frac{2 \cdot 3y}{[\dots]^4}$$

$$\frac{1}{[(p-k)^2 - m_b^2] ((k+z)^2 - m^2)^2 (k^2 - m^2)} = \int dx dy dz \delta(x+y+z-1) \frac{2 \cdot 3x}{[\dots]^4}$$

the denominator in all cases is:

$$\begin{aligned}
 & z(p-k)^2 - m_b^2 + x(k+q)^2 - m^2 + y(k^2 - m^2) \\
 &= k^2 + 2k(-zp + xq) + zp^2 + xq^2 - [zm_b^2 + (x+y)m^2] \\
 &= k^2 - (zp - xq)^2 + zp^2 + xq^2 - [zm_b^2 + (1-z)m^2] \\
 &= k^2 + \underbrace{z(1-z)p^2}_{x+y} + 2p_z x_z + \underbrace{x(1-x)q^2}_{y+z} - [zm_b^2 + (1-z)m^2] \\
 &= k^2 + yzp^2 + xz(p+q)^2 + xyq^2 - (zm_b^2 + (1-z)m^2) \\
 &= k^2 - \{ zm_b^2 + (1-z)m^2 - xyq^2 - xzm_c^2 - yzm_a^2 \}
 \end{aligned}$$

$$k = k - zp + xq$$

$$k = k + zp - xq$$

$$k-p = k - \underbrace{(1-z)p}_{x+y} - xq = k - xp' - yP$$

$$2k+q = 2k + 2zp + \underbrace{(1-2x)q}_{y+z-x} = 2k + z(p+p') + (y-x)q$$

then

$$\int_{\gamma_1} + \int_{\gamma_2} = e g^2 S m^2 \int dx dy dz \delta(x+y+z-1) \int \frac{d^4 k}{(2\pi)^4}$$

$$\cdot 6(x+y) \cdot \bar{u}(p') \frac{(xp' + yP - k + m_b) [2k + z(p+p') + (y-x)q]^2}{[k^2 - \Delta]^4}$$

terms linear in k integrate to 0.

$$= (-ie) \frac{g^2 S m^2}{(4\pi)^2} \int_0^1 dx dy dz \delta(x+y+z-1) \cdot \underbrace{(x+y)}_{= (1-z)}$$

$$\left\{ \bar{u}(p') (-\gamma^\mu) \left(\frac{2}{2} g^{\mu\nu} \right) u(p) \frac{1}{\Delta} \right.$$

$$\left. - \frac{1}{\Delta^2} \bar{u}(p) [x m_e + y m_\mu + m_b] u(p) [z(p'+p)^\mu + (y-x) g^\mu] \right\}$$

$$= (-ie) \frac{g^2 S m^2}{(4\pi)^2} \int dx dy dz \delta(x+y+z-1)$$

$$\left\{ \bar{u} \gamma^\mu u \cdot \left(-\frac{1}{\Delta} \right) (1-z) \right.$$

$$+ \bar{u} g^\mu u \cdot (x-y) (x m_e + y m_\mu + m_b) \frac{1}{\Delta^2} (1-z)$$

$$\left. + \bar{u} (p'+p)^\mu u \cdot (-z) (x m_e + y m_\mu + m_b) \frac{(1-z)}{\Delta^2} \right\}$$

$$= (-ie) \frac{g^2 S m^2}{(4\pi)^2} \int dx dy dz \delta(x+y+z-1) (1-z)$$

$$\cdot \left\{ \bar{u} \gamma^\mu u \left(-\frac{1}{\Delta} + \frac{(m_e + m_\mu)(-z)(x m_e + y m_\mu + m_b)}{\Delta^2} \right) \right.$$

$$+ \bar{u} g^\mu u \frac{(x-y)(x m_e + y m_\mu + m_b)}{\Delta^2}$$

$$\left. + \bar{u} i g^{\mu\nu} \gamma_\nu u \frac{z(x m_e + y m_\mu + m_b)}{\Delta^2} \right\}$$

$$\Delta = z m_b^2 + (1-z) m^2 - xy \not{b}^2 - xz m_e^2 - yz m_\mu^2$$

Now combine the three contributions:

$$\text{---} \text{---} = (-ie) \bar{u} [\gamma^\mu F_1 + \gamma^\mu F_3 + i\sigma^{\mu\nu} q_\nu \bar{F}_2] u$$

where

$$F_1 = \frac{g^2 S m^2}{(4\pi)^2} \left\{ \int dz \frac{(1-z)}{m_\lambda - m_e} \left(\frac{(1-z)m_\lambda + m_b}{[zm_b^2 + (1-z)m^2 - z(1-z)m_\lambda^2]} - \frac{(1-z)m_e + m_b}{[zm_b^2 + (1-z)m^2 - z(1-z)m_b^2]} \right) - \int dx dy dz \delta(x+y+z-1) \left(\frac{(1-z)}{\Delta} + \frac{z(1-z)(m_\lambda + m_e)(x m_e + y m_\lambda + m_b)}{\Delta^2} \right) \right\}$$

$$F_3 = \frac{g^2 S m^2}{(4\pi)^2} \int dx dy dz \delta(x+y+z-1) \frac{(x-y)(x m_e + y m_\lambda + m_b)(1-z)}{\Delta^2}$$

$$\bar{F}_2 = \frac{g^2 S m^2}{(4\pi)^2} \int dx dy dz \delta(x+y+z-1) \frac{z(1-z)(x m_e + y m_\lambda + m_b)}{\Delta^2}$$

with $\Delta = z m_b^2 + (1-z)m^2 - xy z^2 - xz m_e^2 - yz m_\lambda^2$

since Δ is positive \rightarrow the region we are considering is $m_b^2, m^2 \gg m_\lambda^2, m_e^2$

these expressions are nonsingular as $g^2 \rightarrow 0$

i.) Now compute $F_1(0)$ in the limit $m_e \rightarrow 0$ then $m_\mu \rightarrow 0$

$$\begin{aligned}
 \overline{F_1(0)} &= \frac{g^2 S m^2}{(4\pi)^2} \left\{ \int_0^1 dz \frac{(1-z)}{m_\mu} \left[\frac{(1-z)m_\mu}{[zm_b^2 + (1-z)m^2]} + \mathcal{O}(m_\mu^2) \right] \right. \\
 &\quad - \int_0^1 dz \int_0^{1-z} dx \frac{1}{[zm_b^2 + (1-z)m^2]} \\
 &\quad \left. - \mathcal{O}(m_\mu) \right\} \\
 &= \frac{g^2 S m^2}{(4\pi)^2} \int_0^1 dz \left\{ \frac{(1-z)^2}{[zm_b^2 + (1-z)m^2]} - \frac{(1-z) \cdot (1-z)}{[zm_b^2 + (1-z)m^2]} \right\} \\
 &= 0
 \end{aligned}$$

j.) (See below)

k.) For $\mu \rightarrow e\gamma$, the photon has $q^2 = 0$. But $F_1(q^2) = 0$ at $q^2 = 0$. F_3 multiplies g^μ , and $g^\mu \cdot \Sigma^\pm(q) = 0$. So neither term contributes to $\mu \rightarrow e\gamma$.

l.) In the limit $m_e \rightarrow 0$ $m_\mu \rightarrow 0$ $g^2 \rightarrow 0$

$$\begin{aligned}\overline{F}_2(s) &= \frac{g^2 \delta m^2}{(4\pi)^2} \int_0^1 dz \int_0^{1-z} dx \frac{z(1-z)m_b}{[zm_b^2 + (1-z)m^2]^2} \\ &= \frac{g^2 \delta m^2}{(4\pi)^2} \int_0^1 dz \frac{z(1-z)^2 m_b}{[zm_b^2 + (1-z)m^2]^2}\end{aligned}$$

Now this is a single integral:

$$\text{II} = \int_0^1 dz \frac{z(1-z)^2}{[zm_b^2 + (1-z)m^2]^2} = \frac{m^4 + 4m^2 m_b^2 - 5m_b^4 + (2m_b^4 + 4m^2 m_b^2) \log \frac{m_b^2}{m^2}}{2(m^2 - m_b^2)^4}$$

so

$$\begin{aligned}\overline{F}_2(s) &= \frac{1}{4\pi} \cdot \frac{g^2}{(4\pi)} \cdot m_b \cdot \left(\frac{\delta m^2}{m^2} \right) \\ &\quad \cdot \frac{1}{2(1 - m_b^2/m^2)^4} \left(1 + 4 \frac{m_b^2}{m^2} - 5 \frac{m_b^4}{m^4} + 2 \left(\frac{m_b^4}{m^4} + 2 \frac{m_b^2}{m^2} \right) \log \frac{m_b^2}{m^2} \right)\end{aligned}$$

which is of the form

$$\overline{F}_2(s) = G \frac{\alpha}{4\pi} \cdot m_b \frac{1}{m^2}$$

$$\text{with } G \sim \left(\frac{\delta m^2}{m^2} \right)$$

now there is one unfinished piece of business:
the extra-credit part (j)

We need to simplify the expression for $F_1(g)$ on p. 16

First, let $x = (1-z)\xi$ $y = (1-z)(1-\xi)$

$$\int dx dy dz \delta(x+y+z-1) = \int_0^1 dz \int_0^{1-z} dx = \int_0^1 dz \int_0^1 d\xi (1-z)$$

now take the term

$$\int dx dy dz \delta(x+y+z-1) \frac{(1-z)}{\Delta}$$

$$= \int_0^1 dz \int_0^1 d\xi \frac{(1-z)^2}{[zm_b^2 + (1-z)m^2 - z(1-z)(\xi m_e^2 + (1-\xi)m_\mu^2) - (1-z)^2 \xi(1-\xi)g^2]}$$

and integrate by parts in ξ :

$$= \int_0^1 dz \left\{ \frac{(1-z) [(1-z)(\xi m_e + (1-\xi)m_\mu) + m_b]}{[zm_b^2 + (1-z)m^2 - z(1-z)(\xi m_e^2 + (1-\xi)m_\mu^2) - (1-z)^2 \xi(1-\xi)g^2]} \right\} \Bigg|_{\xi=0}^{\xi=1}$$

$$= \int_0^1 dz \frac{(1-z)}{m_e - m_\mu} \left\{ \frac{[(1-z)(\xi m_e + (1-\xi)m_\mu) + m_b] [z(1-z)(m_e^2 - m_\mu^2) + (1-z)^2 (1-\xi)g^2]}{[zm_b^2 + (1-z)m^2 - z(1-z)(\xi m_e^2 + (1-\xi)m_\mu^2) - (1-z)^2 \xi(1-\xi)g^2]^2} \right\}$$

then give a nice identity:

$$\int_0^1 dz \int_0^{1-z} dx \frac{(1-z)}{\Delta}$$

$$= \int_0^1 dz \frac{(1-z)}{(m_\mu - m_e)} \left\{ \frac{(1-z)m_\mu + m_b}{[zm_b^2 + (1-z)m^2 - z(1-z)m_\mu^2]} - \frac{((1-z)m_e + m_b)}{[2m_b^2 + (1+z)m^2 - z(1-z)m_e^2]} \right\}$$

cancels the first term
of F_1

$$- \int_0^1 dz \int_0^{1-z} dx (1-z) \left\{ \frac{(xm_e + ym_\mu + m_b)}{\Delta^2} \right\}$$

$$\cdot \left\{ \frac{1}{m_\mu - m_e} [z(1-z)(m_\mu^2 - m_e^2) - (1-z)(y-x)q^2] \right\}$$

$$= z(1-z)(m_\mu + m_e)$$

so!

$$\frac{F_1(q^2)}{[q^2 \delta m^2 (4\pi)^2]} = \int_0^1 dz \int_0^{1-z} dx (1-z) \frac{q^2}{(m_\mu - m_e)} \frac{(x-y)(xm_e + ym_\mu + m_b)}{\Delta^2}$$

$$= \frac{q^2}{m_\mu - m_e} \int dx dy dz \delta(x+y+z-1) \frac{(x-y)(xm_e + ym_\mu + m_b)}{\Delta^2}$$

$$\therefore F_1(q^2) = \frac{q^2}{m_\mu - m_e} F_3(q^2) \quad \text{Q.E.D.}$$